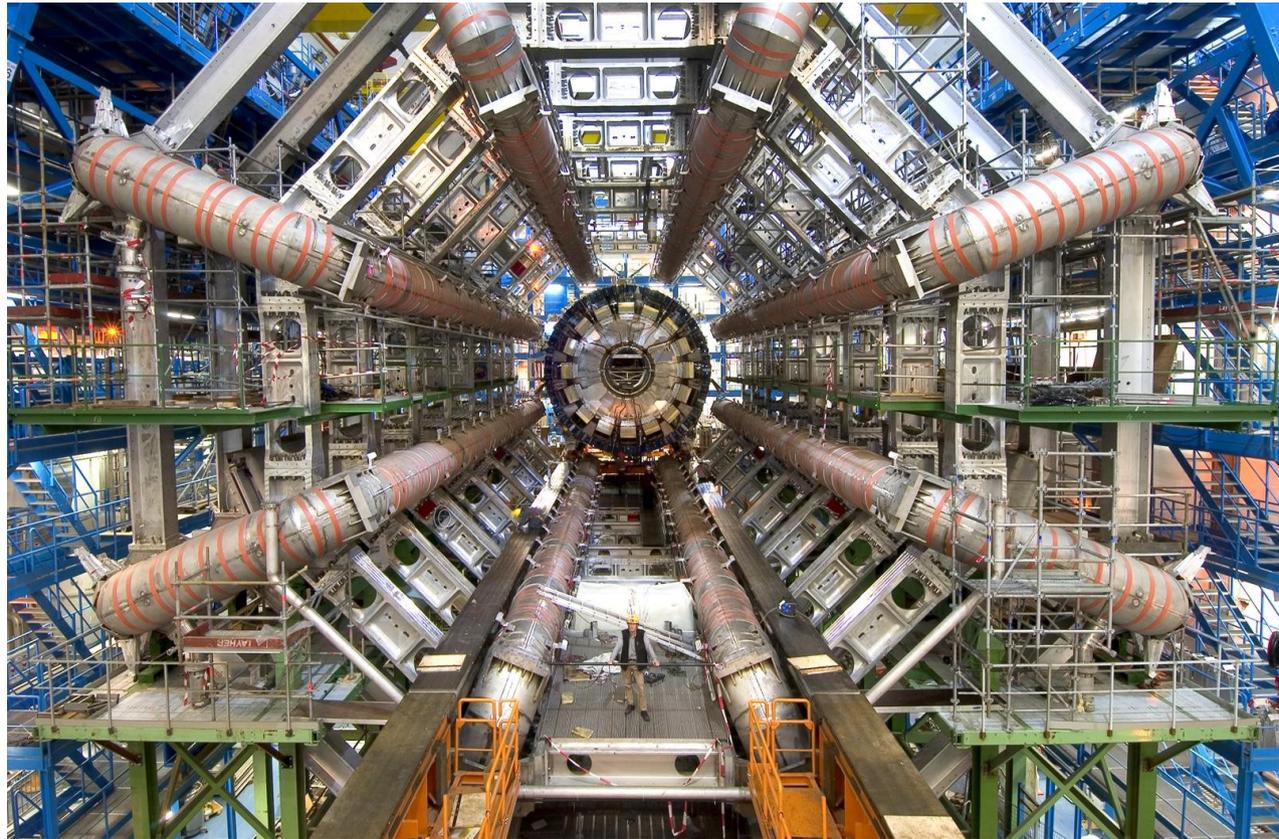
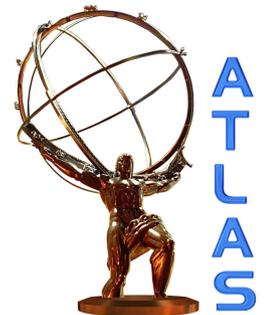


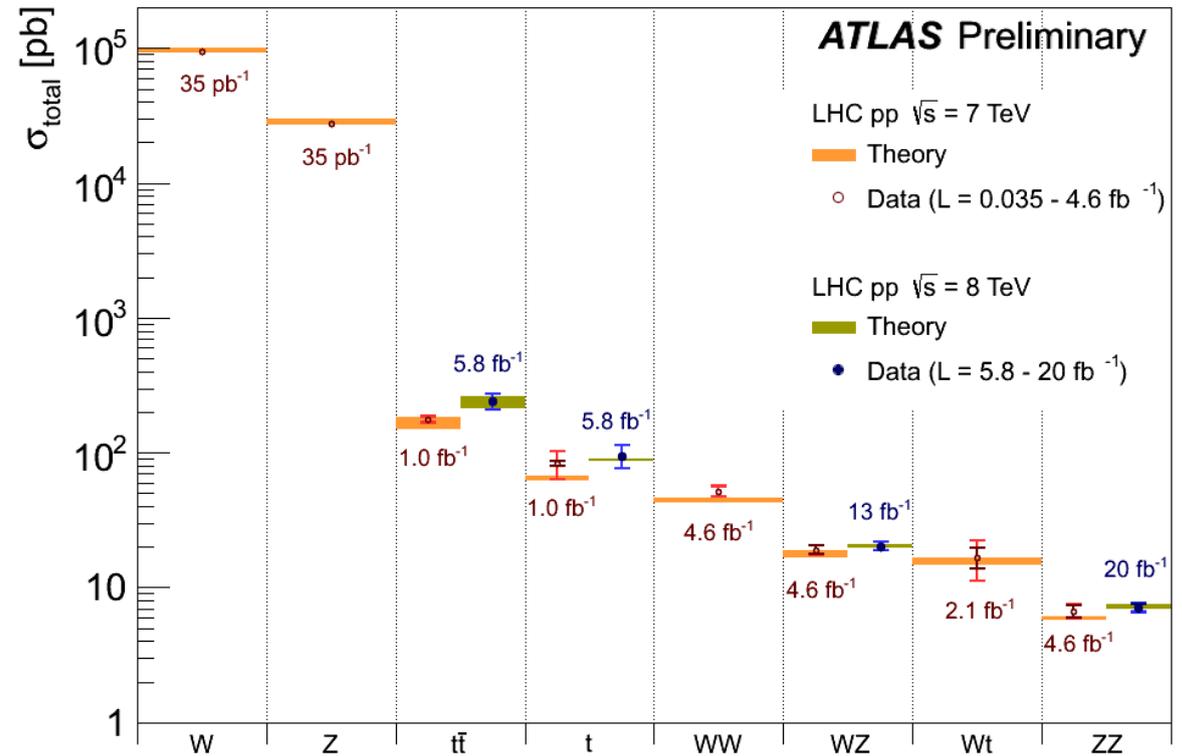
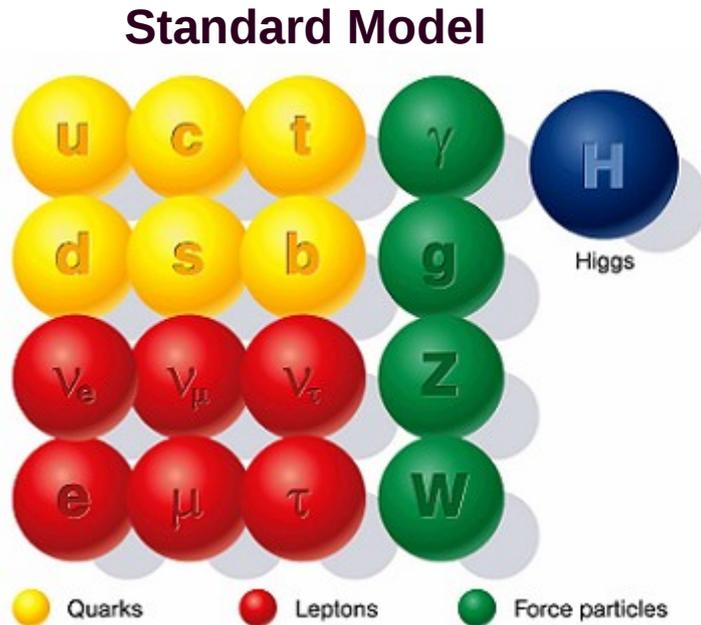
Supersymmetry searches with ATLAS: overview and latest results



Tina Potter
On behalf of the ATLAS Collaboration



The Standard Model



The Standard Model of elementary particles is a very successful theory.

Precise predictions, verified by experiment over many orders of production rate.

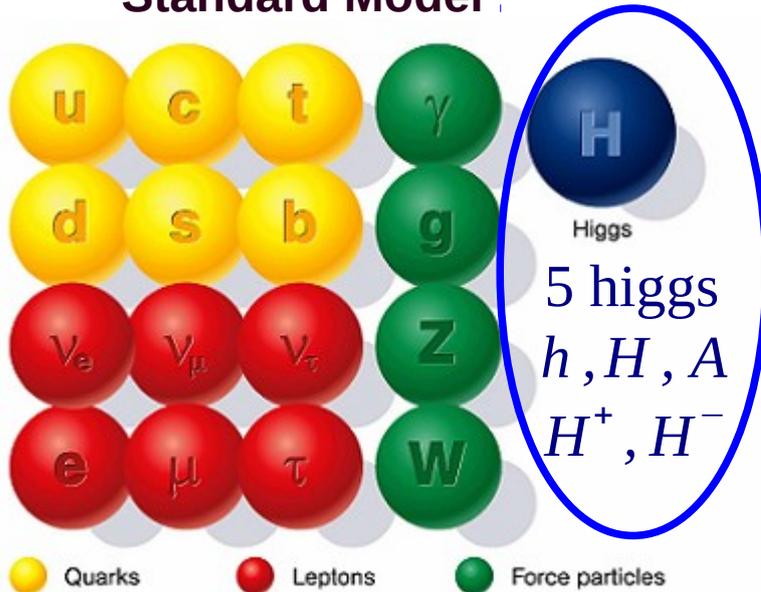
So what's the problem?

- Need high-levels of fine tuning to avoid quadratic divergences in Higgs mass corrections
- No explanation for Dark Matter
- No unification of the forces

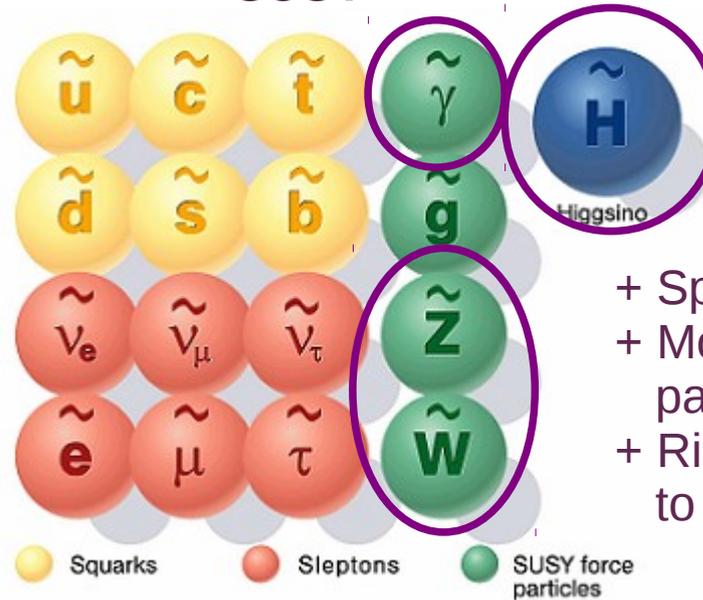
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Supersymmetry

Standard Model



SUSY

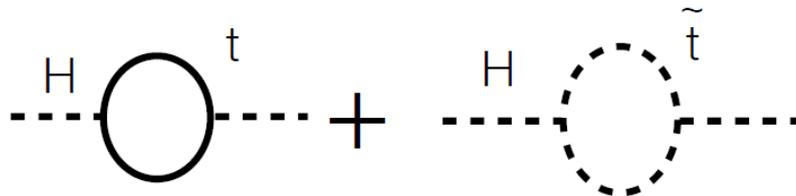


Superpartner for every SM particle

- + Spin differs by one half
- + Mostly heavier than SM partners \rightarrow broken symmetry
- + Rich array of signatures to search for at the LHC

$$W^\pm, W^0, B \xrightarrow{\text{mixing}} W^\pm, Z, \gamma$$

$$\begin{aligned} \tilde{H}_u^0, \tilde{H}_d^0, \tilde{W}^0, \tilde{B}^0 &\xrightarrow{\text{mixing}} \tilde{\chi}_1^0, \tilde{\chi}_2^0, \tilde{\chi}_3^0, \tilde{\chi}_4^0 && \text{neutralinos} \\ \tilde{H}_u^+, \tilde{H}_d^-, \tilde{W}^+, \tilde{W}^- &\xrightarrow{\text{mixing}} \tilde{\chi}_1^\pm, \tilde{\chi}_2^\pm && \text{charginos} \end{aligned}$$



New superpartner loop roughly cancels the SM loop

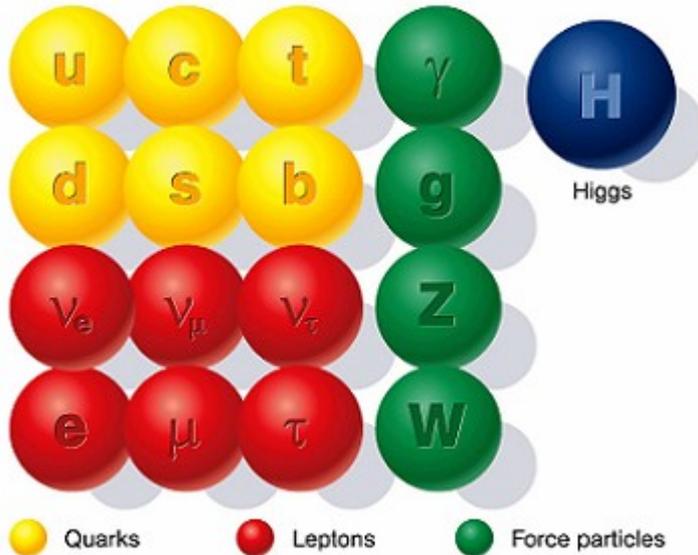
So what's the problem?

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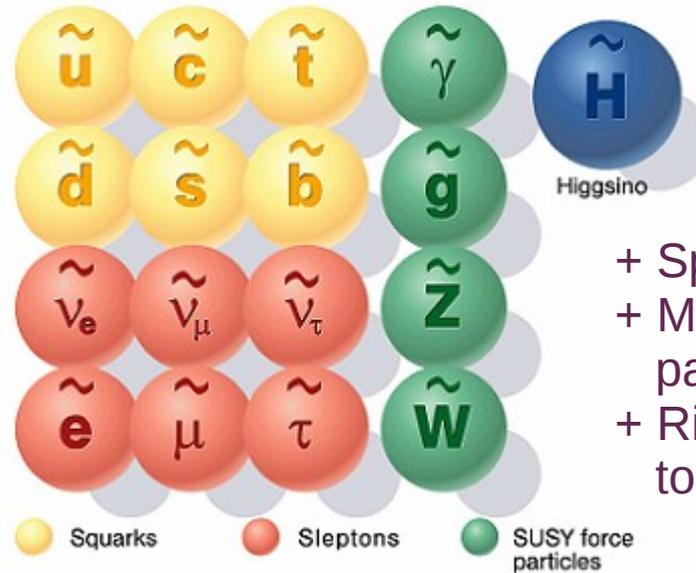
...

Supersymmetry

Standard Model



SUSY



Superpartner for every SM particle

- + Spin differs by one half
- + Mostly heavier than SM partners → broken symmetry
- + Rich array of signatures to search for at the LHC

If R-parity is conserved, lightest SUSY particle (LSP) is stable

$$P_R = (-1)^{3(B-L)+2S}$$

+1 for SM particles, -1 for SUSY particles

In many models, LSP is commonly lightest neutralino $\tilde{\chi}_1^0$

→ good dark matter candidate!

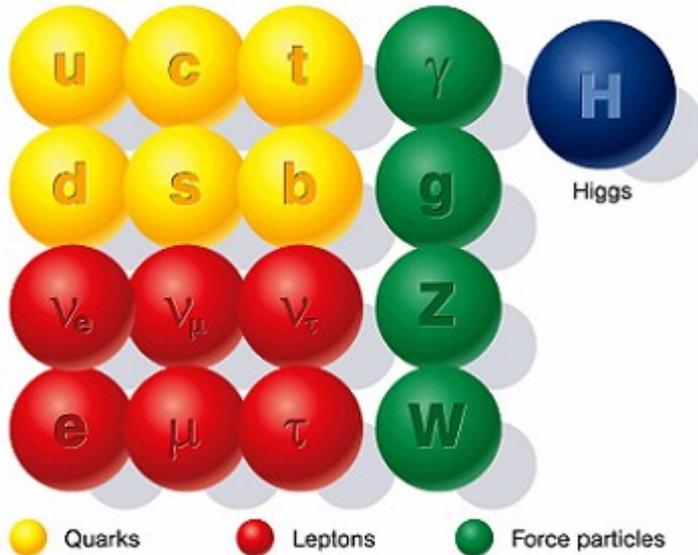
So what's the problem?

- ✓ – Need high-levels of fine tuning to avoid quadratic divergences in Higgs mass corrections
- ✓ – No explanation for Dark Matter
- No unification of the forces

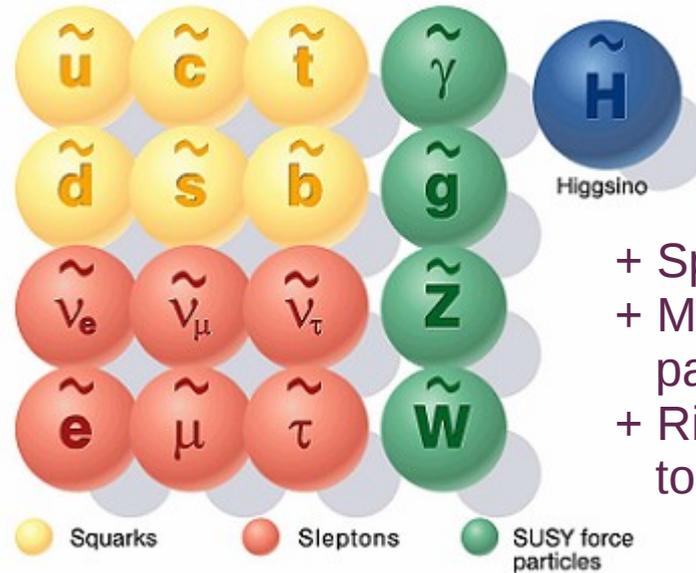
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Supersymmetry

Standard Model

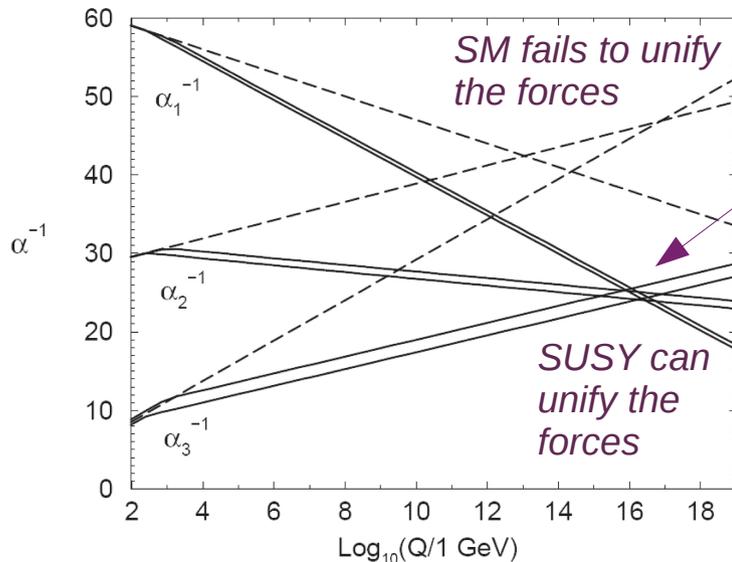


SUSY



Superpartner for every SM particle

- + Spin differs by one half
- + Mostly heavier than SM partners → broken symmetry
- + Rich array of signatures to search for at the LHC



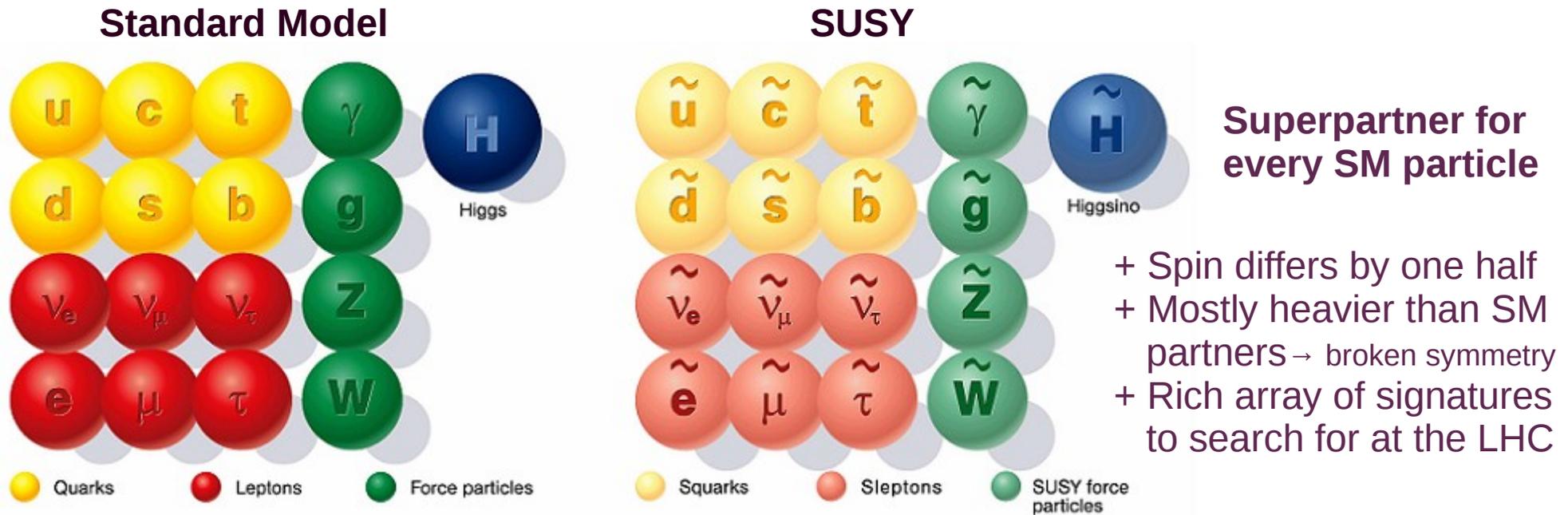
High unification scale also resolves issues with proton decay

So what's the problem?

- ✓ – Need high-levels of fine tuning to avoid quadratic divergences in Higgs mass corrections
- ✓ – No explanation for Dark Matter
- ✓ – No unification of the forces
- ...

arXiv:hep-ph/9709356

Supersymmetry



After the discovery of the higgs with mass ~ 125 GeV, there is renewed interest in accomodating higgs in SUSY models to help stabilise the higgs mass. To avoid high levels of “unnatural” fine tuning, some sparticles need to be light.

$$\frac{m_H^2}{2} = -|\mu|^2 + \dots + \delta m_H^2$$

$$\delta m_H^2|_{stop} \cong -\frac{3y_t^2}{8\pi^2} (m_{Q_3}^2 + m_{U_3}^2 + |A_t|^2) \ln\left(\frac{\Lambda}{TeV}\right)$$

$$\delta m_H^2|_{gluino} \cong -\frac{2y_t^2}{\pi^2} \left(\frac{\alpha_s}{\pi}\right) |M_3|^2 \ln^2\left(\frac{\Lambda}{TeV}\right)$$

For natural SUSY (low levels of fine tuning)

→ Light higgsinos

→ Light stop (< 1 TeV)

→ Light gluinos ($< 1-2$ TeV)

Weak-scale SUSY needed for naturalness should be seen at the LHC!

Supersymmetry

R-Parity $P_R = (-1)^{3(B-L)+2S}$

B, L, S : baryon, lepton, spin
 +1 for SM particles, -1 for SUSY particles

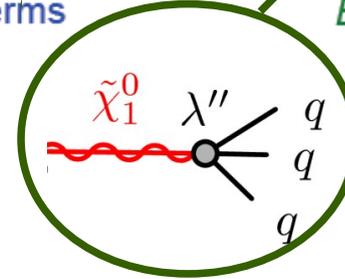
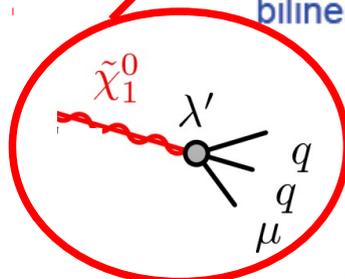
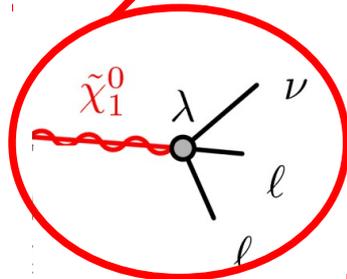
No reason to assume conservation of R-parity
 Can constrain proton decay with lepton or baryon violating SUSY, but not both
 LSP decays \rightarrow no dark matter candidate

The MSSM potential

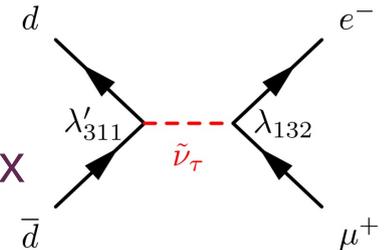
$$W_{RP} = \lambda_{ijk} \hat{L}_i \hat{L}_j \hat{E}_k^C + \lambda'_{ijk} \hat{L}_i \hat{Q}_j \hat{D}_k^C + \epsilon_i \hat{L}_i \hat{H}_u + \lambda''_{ijk} \hat{U}_i^C \hat{D}_j^C \hat{D}_k^C$$

L-number violating terms (bracketed over the first two terms)
bilinear terms (pointing to the third term)
B-number violating terms (bracketed under the fourth term)

RPV couplings $\lambda, \lambda', \lambda''$



RPV can be at the production vertex and/or at decay vertices



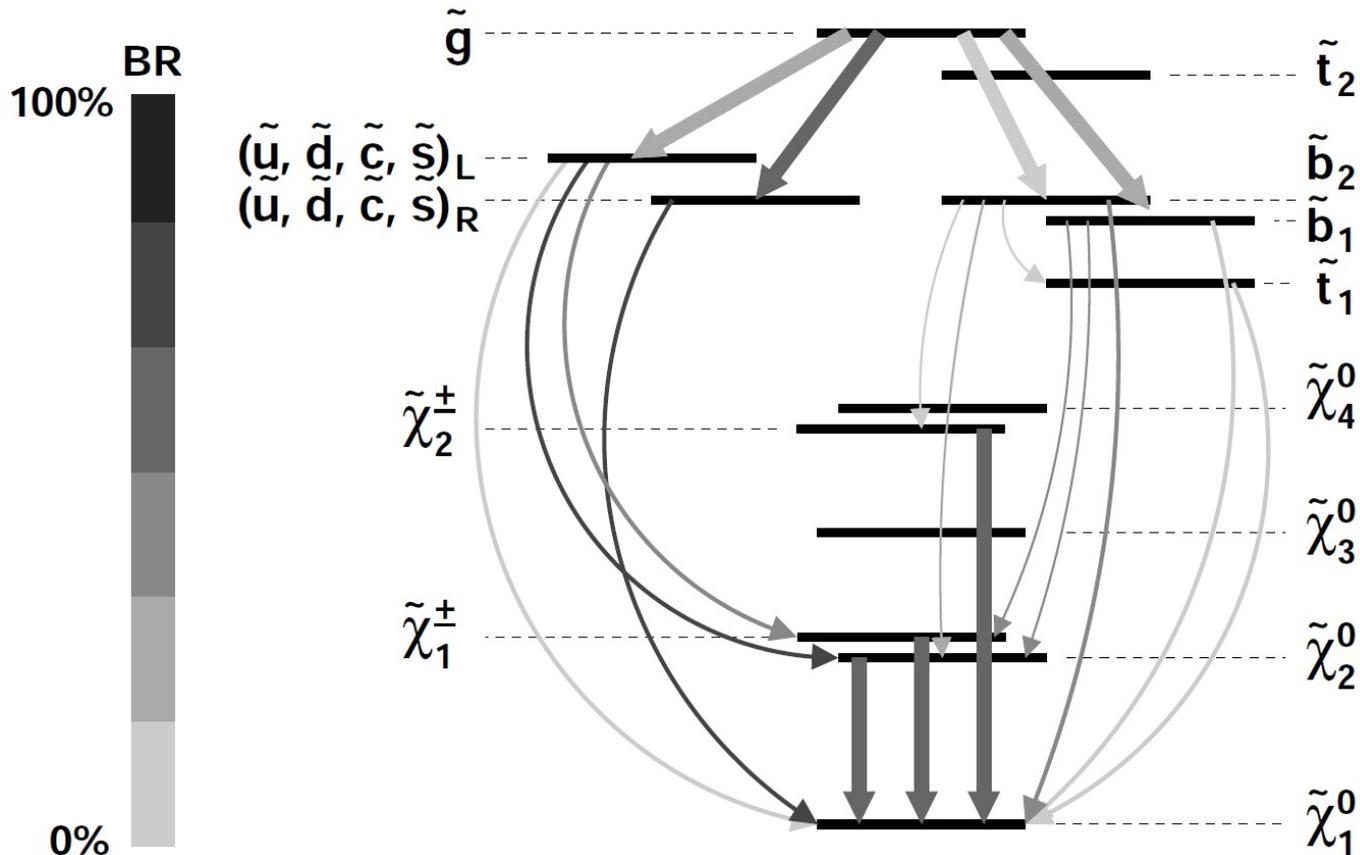
RPV also lead to **non-prompt decays** if λ couplings are small

Long-lived SUSY particles can also arise from

- Heavy mediator sparticles e.g. Split SUSY
- Mass degeneracy
- Weak couplings

Supersymmetry models

Physics model e.g. MSUGRA/CMSSM

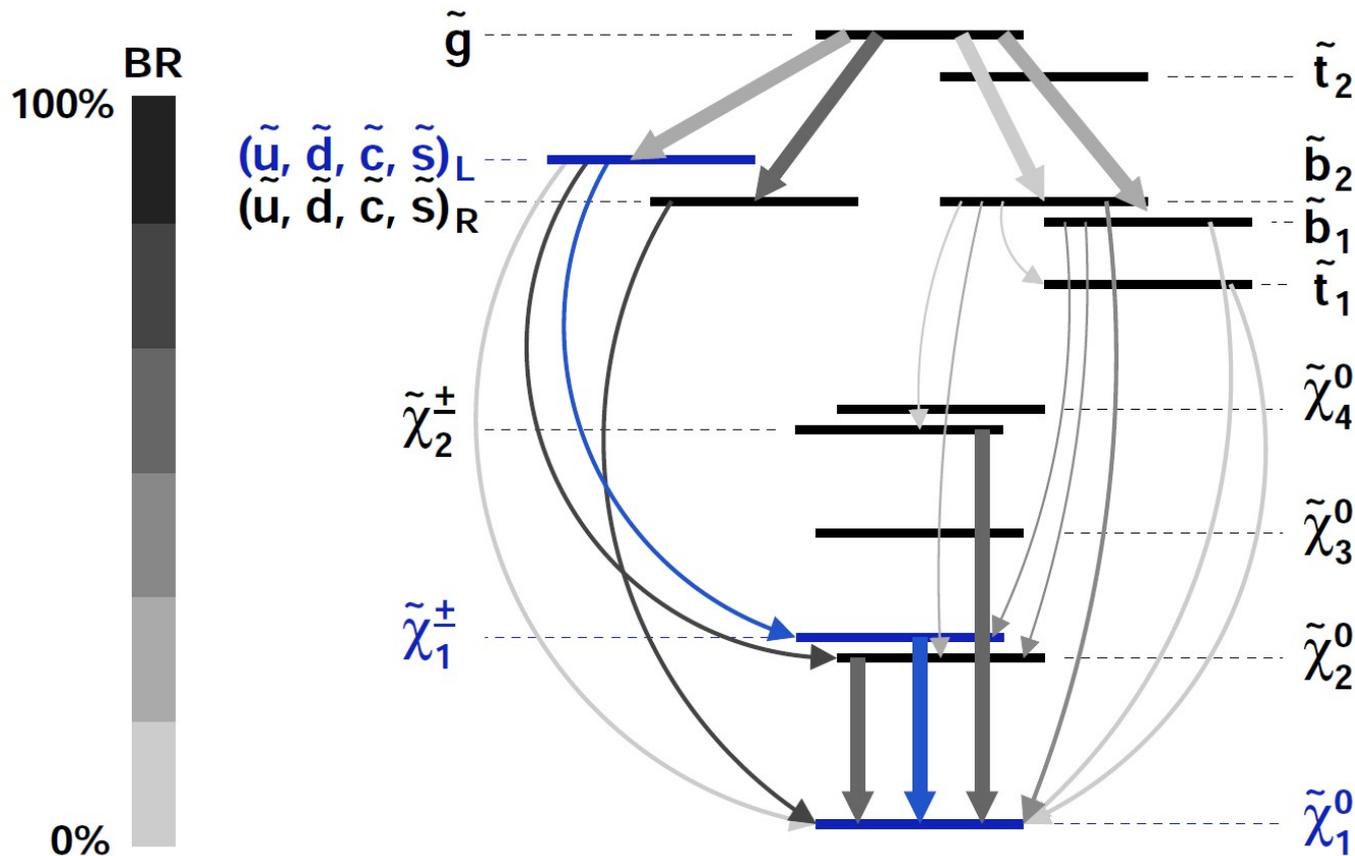


A typical SUSY spectrum involves
+ many sparticles with different masses
+ many different possible ways for each to decay

Where do we start looking?

Supersymmetry models

Physics model e.g. MSUGRA/CMSSM

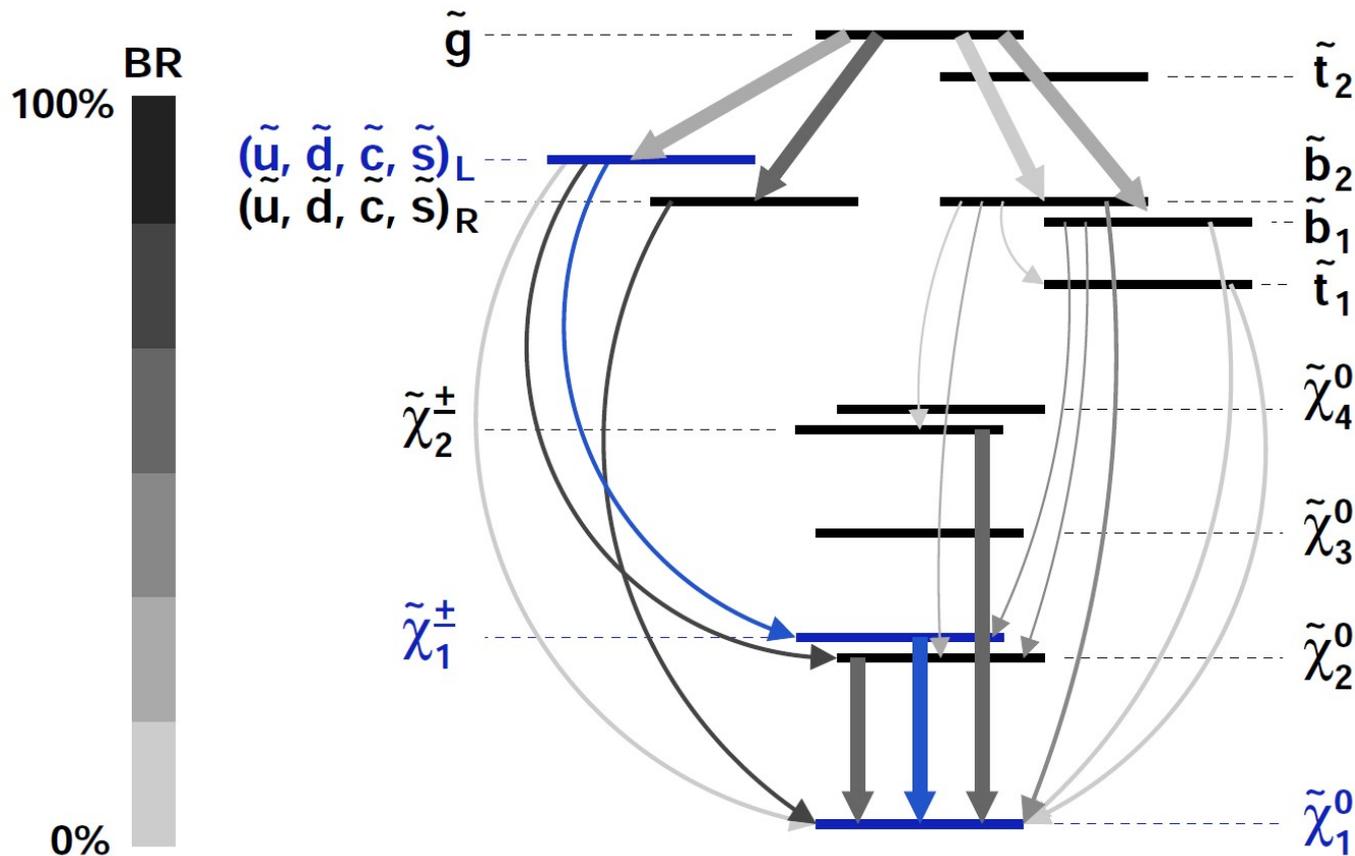


A typical SUSY spectrum involves
+ many sparticles with different masses
+ many different possible ways for each to decay

Focus on process of interest

Supersymmetry models

Physics model e.g. MSUGRA/CMSSM



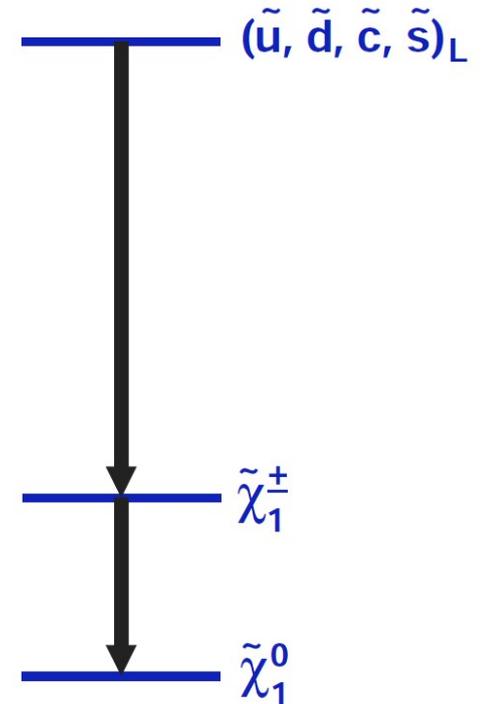
A typical SUSY spectrum involves

- + many sparticles with different masses
- + many different possible ways for each to decay

Focus on process of interest

Simplified model

Study a specific decay chain



- + Small number of sparticles, assumed BR – usually 100%.
- + Described by masses and cross-sections.
- + Simple and broad approach for designing SUSY searches.

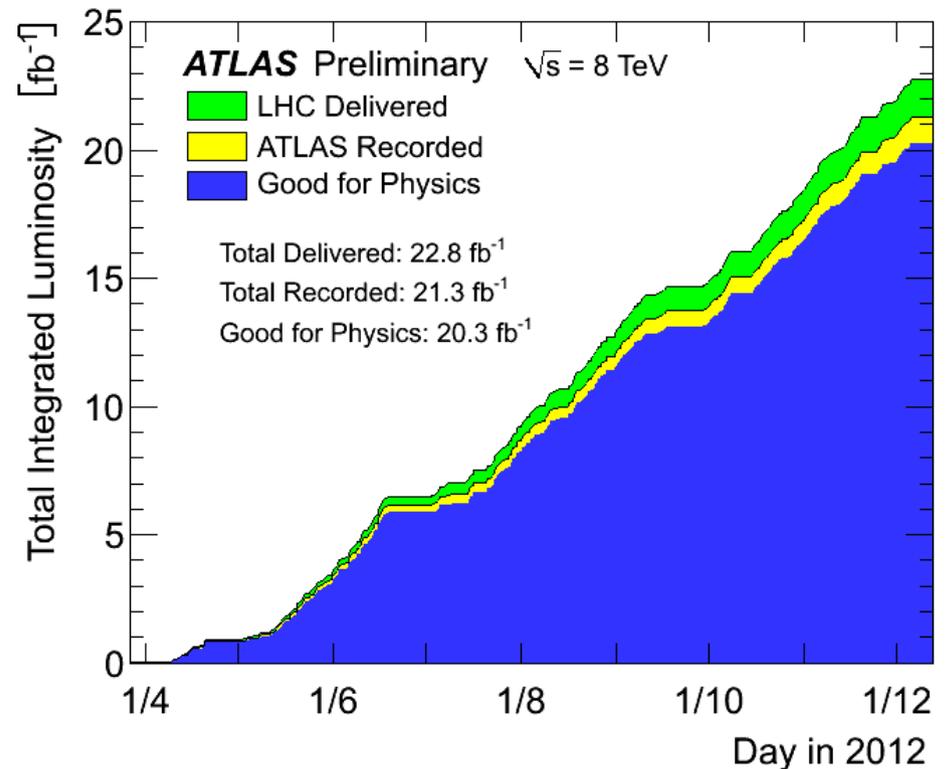
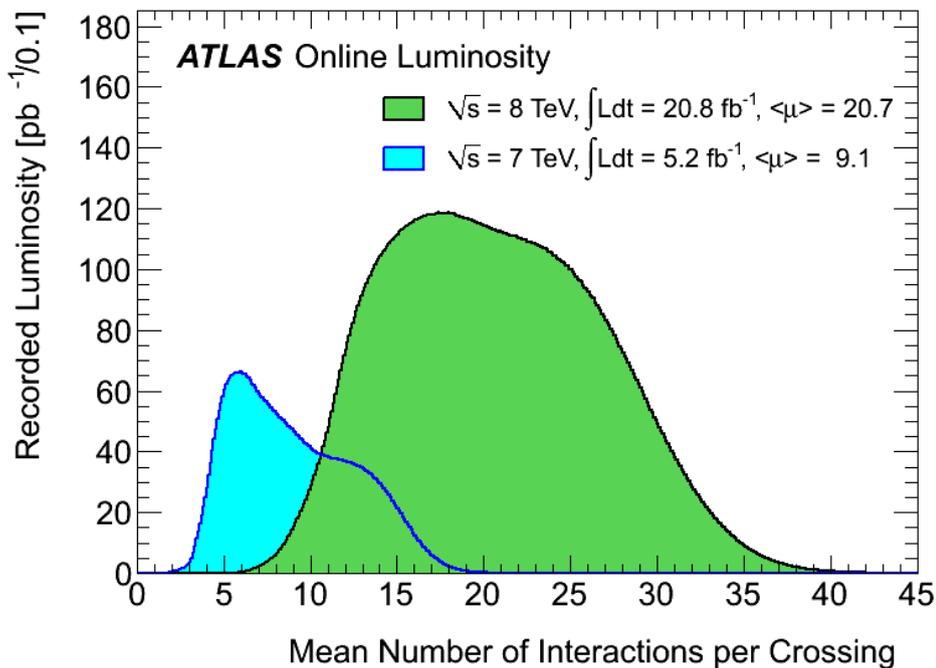
Experimental setup: Luminosity and pileup

~ 22 fb⁻¹ collected at $\sqrt{s} = 8$ TeV

~ 5 fb⁻¹ collected at $\sqrt{s} = 7$ TeV

with ~ 90% of the delivered data being good for physics

Most results presented here use the full 8 TeV dataset



Large luminosity results in large pileup (number of interactions per bunch crossing)

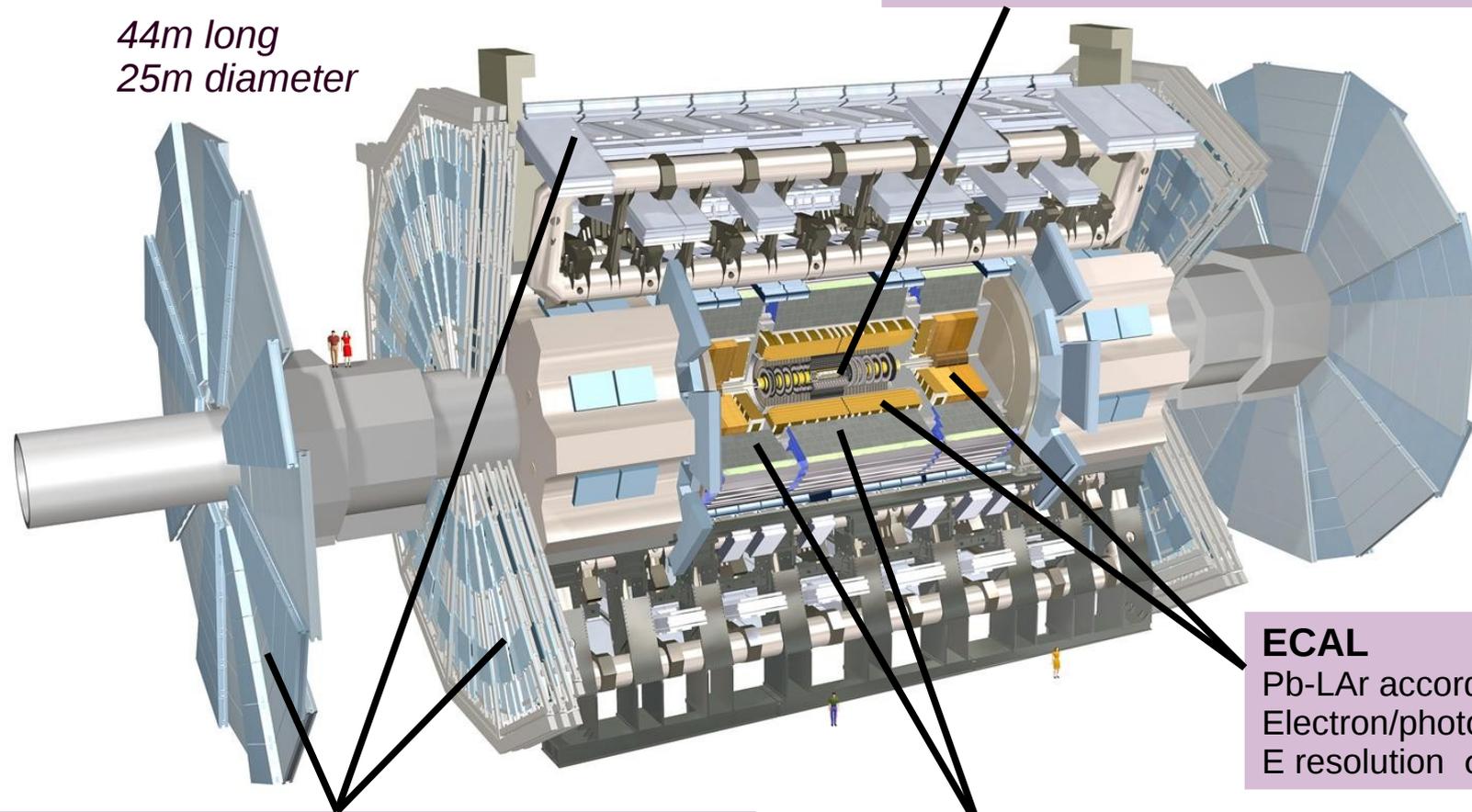
Pileup suppression strategies have been carefully developed

Experimental setup: The ATLAS detector

3-level trigger

Rate 40 MHz \rightarrow \sim 400 Hz

44m long
25m diameter



Inner Detector

Silicon pixels & strips + TRT straws

Precise tracking and vertexing, electron/pion separation

p resolution $\sigma/p_T \sim 3.8 \times 10^{-4} p_T$ (GeV) \oplus 0.015

ECAL

Pb-LAr accordion

Electron/photon id & measurement

E resolution $\sigma/E \sim 10\%/\sqrt{E}$

Muon Spectrometer

Air-core toroids with gas-based muon chambers

Muon measurement

p resolution $\sigma/p < 10\%$ up to $p \sim 1$ TeV

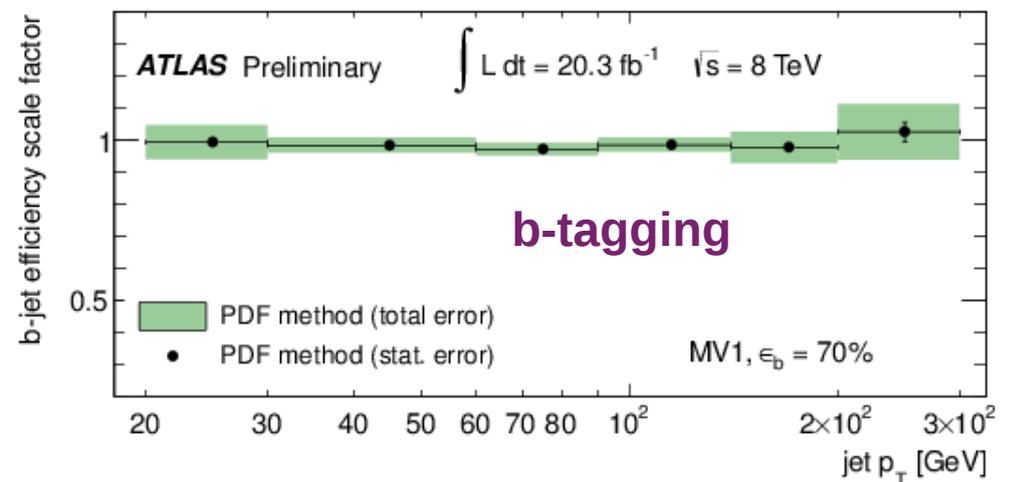
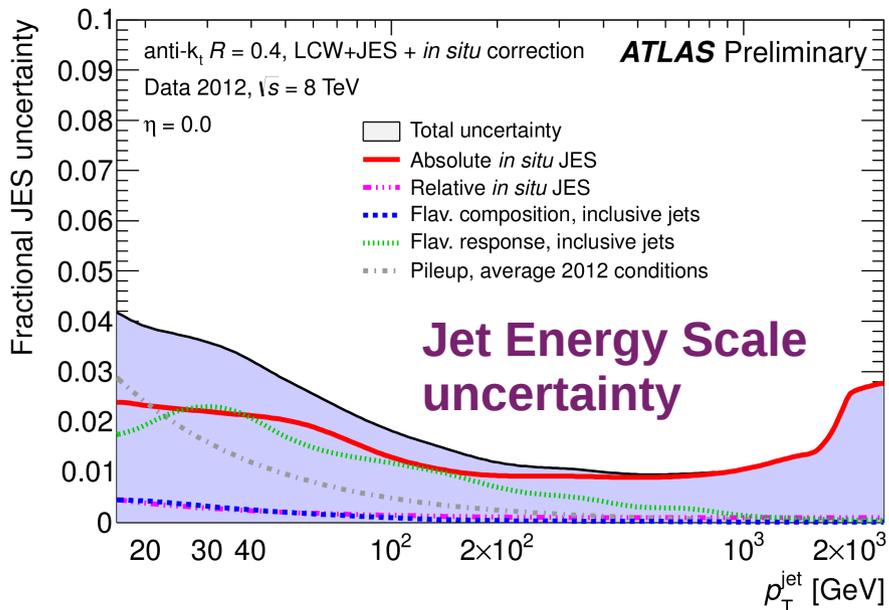
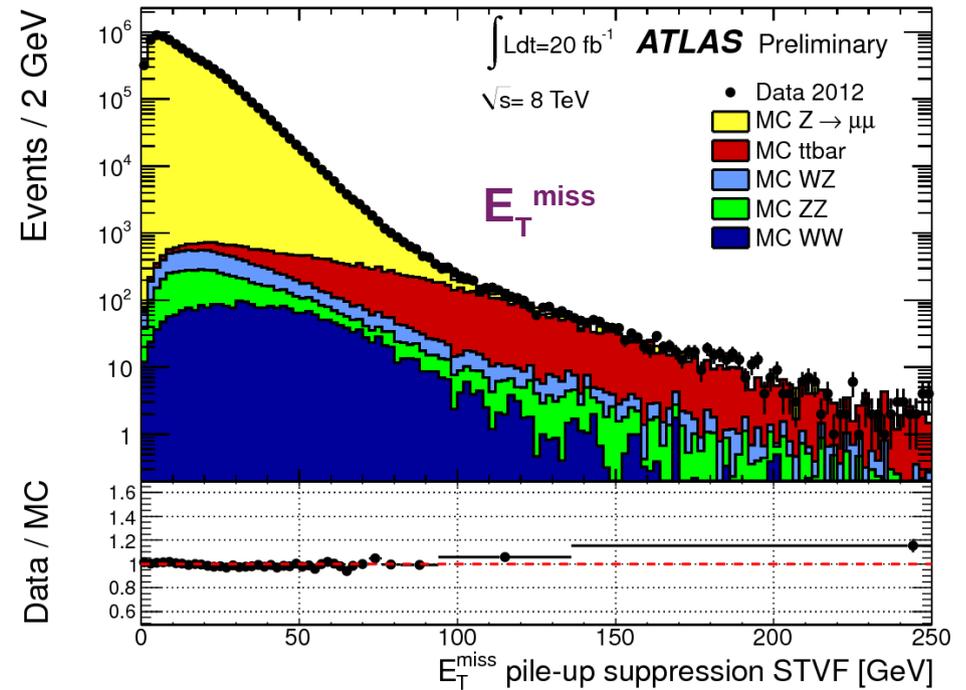
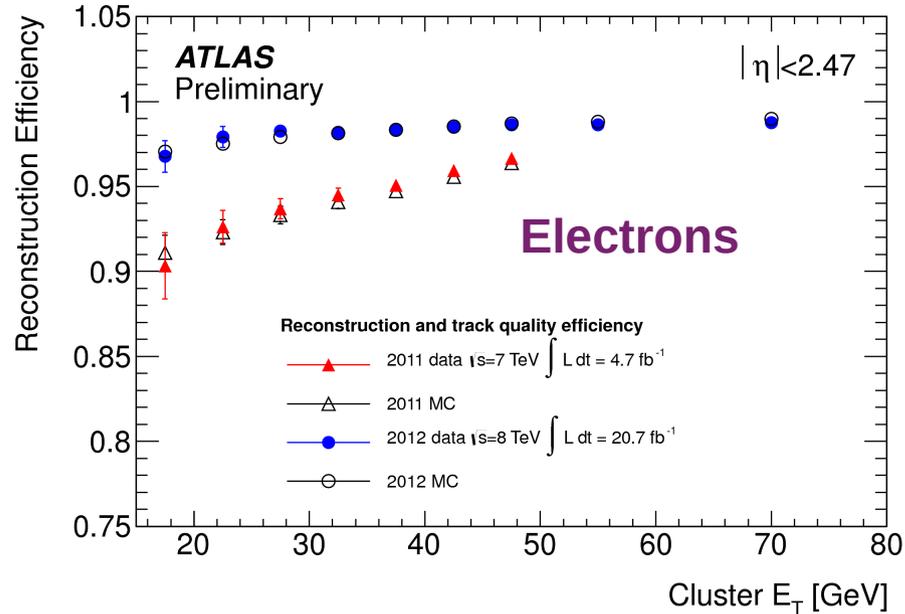
HCAL

Fe/scintillator tiles (central), Cu/Q-LAr (fwd)

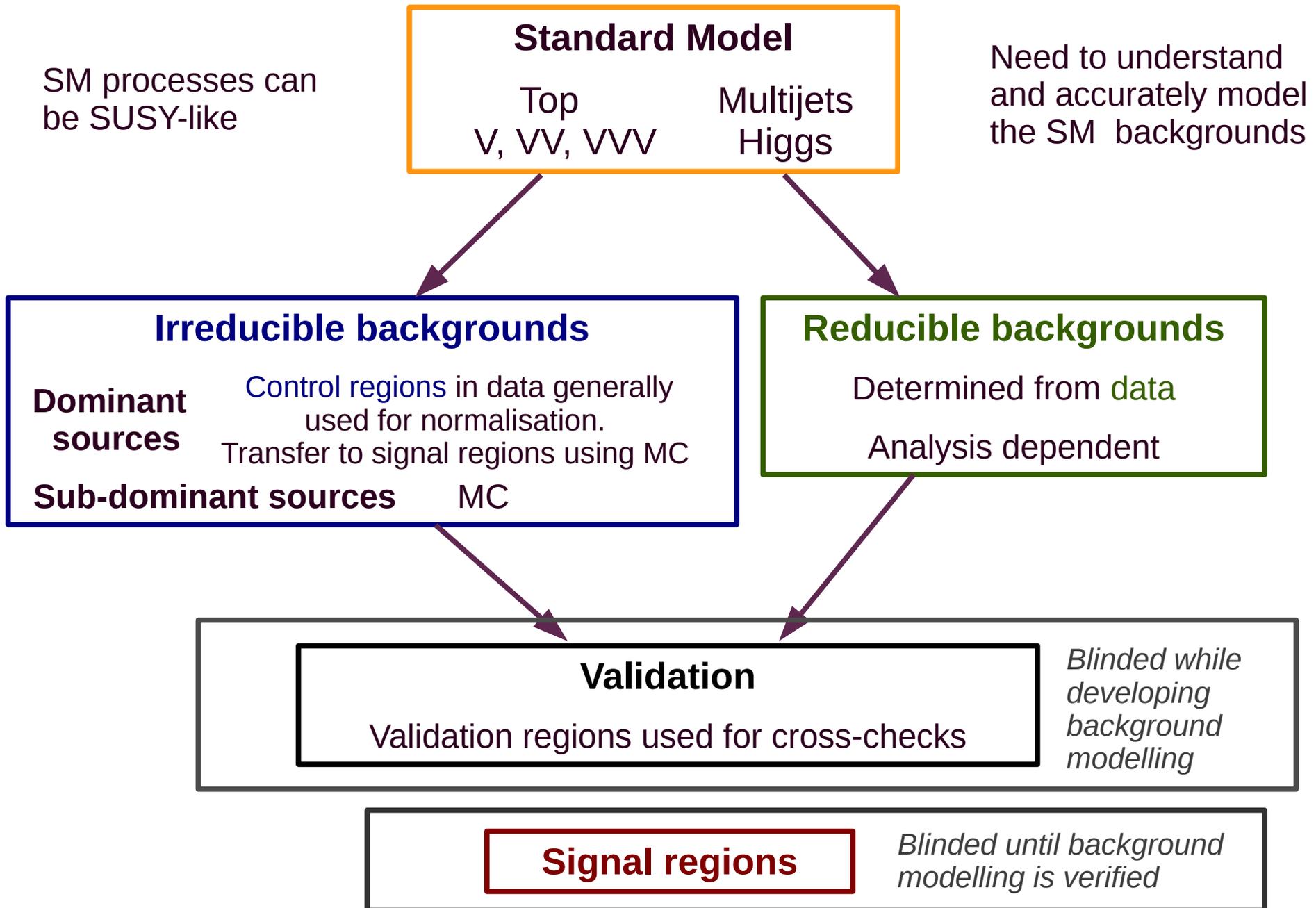
Measurement of jets & missing E_T

E resolution $\sigma/E \sim 50\%/\sqrt{E} \oplus 0.03$

Experimental setup: Object Performance



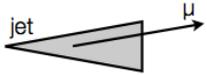
Standard Model Background Modelling



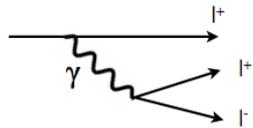
Reducible background determination

Fake leptons

HF/LF fakes from jets



Conversion leptons from photon radiation



A few examples....

Large, fake E_T^{miss} can be induced by a jet mis-measurement

Matrix Method

Estimate from data: set of linear equations relating the kinematic properties of the leptons to the real and fake lepton composition of the data sample

e.g. for a 1-lepton signature

Define “loose” (pre-selected) lepton

Define “tight” (signal) lepton

Solve set of equations

$$N^{\text{loose}} = N_{\text{real}}^{\text{loose}} + N_{\text{fake}}^{\text{loose}}$$

$$N^{\text{tight}} = \epsilon_{\text{real}} N_{\text{real}}^{\text{loose}} + \epsilon_{\text{fake}} N_{\text{fake}}^{\text{loose}}$$

Measure independently from data

Count in data

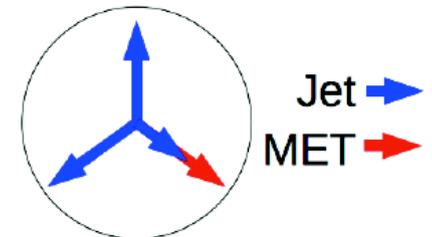
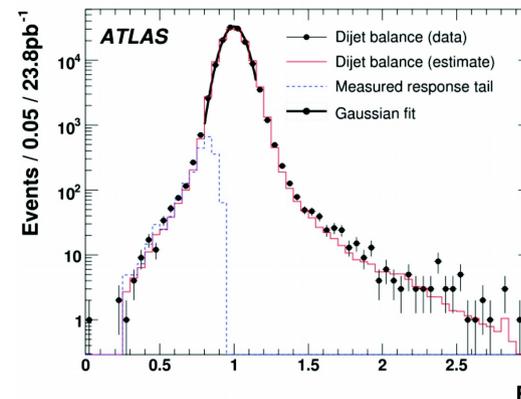
$$N_{\text{fake}}^{\text{tight}} = \frac{\epsilon_{\text{fake}}}{\epsilon_{\text{real}} - \epsilon_{\text{fake}}} (\epsilon_{\text{real}} N_{\text{loose}} - N_{\text{tight}})$$

Jet Smearing

Derive a “jet response function” from MC and adapt it to data

core: p_T balance in di-jet events

tail: three-jet (Mercedes) events



Use response function to smear jets in real data events with low E_T^{miss}

→ Obtain events with large “fake” E_T^{miss}

Irreducible background determination using control regions

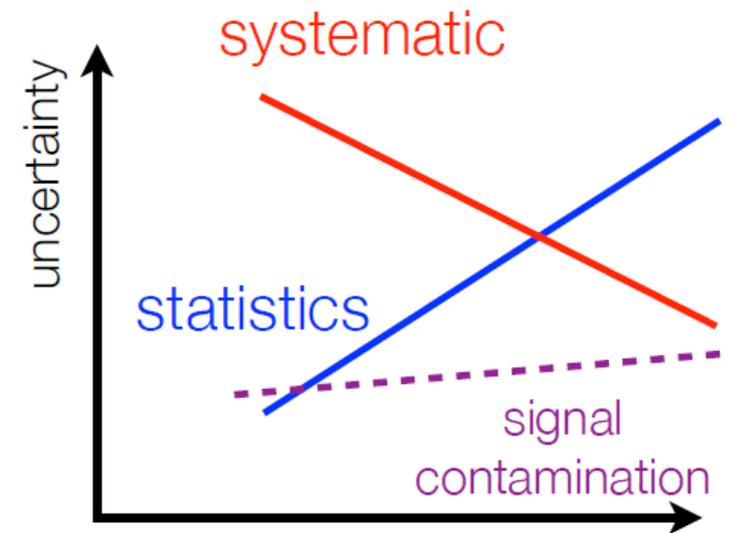
Irreducible backgrounds normalised to data in dedicated control regions

$$N_{SR}^i = \frac{N_{SR}^{i,MC}}{N_{CR}^{i,MC}} \left(N_{CR}^{i,data} - \sum_{j=process} N_{CR}^{j,MC} \right)$$

MC
Count in data
MC or other CR

“Transfer factor”
MC or other CR

If contamination from other processes is small, all systematic uncertainty is associated to transfer factor



Closeness to signal region

Need to be careful in the choice of CR

Typical uncertainties

Experimental

- Trigger efficiency
- Jet energy scale, resolution
- Lepton energy scale, efficiency
- E_T^{miss} soft component
- b-tagging
- Luminosity
- pileup modelling

Theory

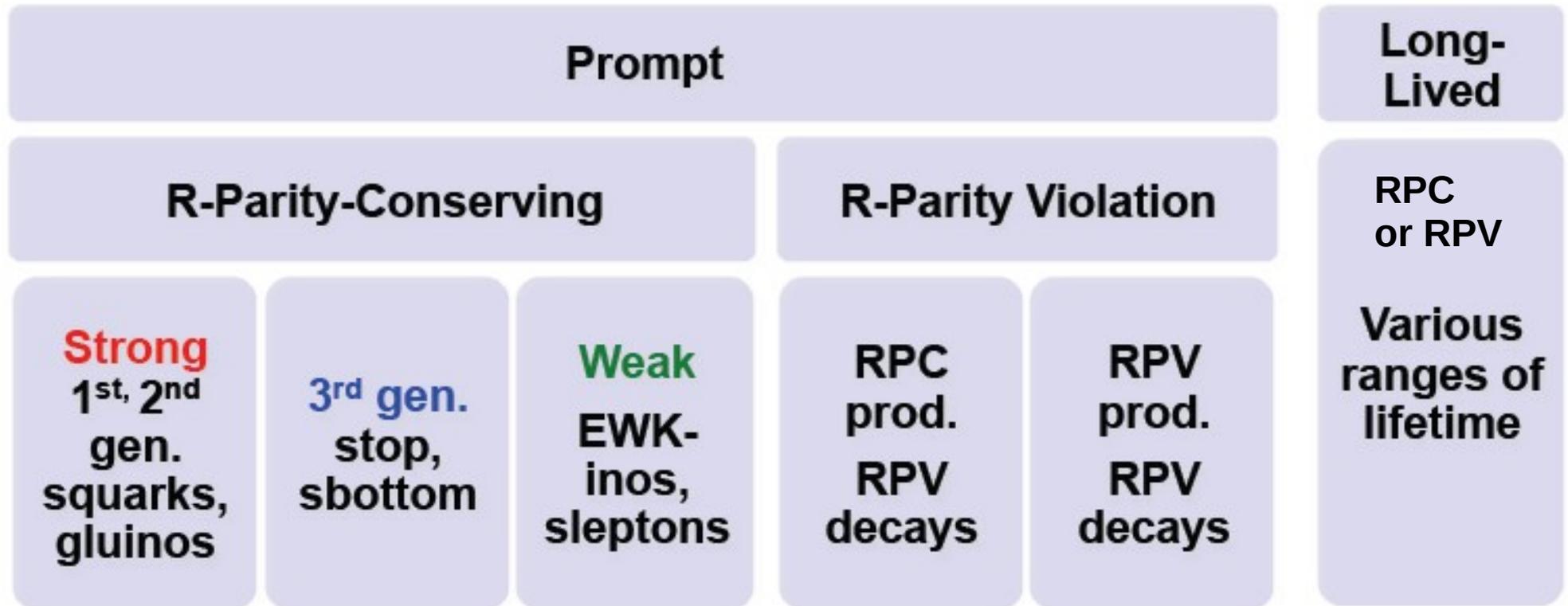
- Generator modelling
 $\mu_F, \mu_R,$
ME/PS matching,
 α_s scale choice
- PS uncertainties
typically compare Pythia
and Herwig
- PDF choice

Background determination **verified in validation regions**

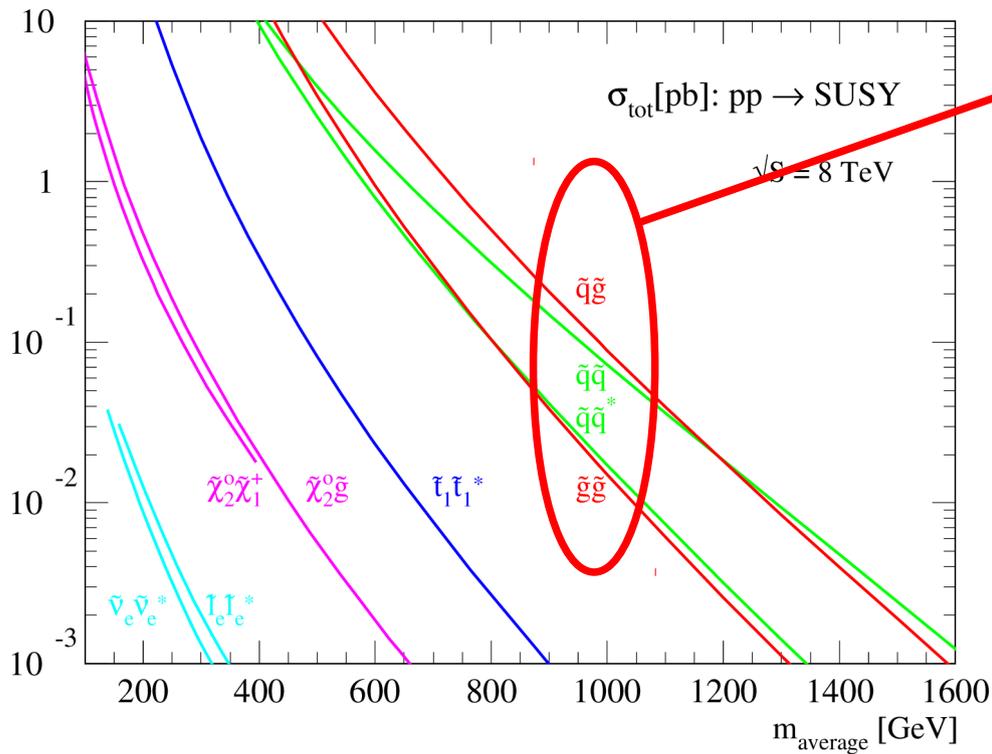
Calculation done **performing a combined fit** to all control and signal regions.
Account for signal contamination for exclusion

Search Strategy

Search strategy designed to provide coverage for a broad class of SUSY models



For each search, a number of signal regions is optimised based on a variety of models



High cross-section for strongly produced SUSY particles at the LHC

Large yield even in small datasets

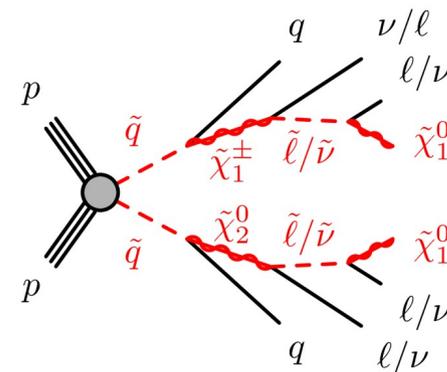
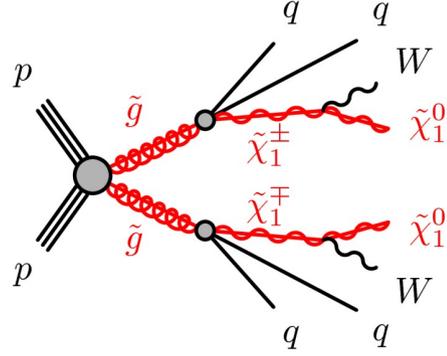
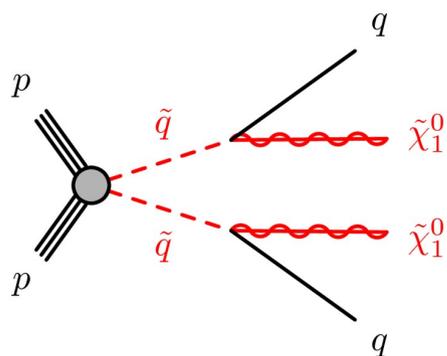
Dominates total SUSY cross-section in many SUSY models

e.g. MSUGRA/CMSSM

Inclusive search channels

- » 0-lepton + 2-6 jets + E_T^{miss}
- » 0-lepton + 7-10 jets + $E_T^{\text{miss sig}}$
- » 1-2 leptons + jets + E_T^{miss}
- » 2-lepton + 2-6 jets + E_T^{miss}
- 1-2 taus + jets + E_T^{miss} ATLAS-CONF-2013-026
- Z(l) + jets + E_T^{miss} ATLAS-CONF-2013-152
- Photon + lepton + E_T^{miss} ATLAS-CONF-2013-144

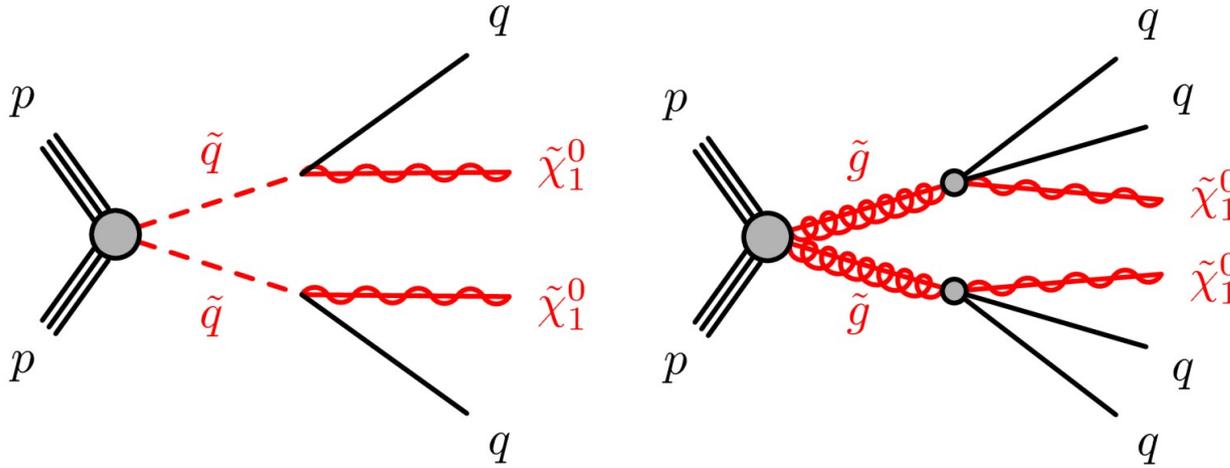
Final state depends on decay of squark/gluino



>> new result to be discussed

(since last ATLAS SUSY seminar on 26th March 2013)

Consider SUSY processes with decays to jets + LSP + no leptons



10 signal regions

≥ 2 to ≥ 6 jets + E_T^{miss}

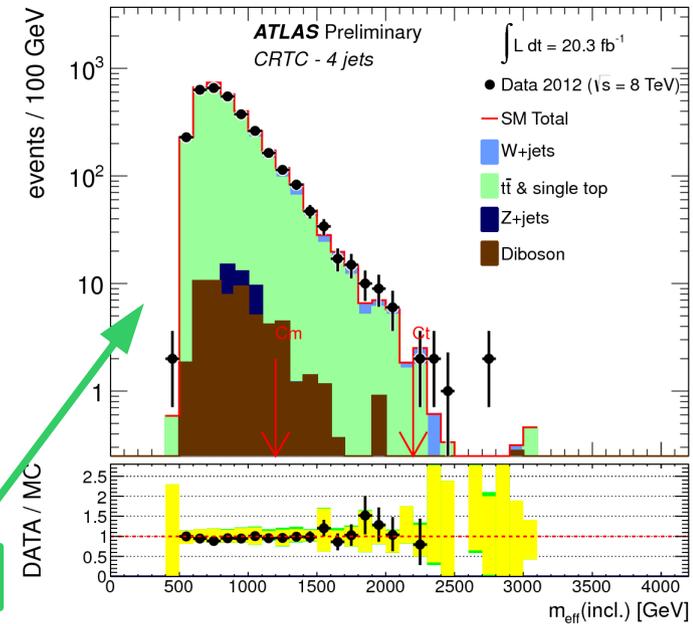
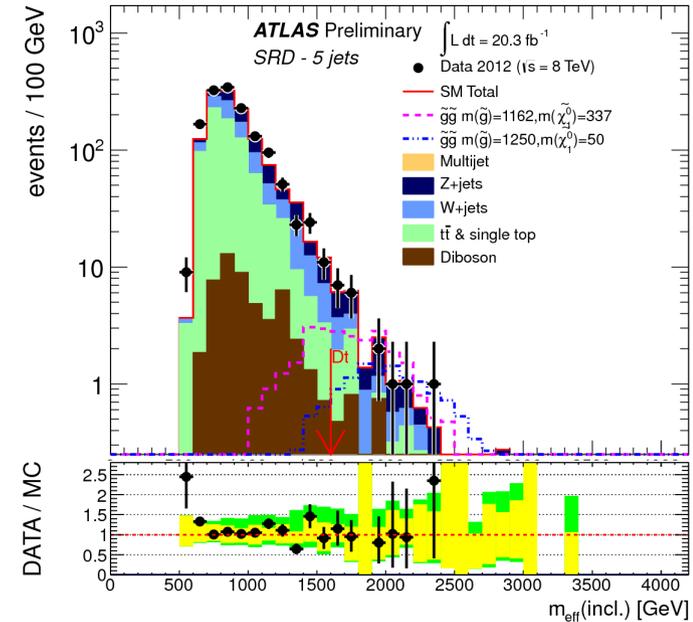
$m_{eff} > 1000$ to > 2200 GeV

$E_T^{miss} / m_{eff} > 0.15$ to 0.4

Jet + E_T^{miss} trigger

$$m_{eff} = E_T^{miss} + \sum p_T^{jets}$$

No significant excess seen



4 control regions per signal region

CR	SR background	CR process	CR selection
CRY	$Z(\rightarrow \nu\nu)+jets$	$\gamma+jets$	Isolated photon
CRQ	multi-jets	multi-jets	Reversed $\Delta\phi(jet, E_T^{miss})_{min}$ and $E_T^{miss}/m_{eff}(Nj)$ requirements ^a
CRW	$W(\rightarrow \ell\nu)+jets$	$W(\rightarrow \ell\nu)+jets$	$30 \text{ GeV} < m_T(\ell, E_T^{miss}) < 100 \text{ GeV}$, b -veto
CRT	$t\bar{t}$ and single- t	$t\bar{t} \rightarrow bbqq'\ell\nu$	$30 \text{ GeV} < m_T(\ell, E_T^{miss}) < 100 \text{ GeV}$, b -tag

Meaning of lines in interpretations

- - Expected limit
- Yellow band $\pm 1\sigma$ experimental uncertainties
- Red line: Observed limit
- - Dashed lines $\pm 1\sigma$ signal theory uncertainties

Signal uncertainties considered

In yellow band

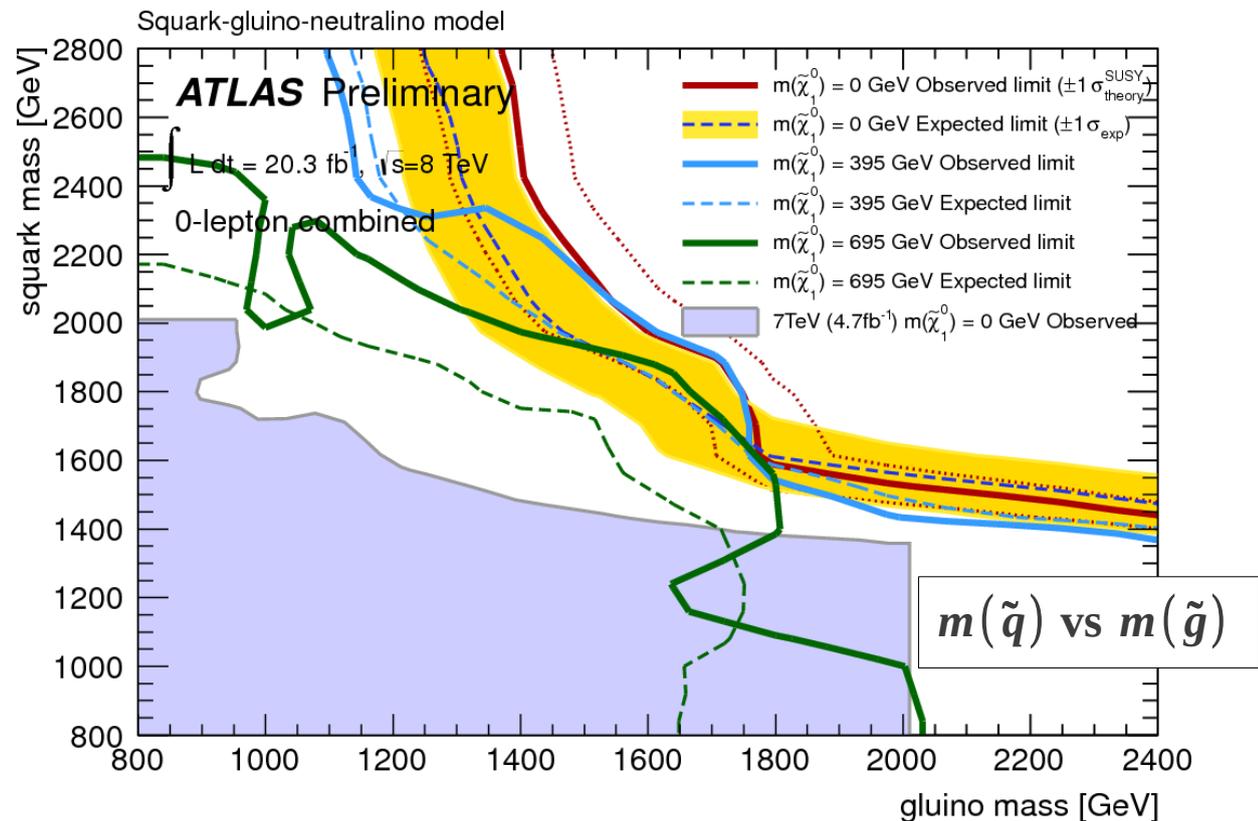
- Experimental uncertainties
- ISR uncertainty on signal MC
- Up to 30% in some regions with small Δm

In red dashed lines

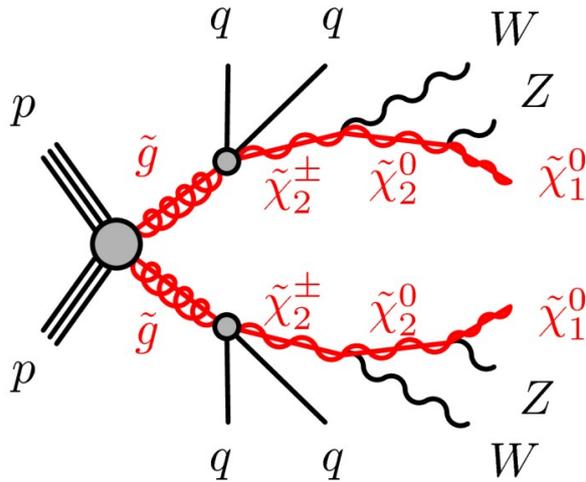
- Cross-section uncertainties (PDF, renormalisation/factorization scales)

\tilde{q} and \tilde{g} production

- $m(\tilde{\chi}_1^0) = 0$ GeV
- $m(\tilde{\chi}_1^0) = 395$ GeV
- $m(\tilde{\chi}_1^0) = 695$ GeV



Consider SUSY processes with longer decay chains \rightarrow more jets



Complimentary to previous analysis

Multijet trigger, possibility to look at lower E_T^{miss}

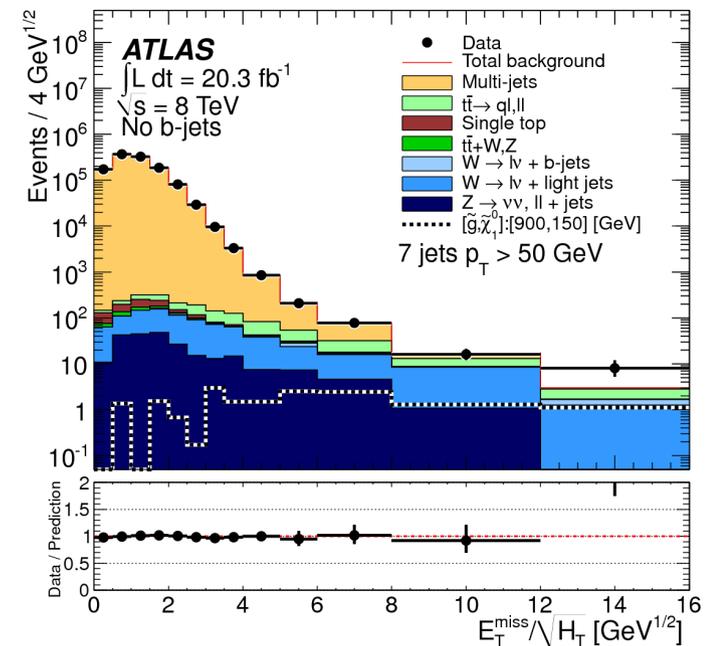
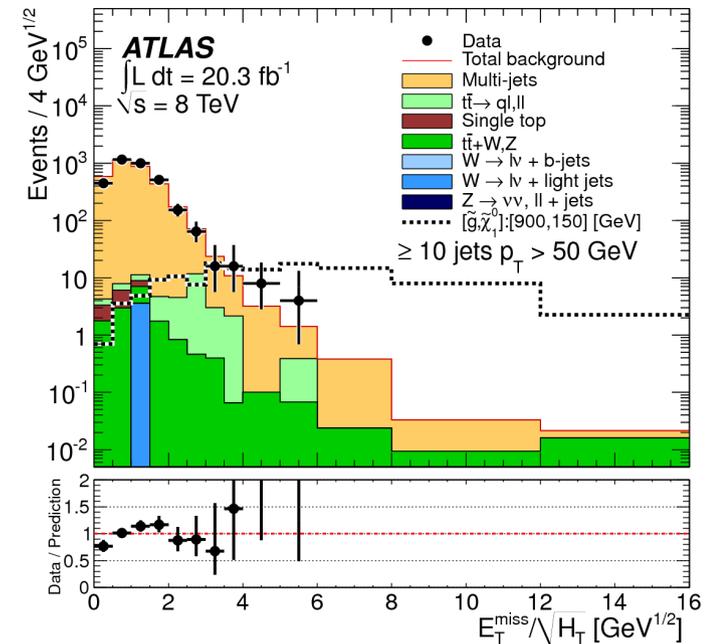
Signal regions $E_T^{miss} / \sqrt{H_T} : > 4 \text{ GeV}^{1/2}$
 High # jets + b-jets or High # jets + high mass composite jets - "fat jets"

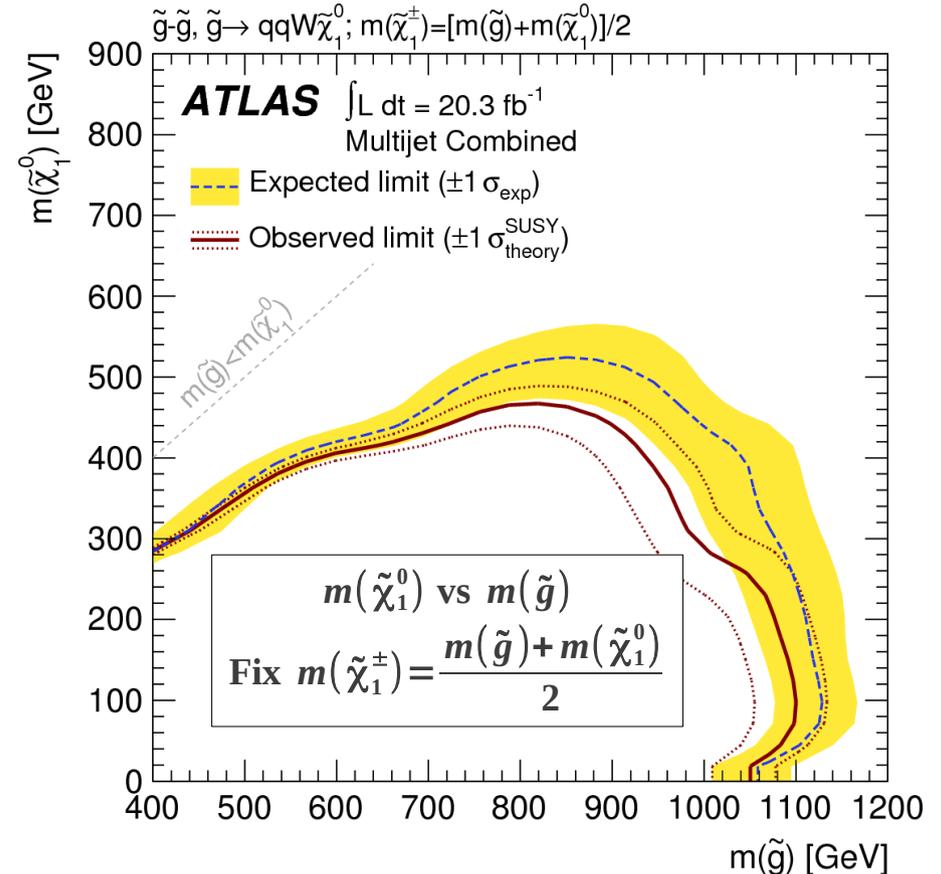
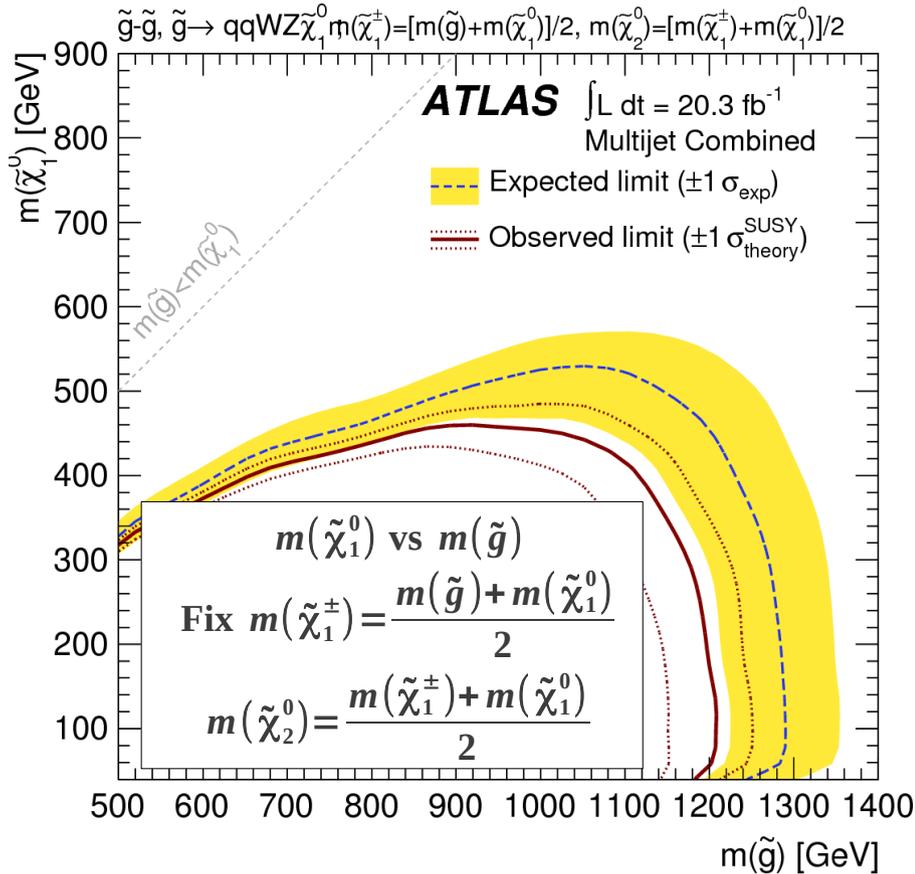
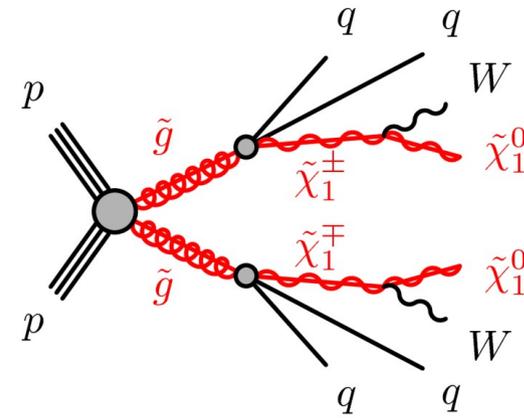
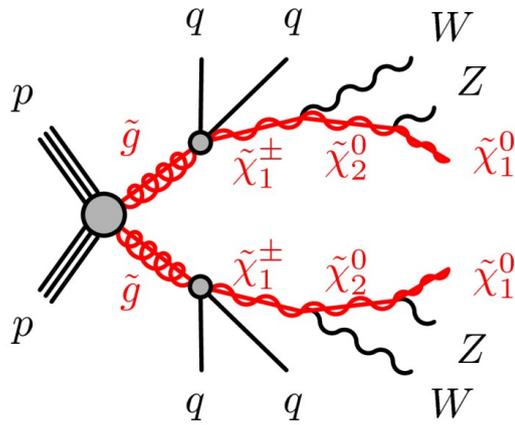
$$H_T = \sum p_T^{jets \ p_T > 40 \text{ GeV}}$$

Dominated by multijet background

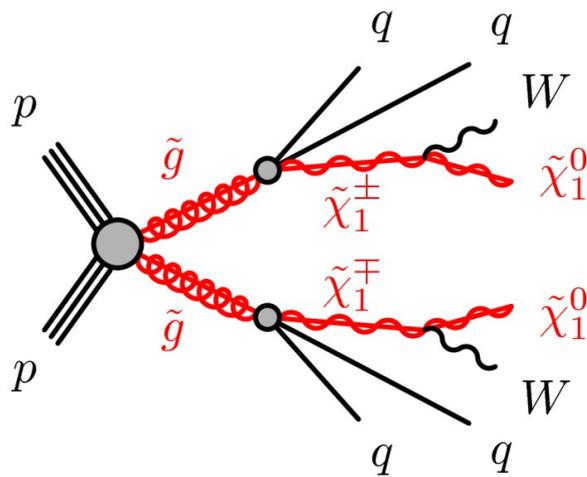
- Use $E_T^{miss} / \sqrt{H_T}$ shape in data from lower jet multiplicities
- Use same number of b-jets in CR and SR to get the template for every SR
- Only need to adjust out of cone energy for every Njets bin

No significant excess seen





Consider SUSY processes with leptons in the decay chain



Signal regions

High p_T jets and E_T^{miss}

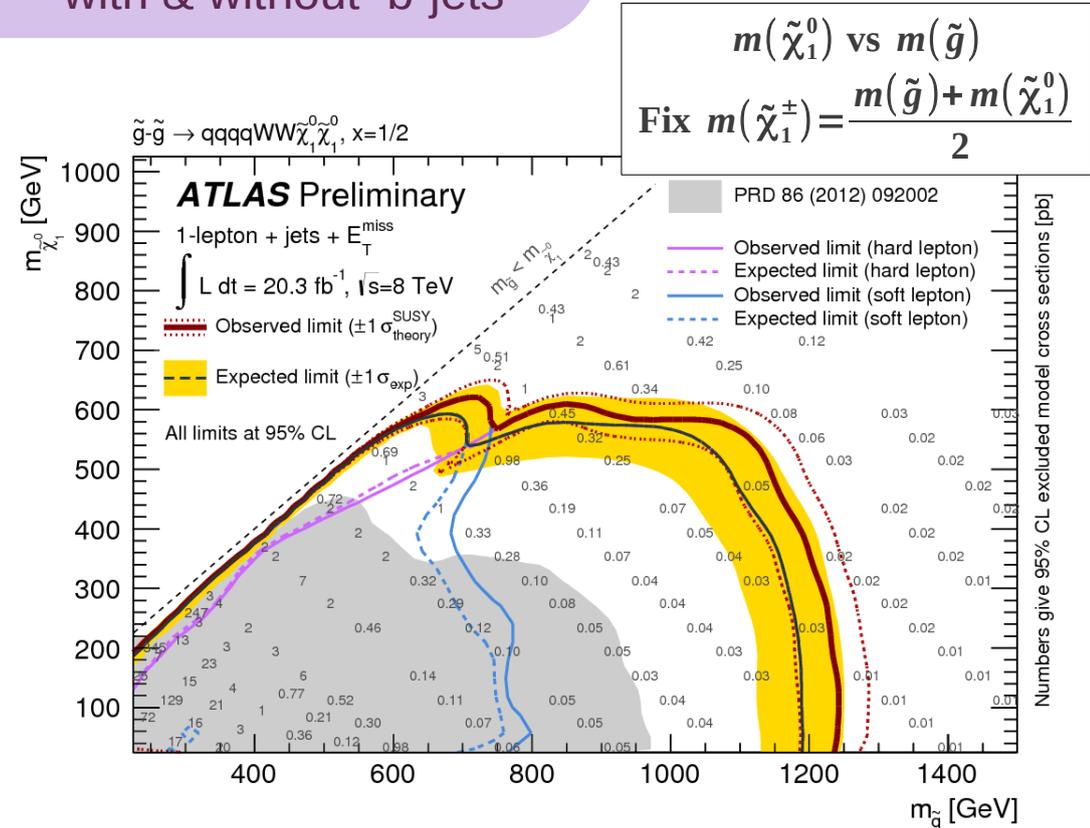
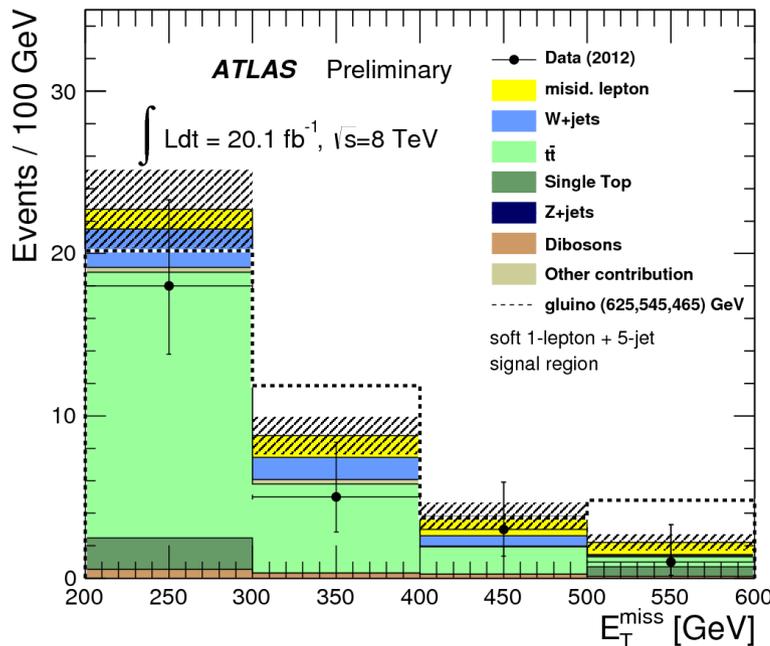
“Hard lepton” regions $p_T > 25$ GeV

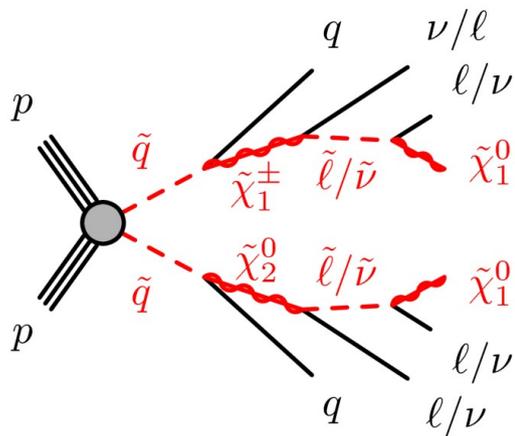
“Soft lepton” regions

$6(7) < p_T e(\mu) < 25$ GeV
with & without b-jets

Ttbar and W+jets dominated

No significant excess seen





Consider SUSY processes with decays to jets + LSP + 2 leptons
 Rather than use m_{eff} , exploit set of “Razor” variables
 Build mega-jets $j_1 j_2$ from visible decay products

$$M'_R = \sqrt{(j_{1,E} + j_{2,E})^2 - (j_{1,p_L} + j_{2,p_L})^2}$$

$$M_T^R = \sqrt{\frac{|\vec{E}_T^{\text{miss}}| (|\vec{j}_{1,p_T}| + |\vec{j}_{2,p_T}|) - \vec{E}_T^{\text{miss}} \cdot (\vec{j}_{1,p_T} + \vec{j}_{2,p_T})^2}{2}}$$

$$R = \frac{M_T^R}{M'_R}$$

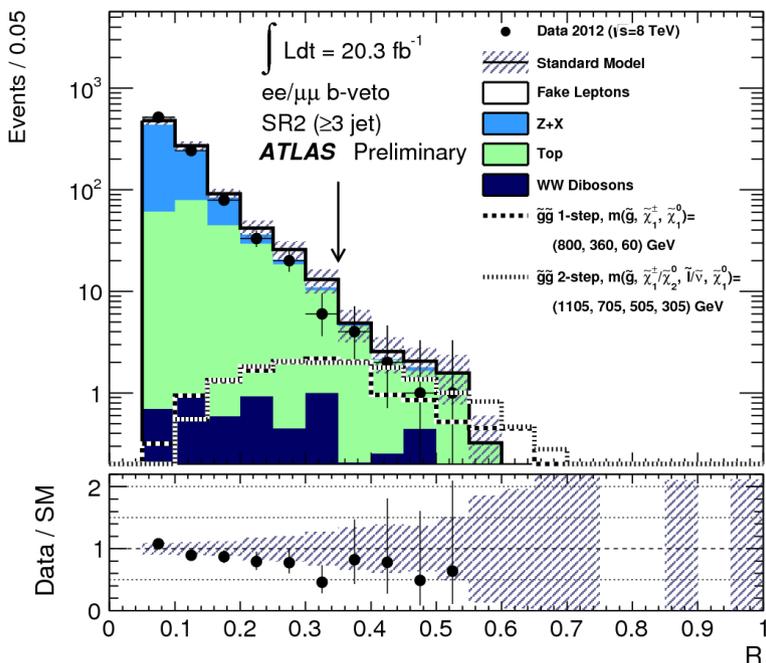
Signal regions

Z veto

< 3 or ≥ 3 jets

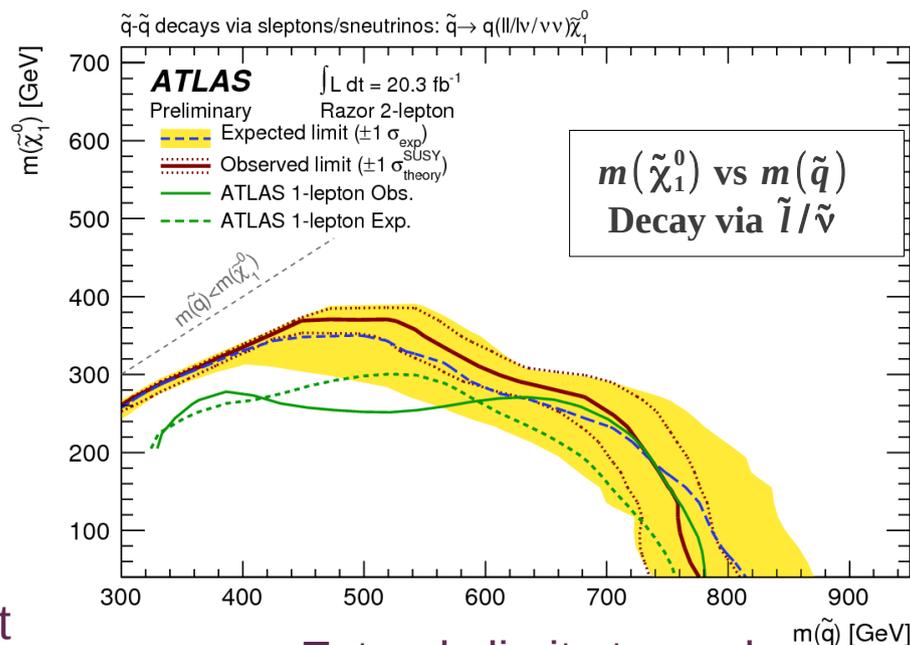
Large R

Large M'_R



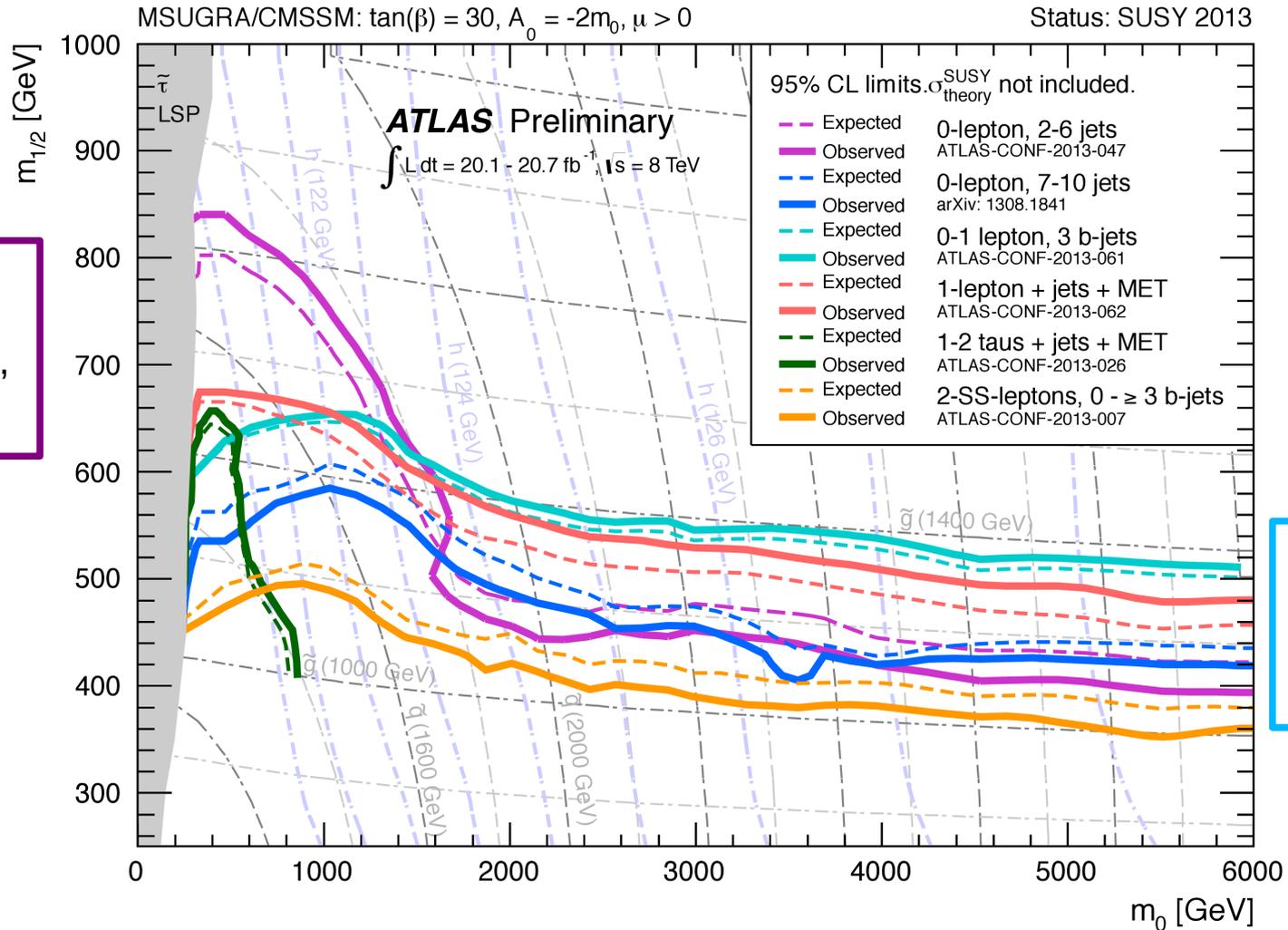
Ttbar
and Z
dominated

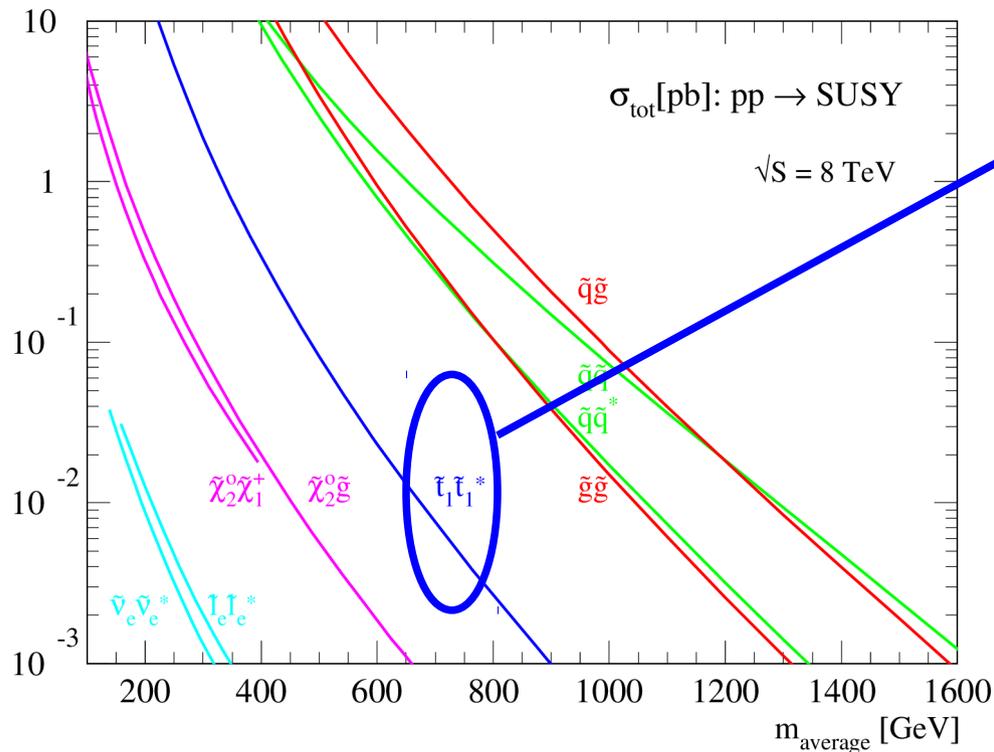
No significant
excess seen



Extends limits towards
smaller mass splittings

A specific SUSY framework like mSUGRA can be used to compare the performance of search channels





Smaller cross-section than
1st/2nd gen. squarks

Light stop preferred by
naturalness

Final state depends on decay
of stop/sbottom

3rd gen. search channels

Glauino mediated production

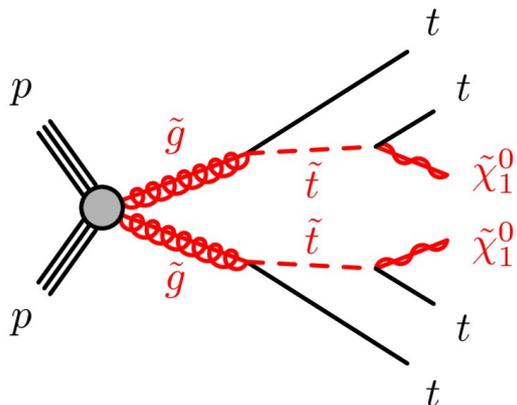
- » 0-1 leptons + ≥ 3 b-jets + E_T^{miss}
ATLAS-CONF-2013-024
- 2 SS leptons (+ b-jets) + E_T^{miss}
ATLAS-CONF-2013-007
- 3 leptons + jets + E_T^{miss}
ATLAS-CONF-2012-151

>> new result to be discussed
(since last ATLAS SUSY seminar on 26th March 2013)

Direct production

- 0-leptons + 6-jets + 2 b-jets + E_T^{miss}
ATLAS-CONF-2013-024
- » 0-leptons + 2 b-jets + E_T^{miss}
1-leptons + 4-jets (1 b-jets) + E_T^{miss}
ATLAS-CONF-2013-037
- » 2-leptons (+ 2 b-jets) + E_T^{miss}
- » Charm/mono-jet + E_T^{miss}
Z(ll) + 2 b-jets + E_T^{miss} ATLAS-CONF-2012-025

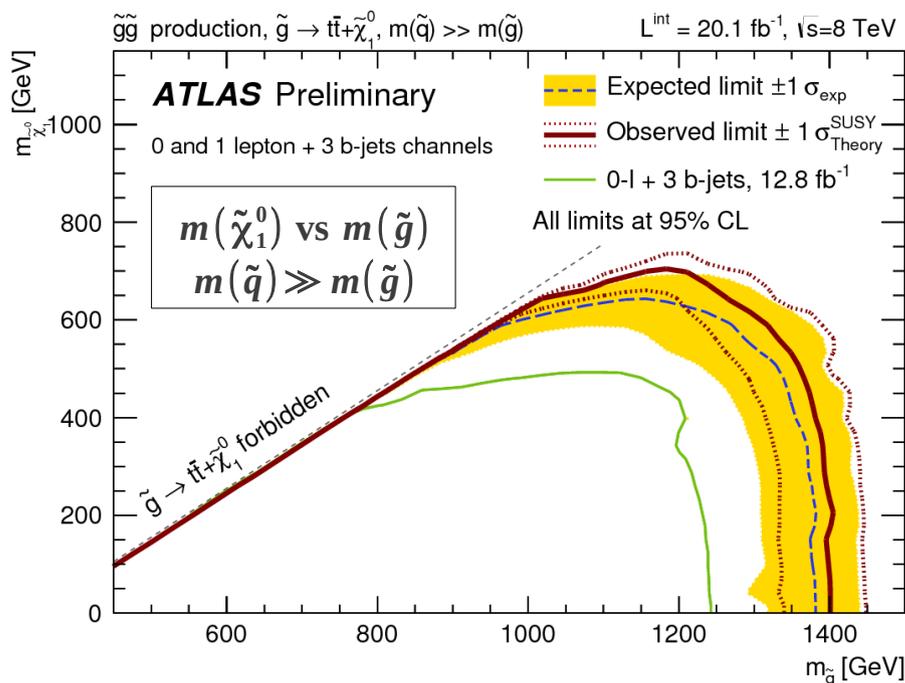
Consider SUSY processes with decays to many (b)jets
+ LSP + 0-1 leptons



9 signal regions

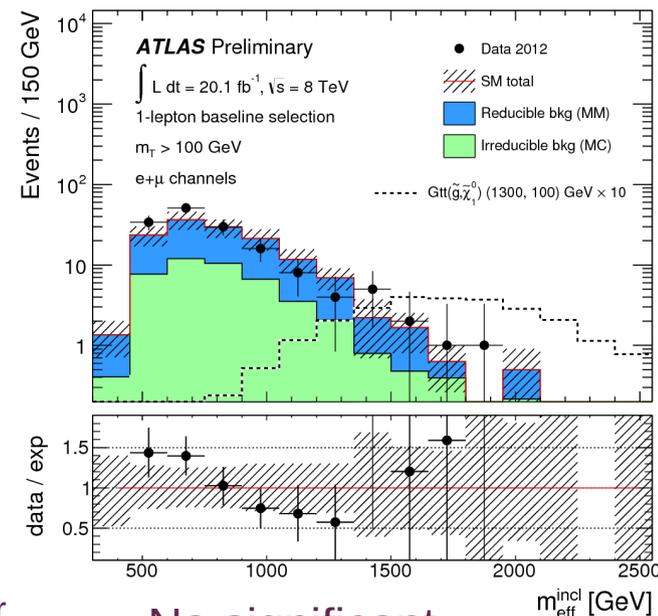
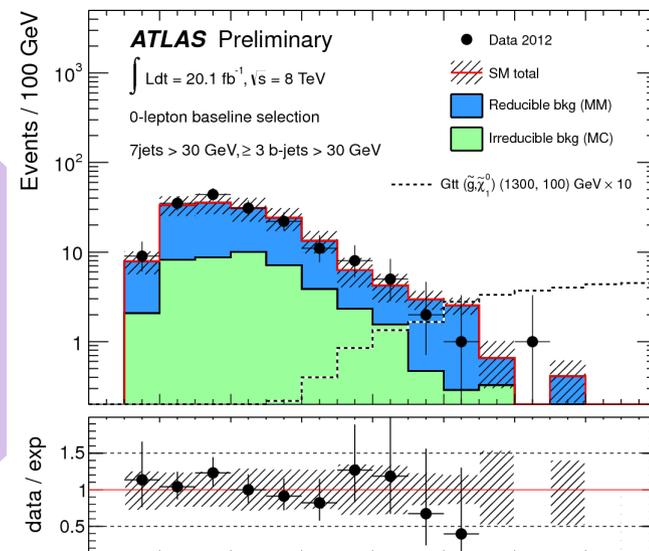
- 0 or 1 lepton
- ≥ 3 b-tagged jets
- ≥ 4 to ≥ 7 jets
- High E_T^{miss} , m_{eff} , $(m_T, E_T^{miss} / \sqrt{H_T})$

Reducible $t\bar{t}$ background (≥ 1 mis-tagged b-jet)
Irreducible $t\bar{t}$ +b/bb/Z/h

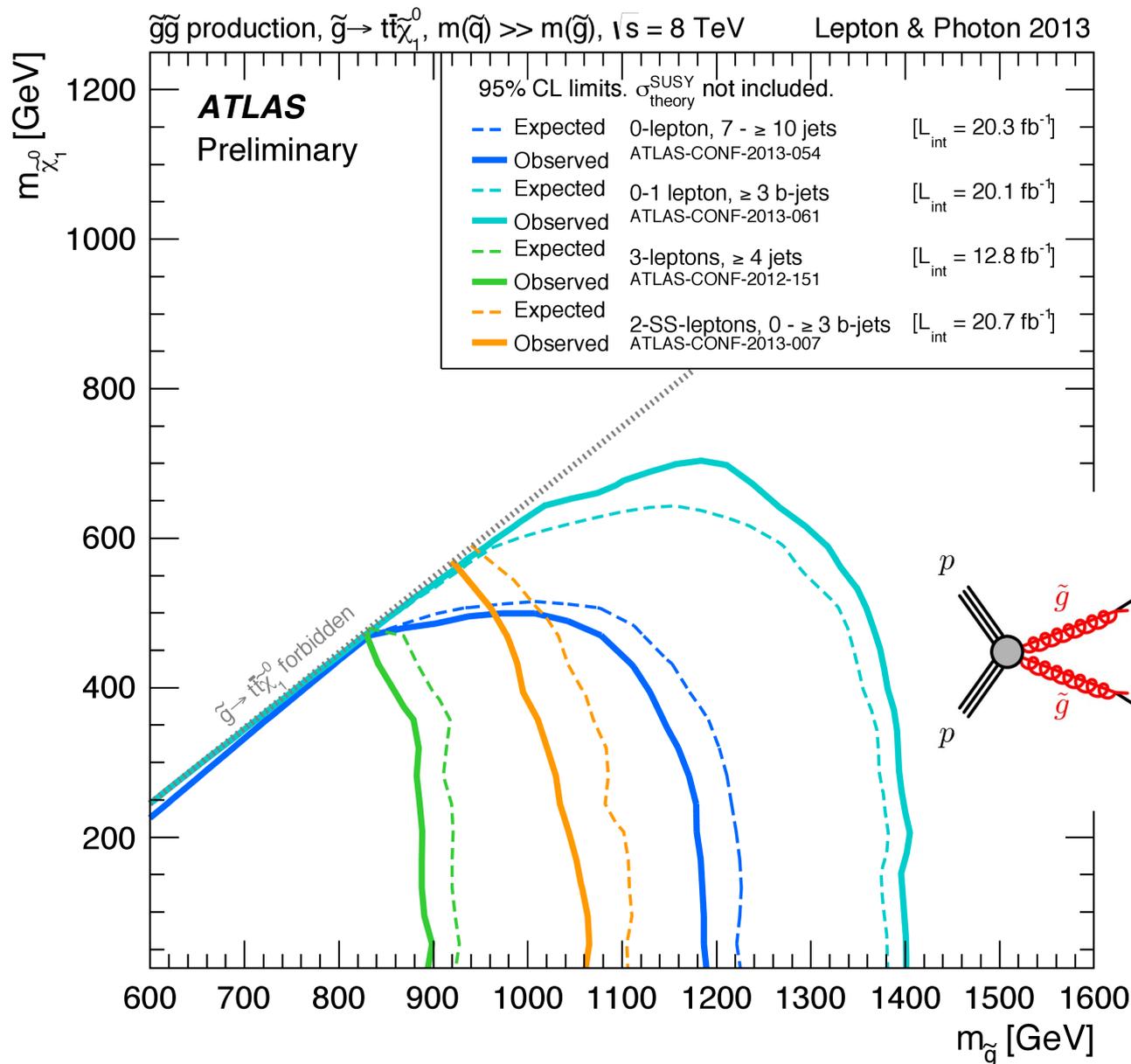


Interpretation in
gluino-mediated 3rd
gen. production

+ more interpretations for
decays via heavy
neutralinos



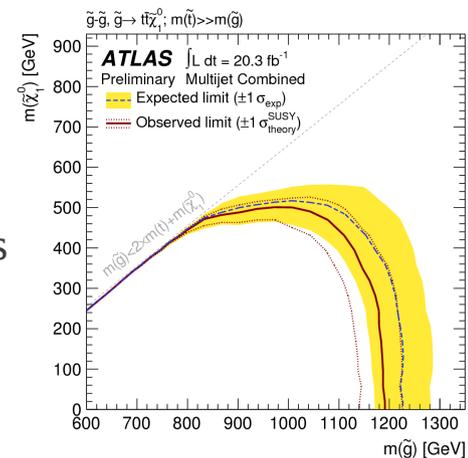
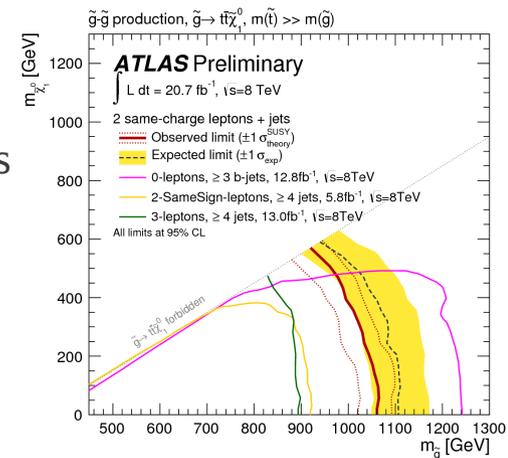
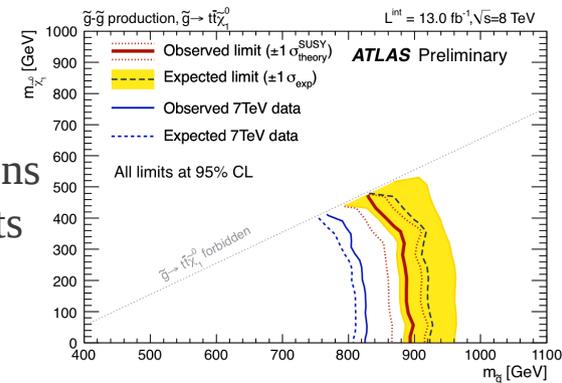
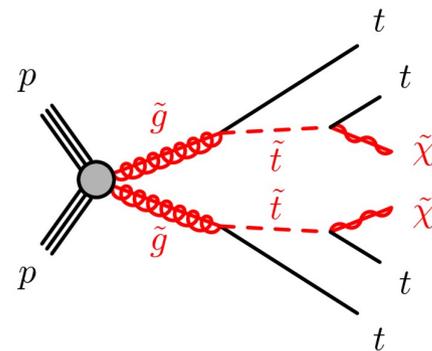
No significant
excess seen



3-leptons
 ≥ 4 jets

SS-leptons
 $\geq 0 - 3$ b-jets

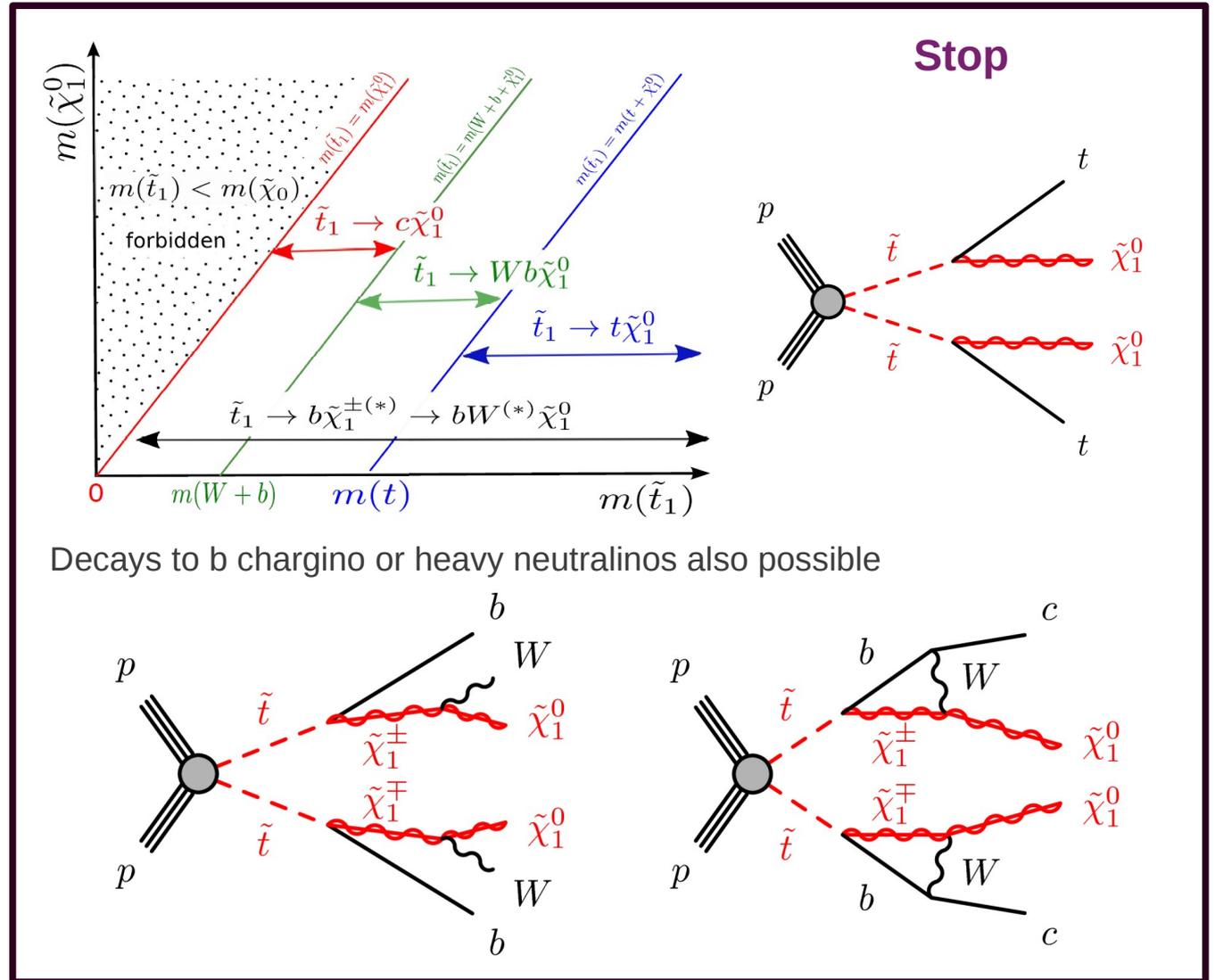
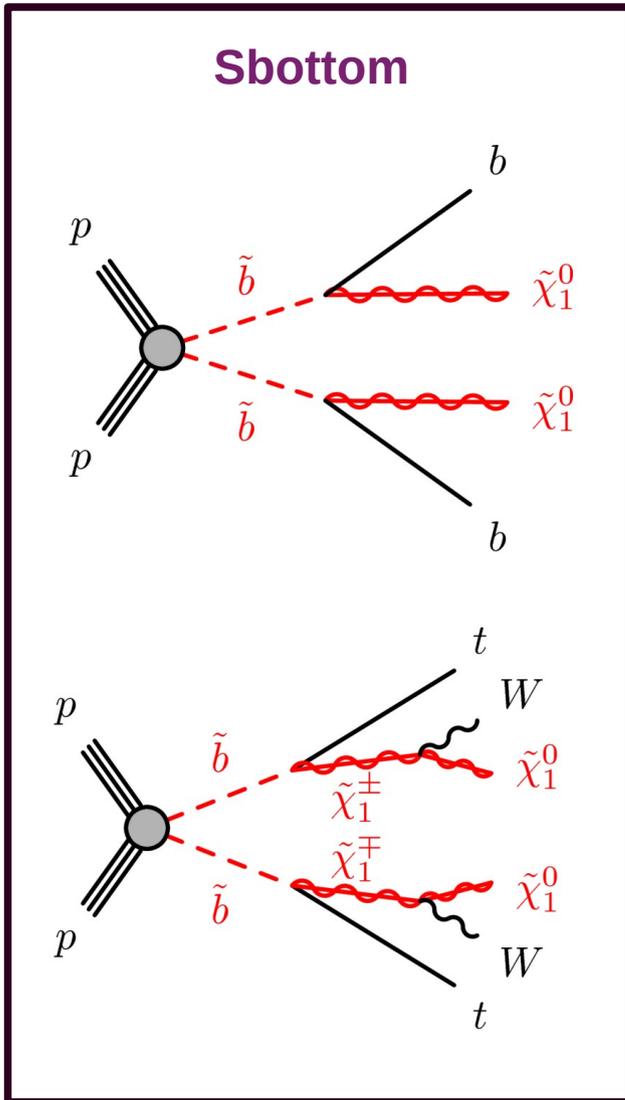
0-lepton
 $\geq 7 - 10$ jets



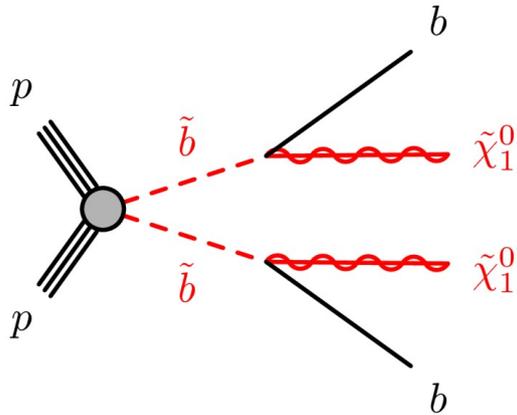
3rd gen. squark searches

Stop/sbottom quark searches target many different scenarios

Sensitivity is dependent on sparticle mass differences and decay channels



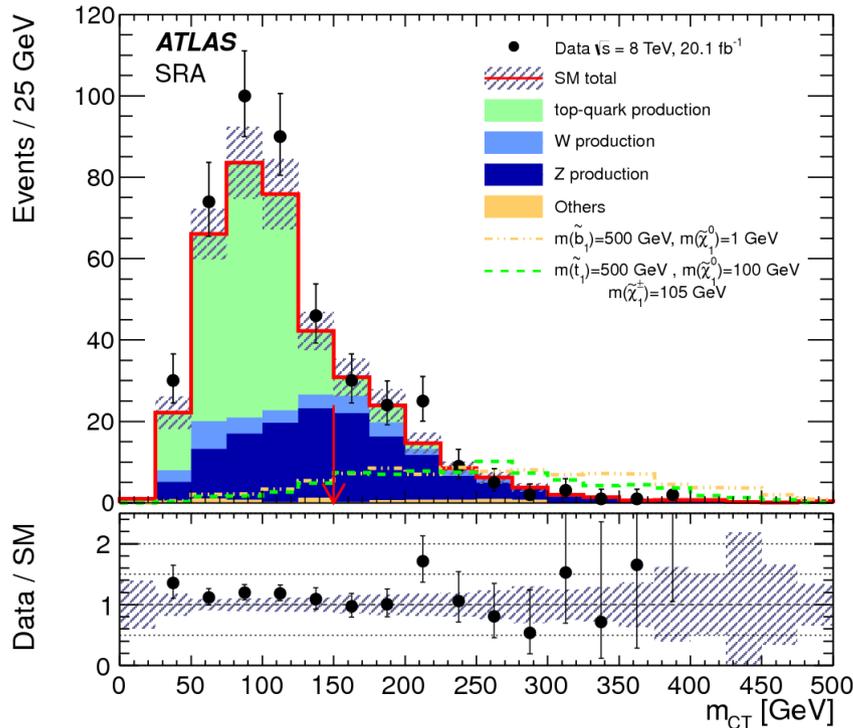
Consider SUSY processes with decays to b-jets + LSP + no leptons



2 signal regions

Large $\Delta m(\tilde{b}, \tilde{\chi}_1^0)$
 = 2 b-tagged jets
 = 2 high p_T jets
 high m_{CT}, m_{bb}

Small $\Delta m(\tilde{b}, \tilde{\chi}_1^0)$
 = 2 b-tagged jets
 additional ISR jet
 + low additional hadronic energy



“Contransverse mass”

$$m_{CT}^2 = (E_T^{b1} + E_T^{b2})^2 - |\mathbf{p}_T^{b1} - \mathbf{p}_T^{b2}|^2$$

For $\tilde{b}\tilde{b}$ events, end-point defined by:

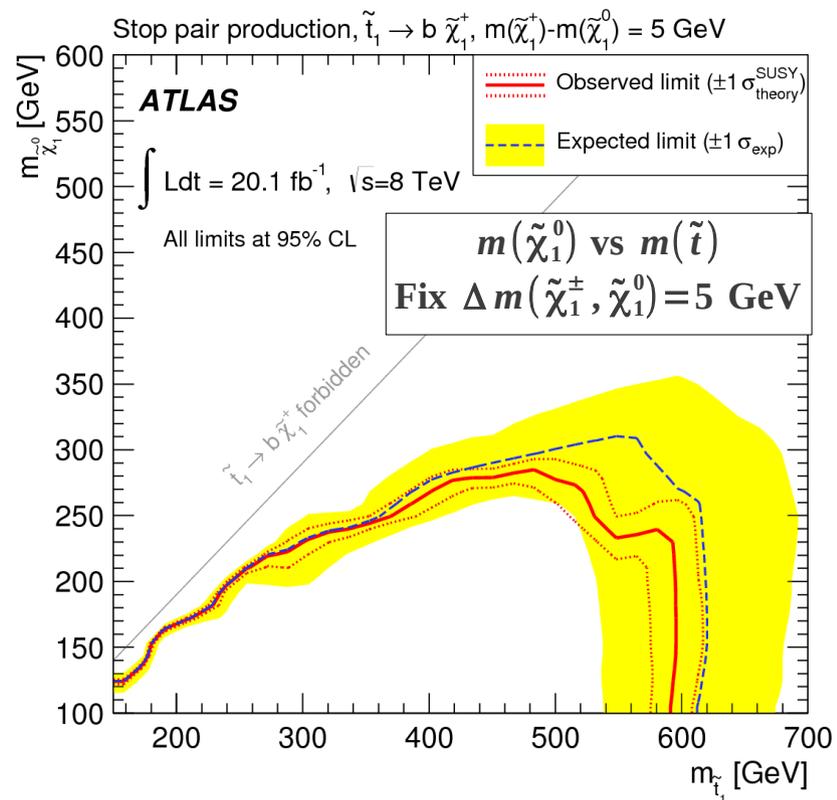
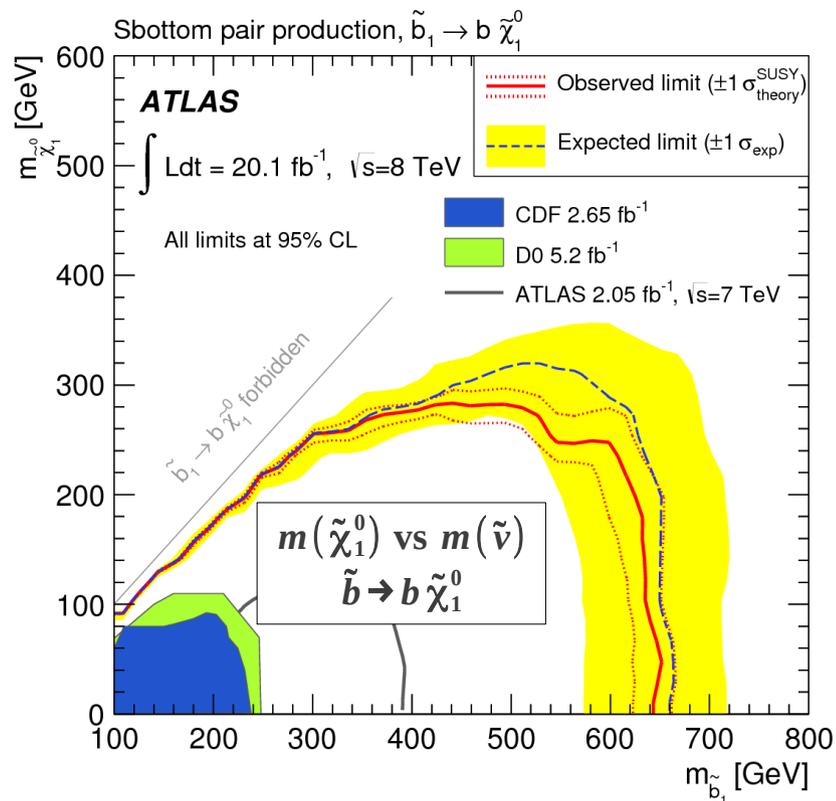
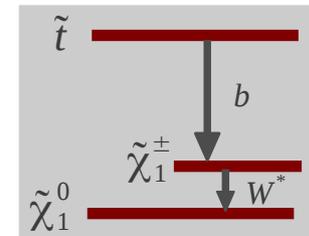
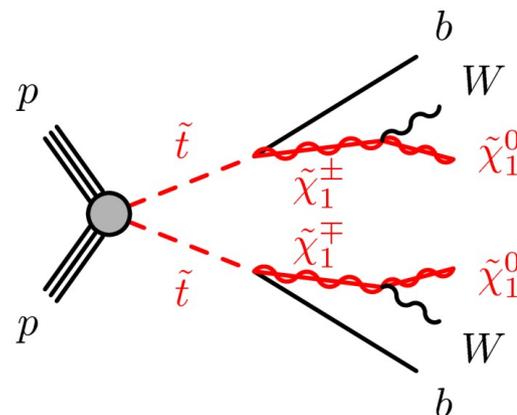
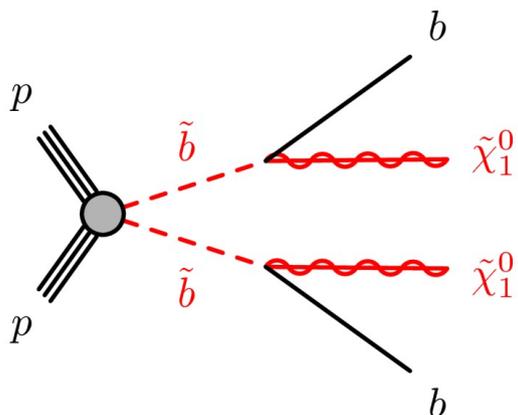
$$\frac{m(\tilde{b})^2 - m(\tilde{\chi}_1^0)^2}{m(\tilde{b})}$$

For $t\bar{t}$ events, end-point ~ 135 GeV

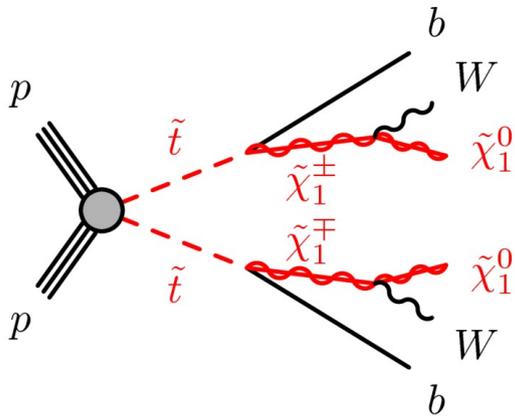
Ttbar, W+bjets, Z+bjets dominated

No significant excess seen

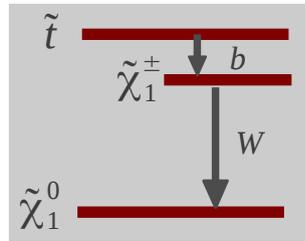
Analysis is reinterpreted for stop production



Consider SUSY processes with decays to two leptons + jets + LSP

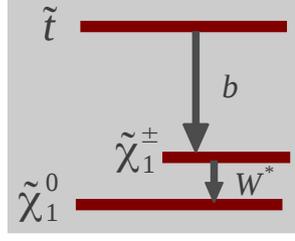


Large $\Delta m(\tilde{\chi}_1^\pm, \tilde{\chi}_1^0)$



4 signal regions
 $m_{T2}^{ll} > 90$ to 120 GeV
 ≥ 2 jets
 one region without jets for sensitivity to small $\Delta m(\tilde{t}, \tilde{\chi}_1^\pm)$

Large $\Delta m(\tilde{t}, \tilde{\chi}_1^\pm)$



1 signal region
 Low p_T leptons
 $= 2$ b-tagged jets
 $m_{T2}^{b\text{-jet}} > 160$ GeV

m_{T2} variable "Stransverse mass"
 Kinematic endpoint
 \sim mass of semi-invisibly decaying particle

$$m_{T2}^{ll}(l_1, l_2, E_T^{\text{miss}})$$

Bounded by W mass for WW, Wt, ttbar

$$m_{T2}^{b\text{-jet}}(b_1, b_2, l_1 + l_2 + E_T^{\text{miss}})$$

Bounded by top mass for ttbar

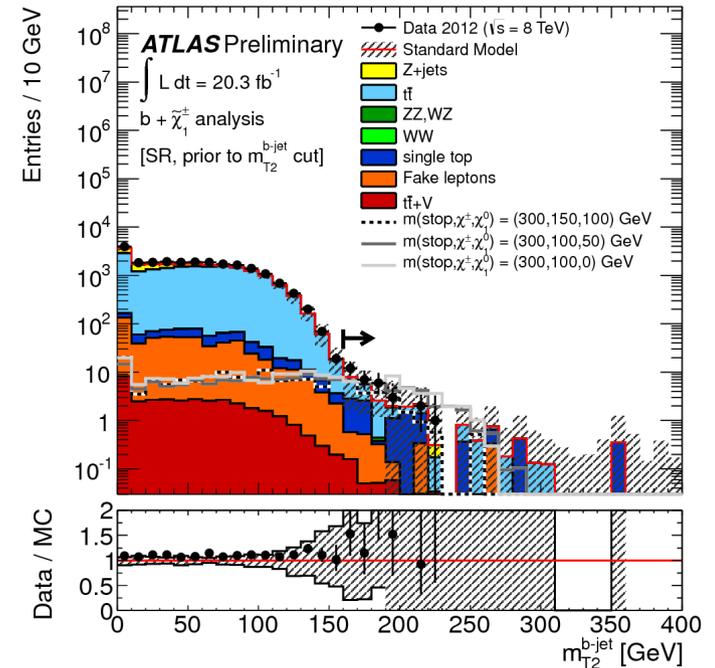
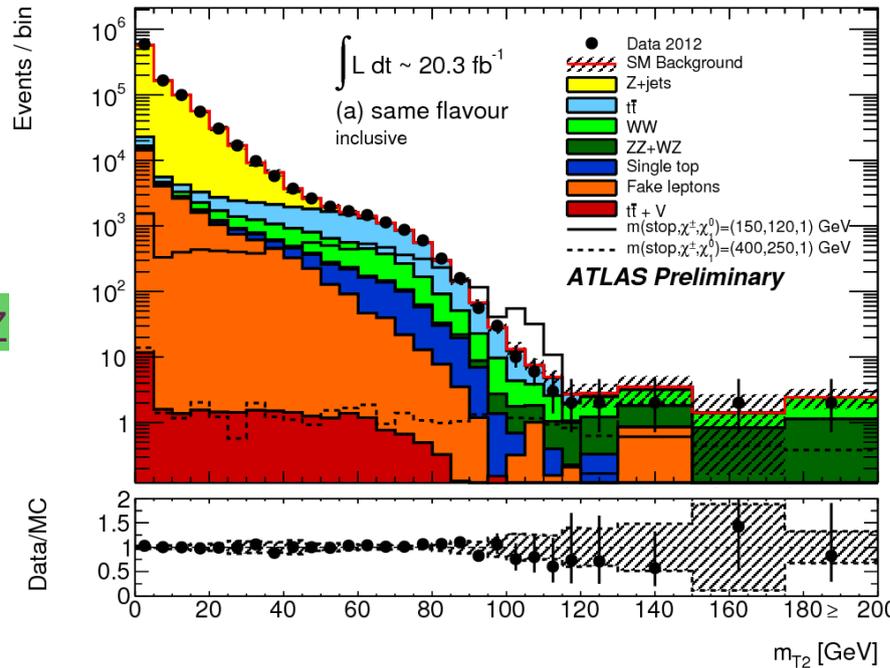
Backgrounds

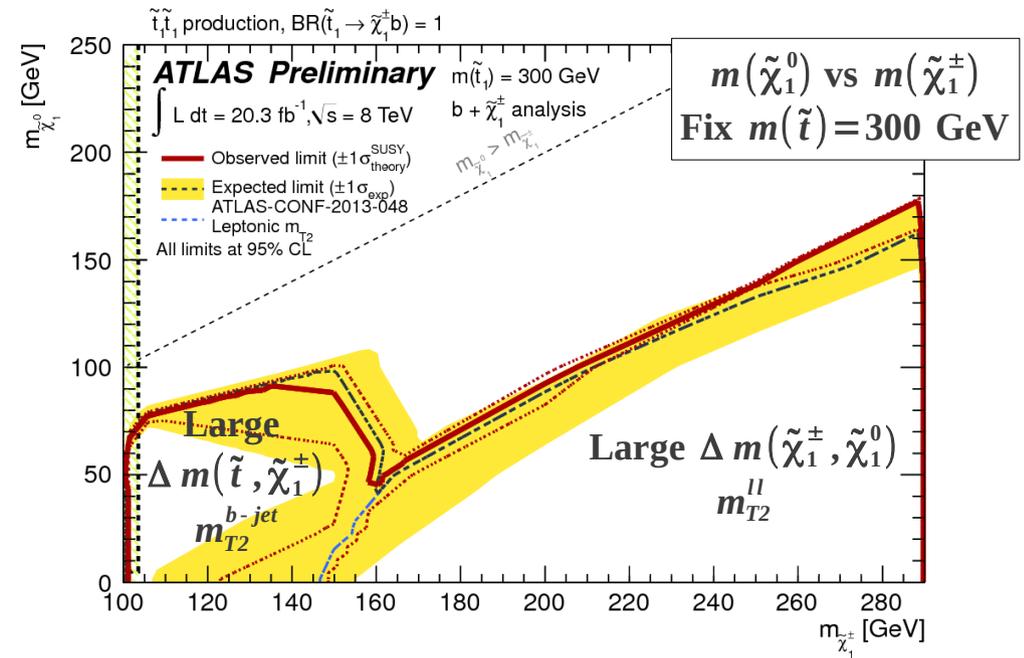
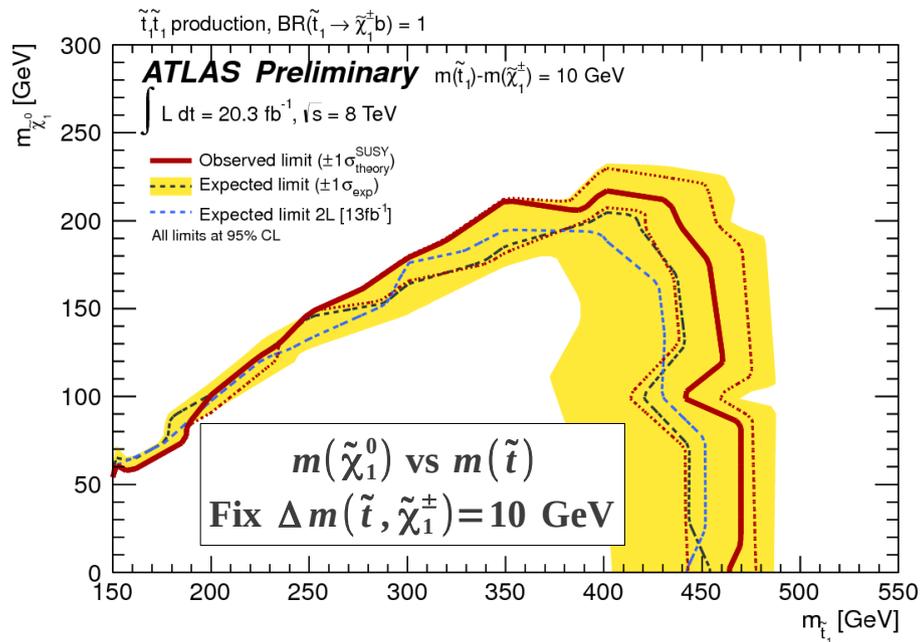
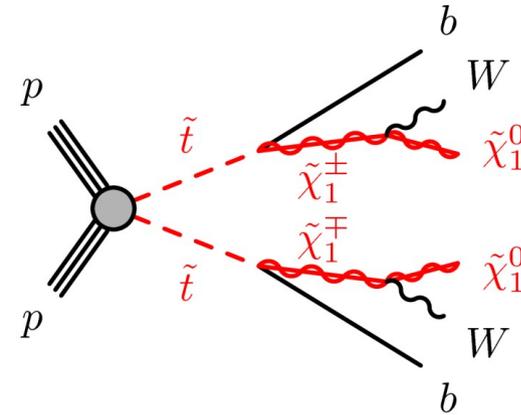
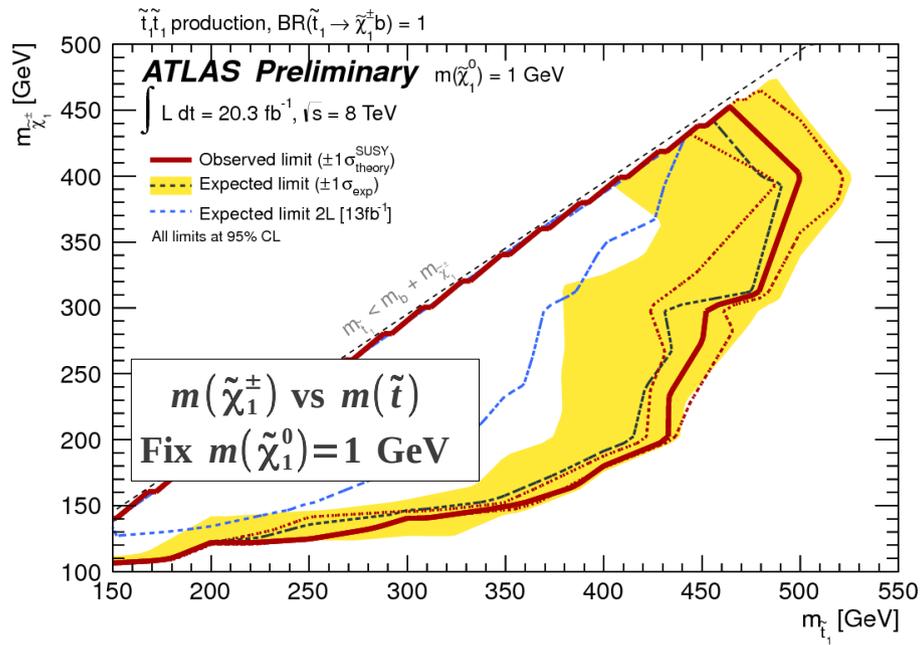
Large $\Delta m(\tilde{\chi}_1^\pm, \tilde{\chi}_1^0)$

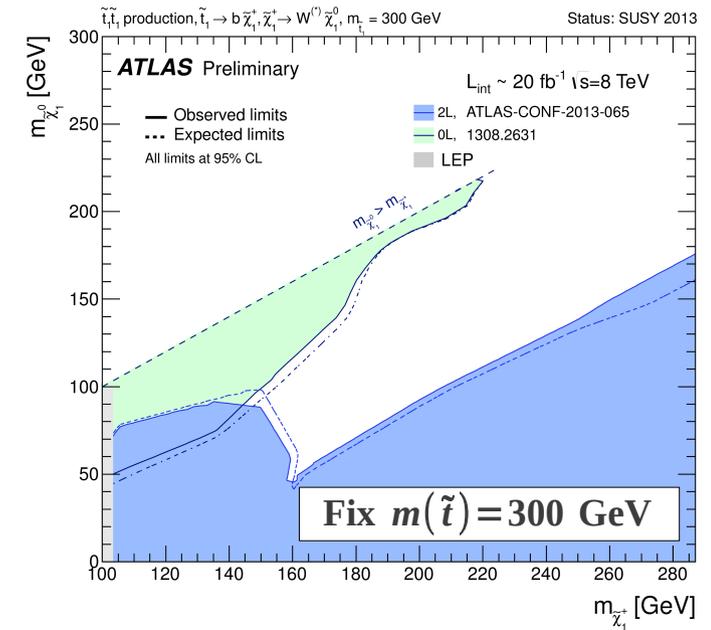
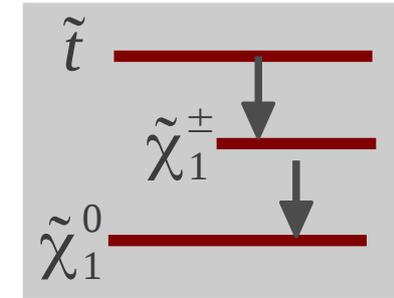
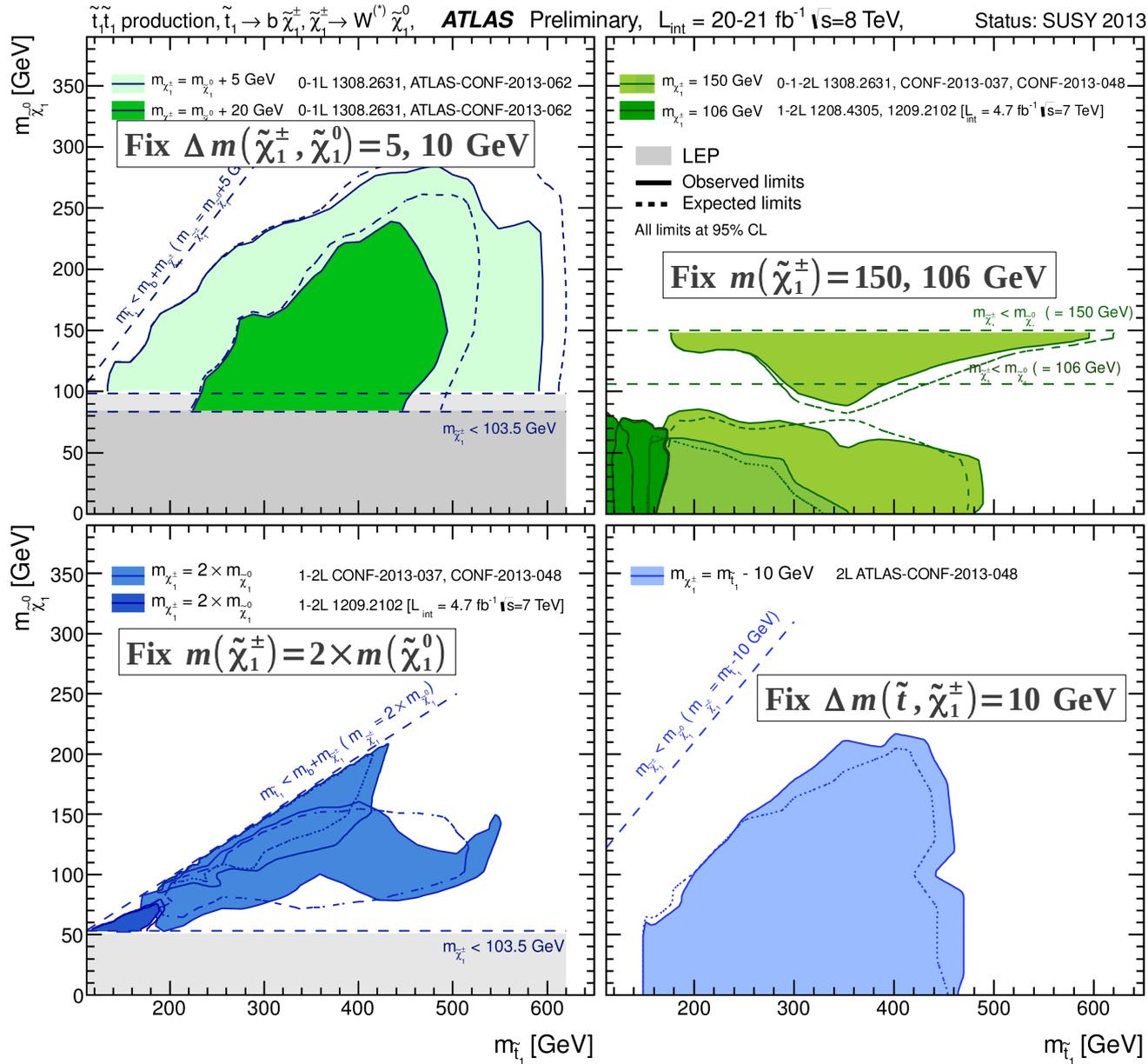
Ttbar, WW, WZ, ZZ

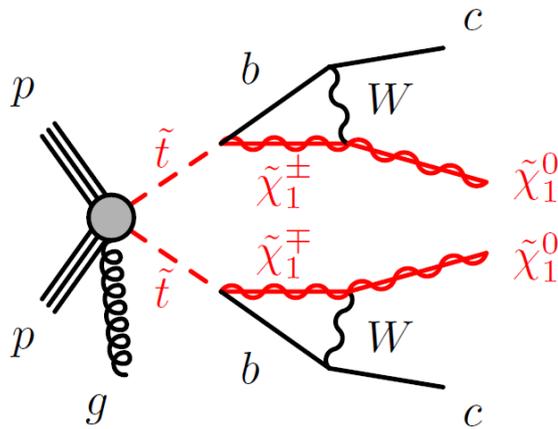
Large $\Delta m(\tilde{t}, \tilde{\chi}_1^\pm)$

Ttbar, Wt, Z+jets









For $\Delta m(\tilde{t}, \tilde{\chi}_1^0) < m_W$ and $m(\tilde{t}) < m(\tilde{\chi}_1^\pm)$

stop can decay via loop to $\tilde{t} \rightarrow c \tilde{\chi}_1^0$

Soft c-jets and small E_T^{miss}

→ ISR/FSR jet for trigger and signal/background separation

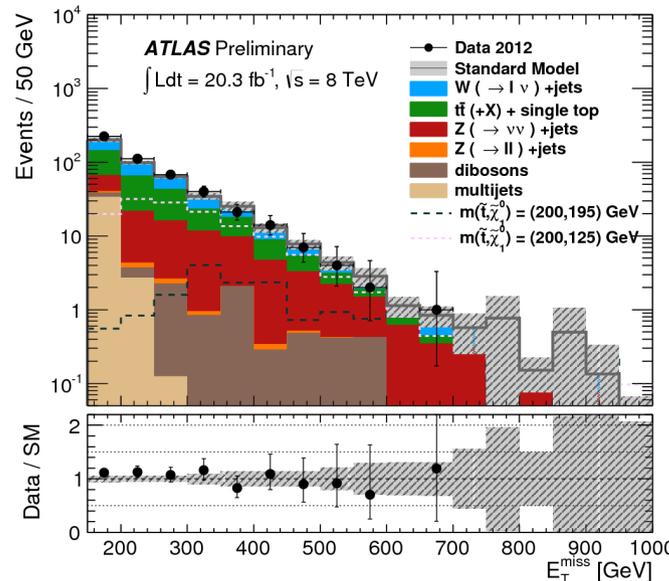
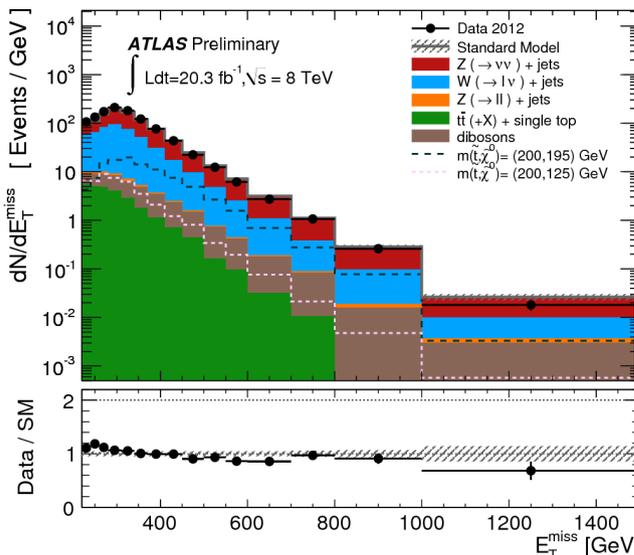
2 signal regions: High E_T^{miss} , High p_T ISR/FSR jet

Small $\Delta m(\tilde{t}, \tilde{\chi}_1^0)$

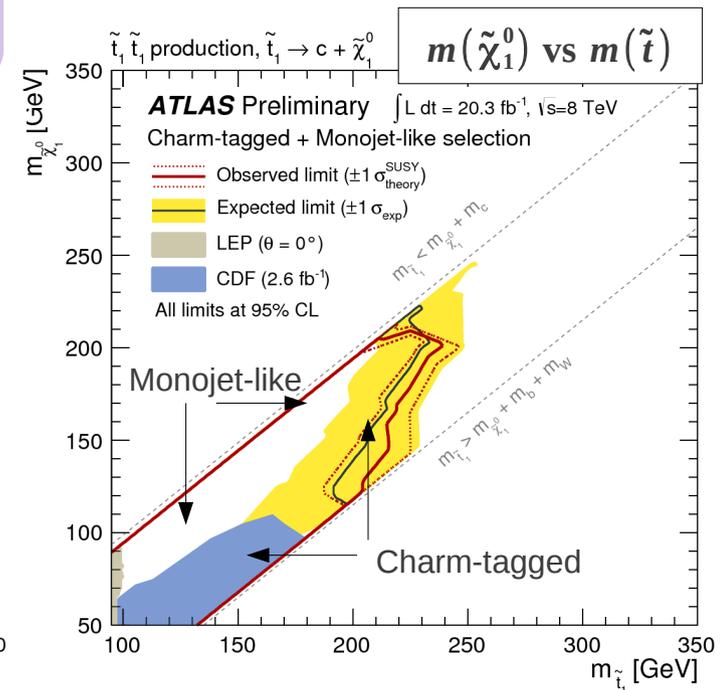
≥ 1-3 jets, no c-tagging
“Monojet-like”

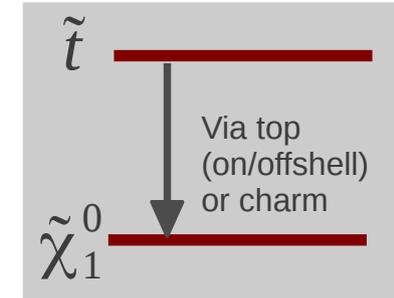
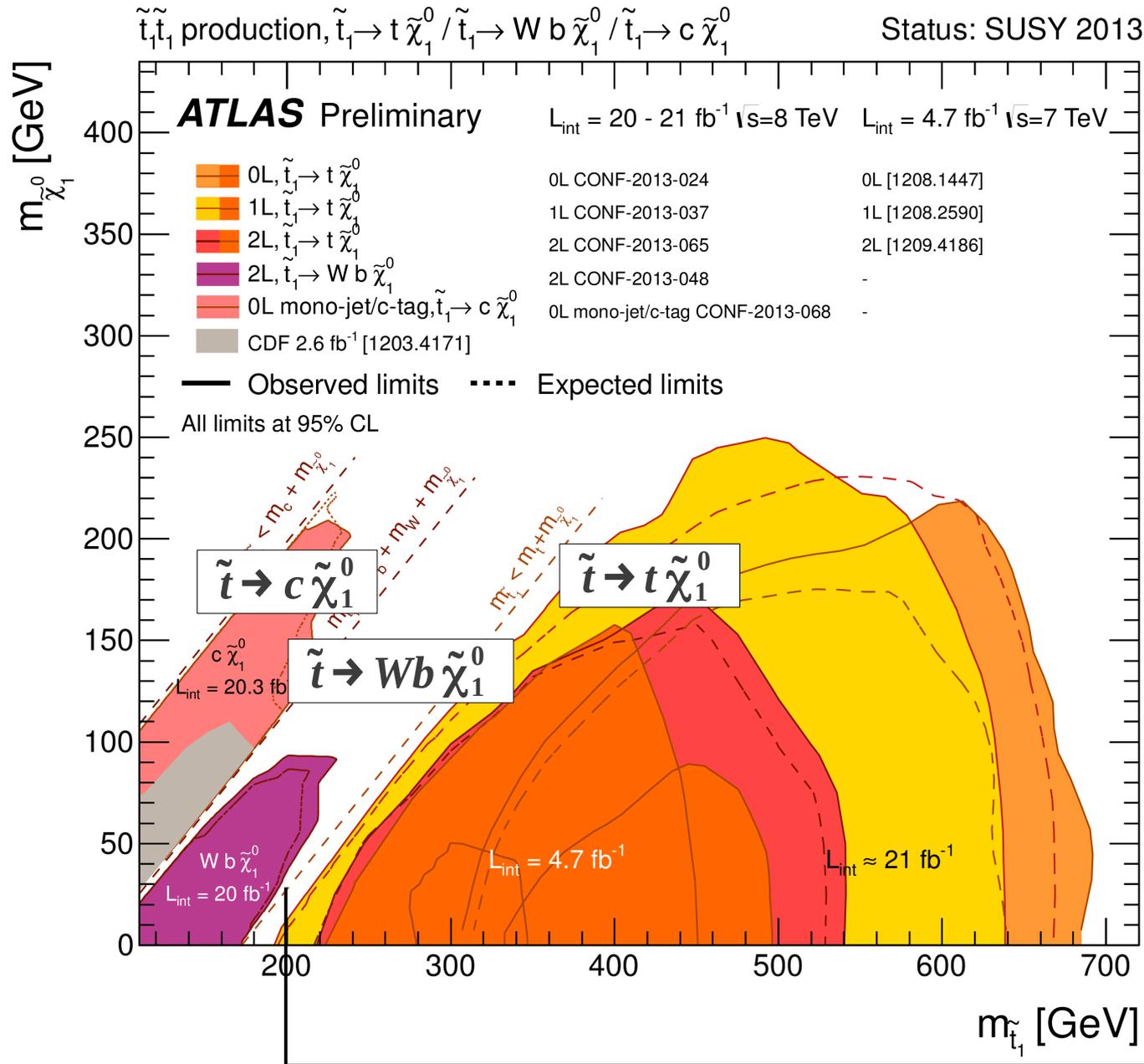
Moderate $\Delta m(\tilde{t}, \tilde{\chi}_1^0) \sim 20-80$ GeV

≥ 4 jets, some c-tagged
“Charm-tagged”



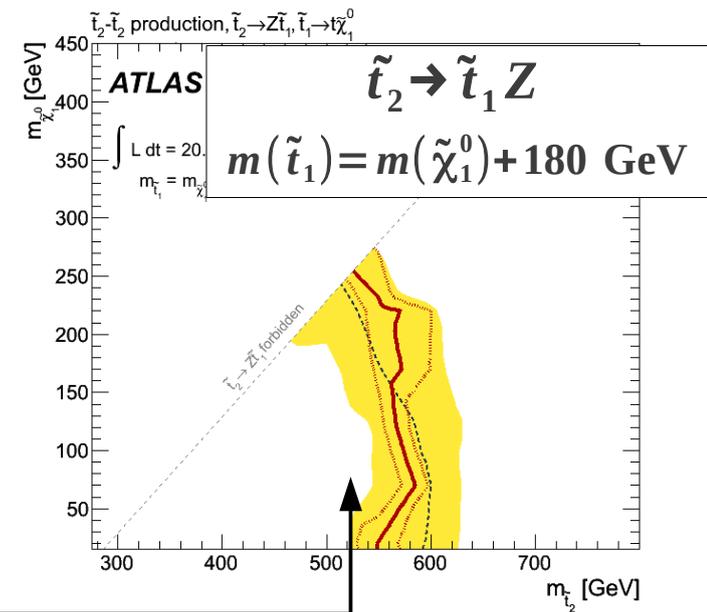
Ttbar, W+bjets, Z+bjets
dominated

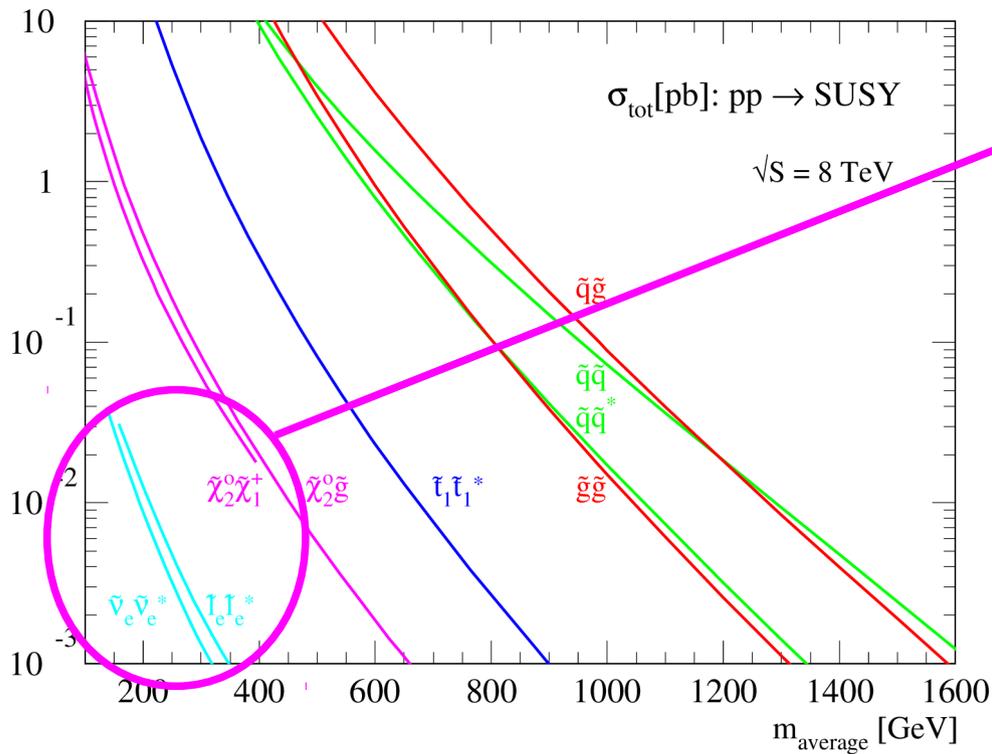




Very good coverage in stop searches, using many different channels and techniques.

Some challenging gaps remain!





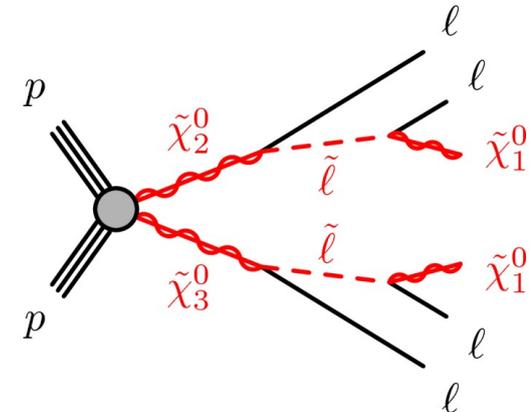
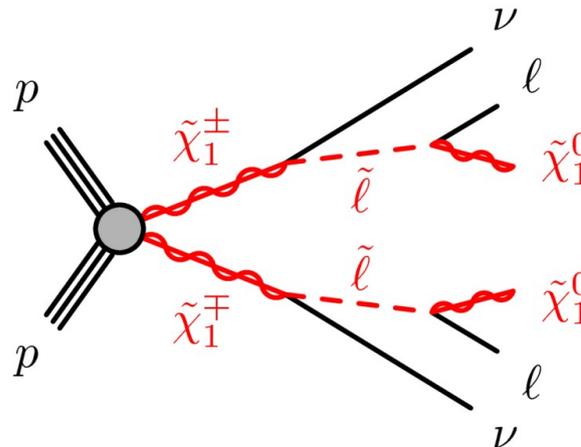
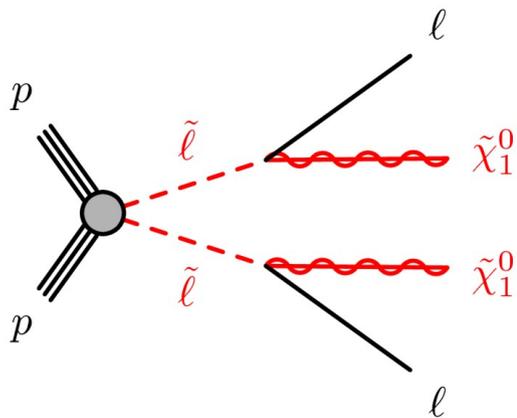
Electroweak SUSY processes have small cross-sections
 Light higgsinos preferred by naturalness

EWKino search channels

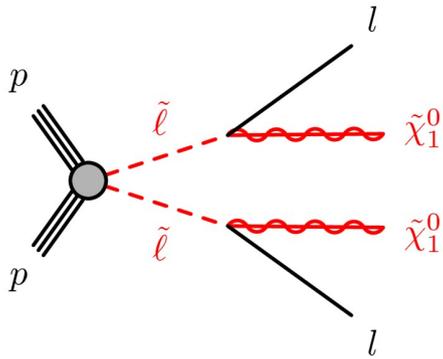
- » 2-leptons + $E_{\text{T}}^{\text{miss}}$
- 3-leptons + $E_{\text{T}}^{\text{miss}}$ ATLAS-CONF-2013-035
- 4-leptons + $E_{\text{T}}^{\text{miss}}$ ATLAS-CONF-2013-036
- 2-taus + $E_{\text{T}}^{\text{miss}}$ ATLAS-CONF-2013-028
- » 1-lepton + 2 b-jets + $E_{\text{T}}^{\text{miss}}$

>> new result to be discussed

(since last ATLAS SUSY seminar on 26th March 2013)



Consider SUSY processes with decays to 2 leptons + LSP



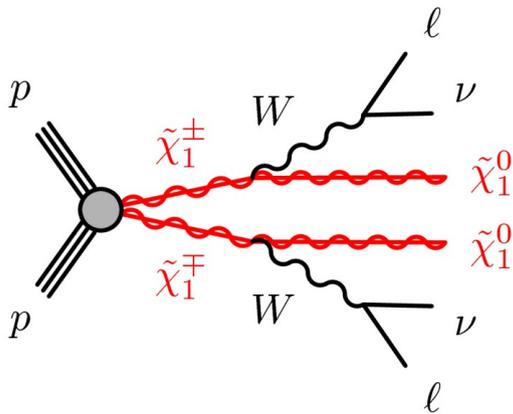
2 signal regions

Target slepton decays
(direct or via chargino decays)

SFOS or DFOS

Z veto, jet veto

Large $E_T^{miss,rel}$ and m_{T2}



3 signal regions

Target charginos decaying via W
Difficult due to low $BR(W \rightarrow l\nu)$

DFOS only

$m(l\bar{l})$ and $p_T(l\bar{l})$ selections

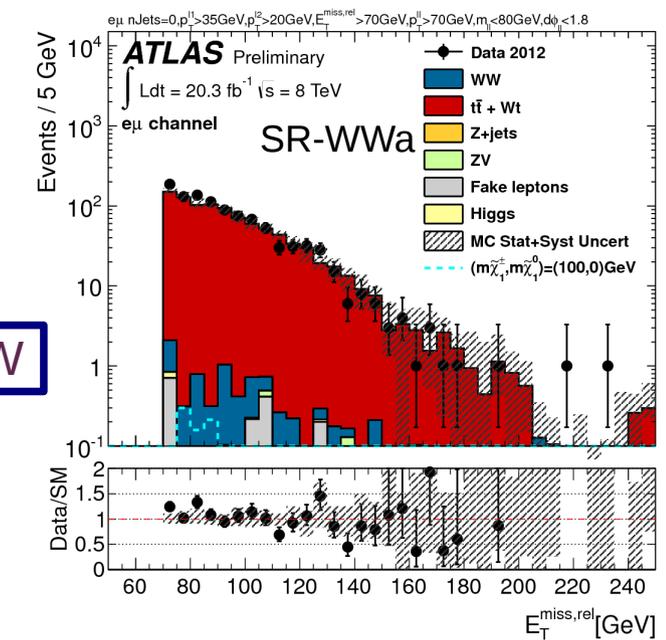
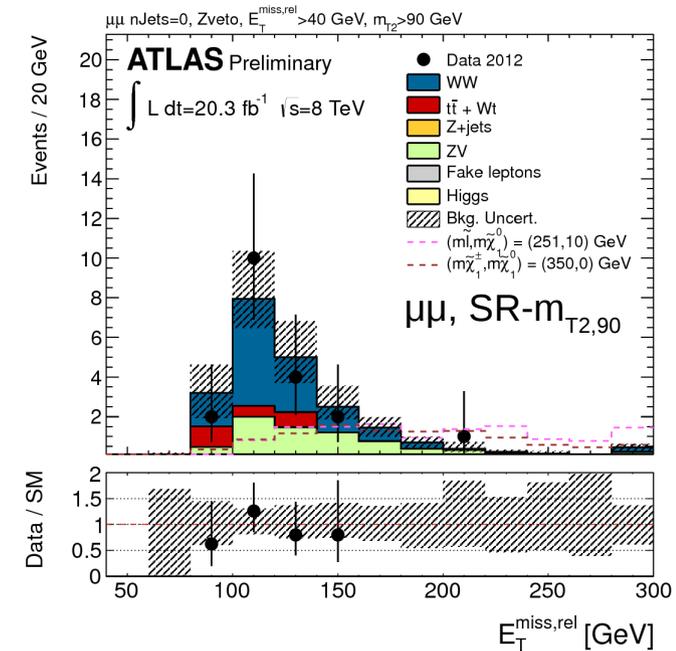
Large $E_T^{miss,rel}$ or m_{T2}

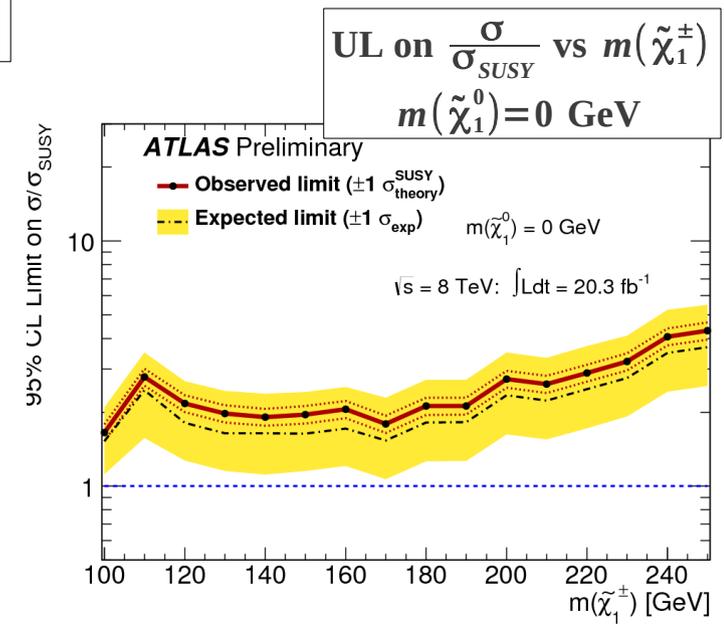
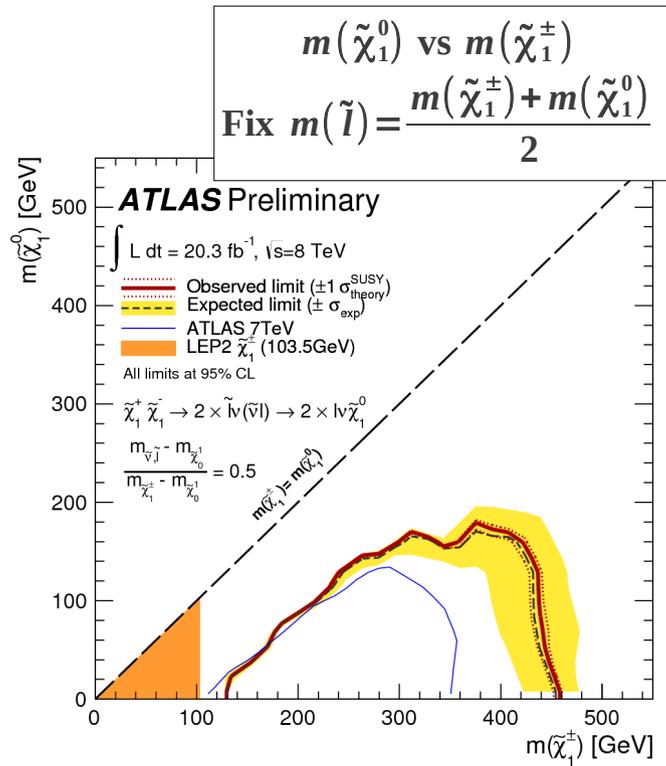
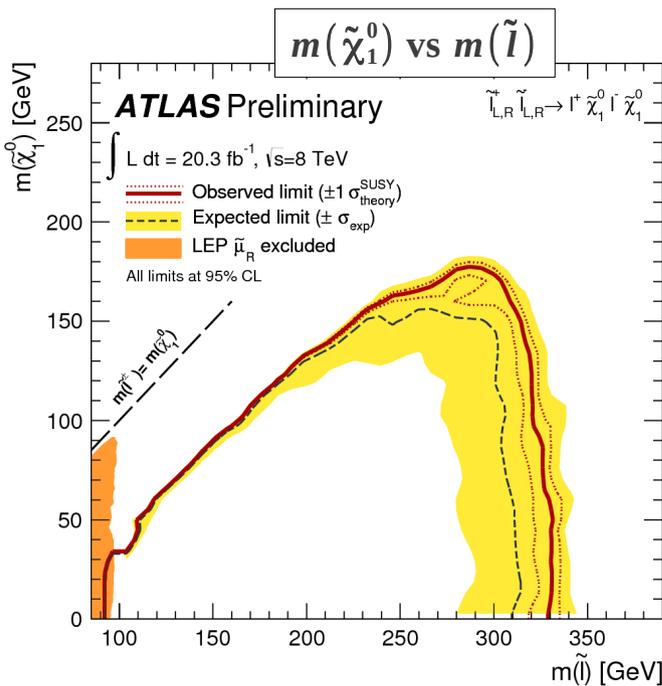
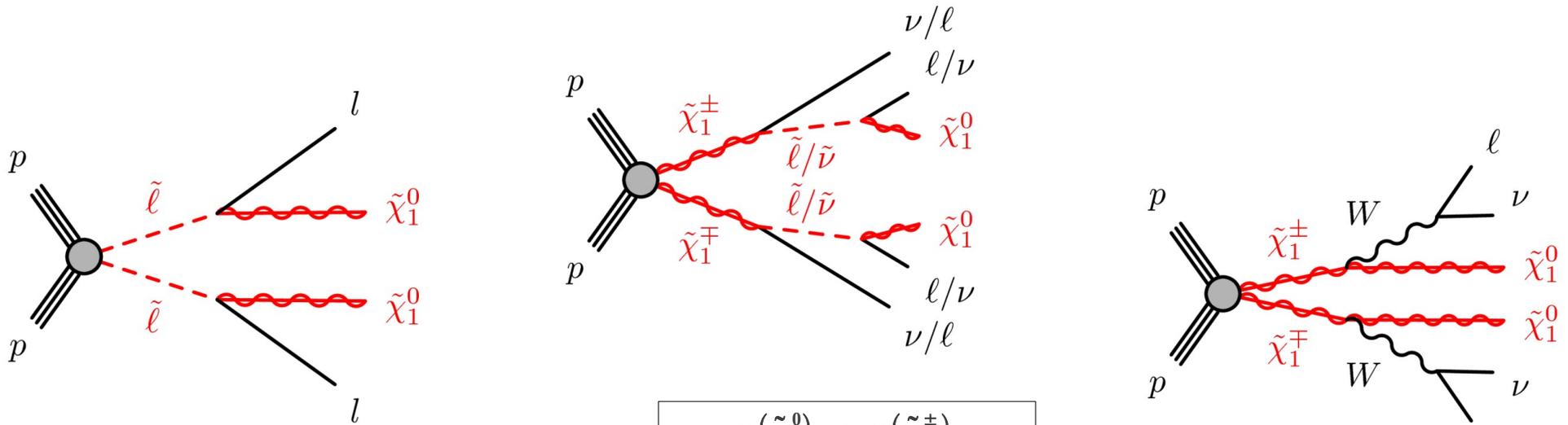
$E_T^{miss,rel}$ projection of E_T^{miss} on perpendicular axis

Reduces mis-measured E_T^{miss}

$$E_T^{miss,rel} = \begin{cases} E_T^{miss} & \text{if } \Delta\phi_{\ell,j} \geq \pi/2 \\ E_T^{miss} \times \sin \Delta\phi_{\ell,j} & \text{if } \Delta\phi_{\ell,j} < \pi/2 \end{cases}$$

Ttbar and **WW** dominated

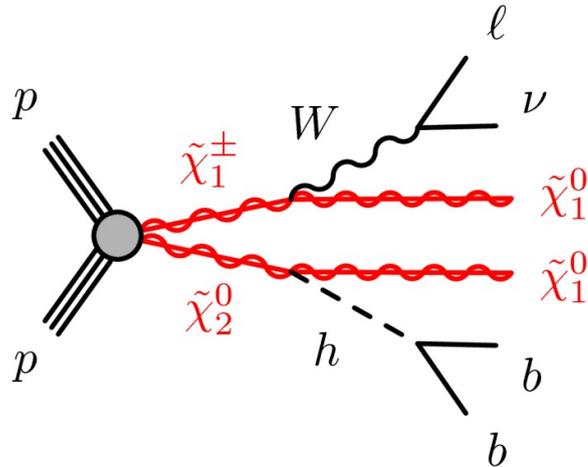




Consider SUSY processes with decays via 125 GeV SM-like higgs

Where kinematically allowed, $\tilde{\chi}_2^0 \rightarrow h \tilde{\chi}_1^0$ could be a significant fraction of the BR

125 GeV SM-like lightest higgs decays to pair of bottom quarks with highest BR



2 signal regions

=2 b-tagged jets $\sim m_H$

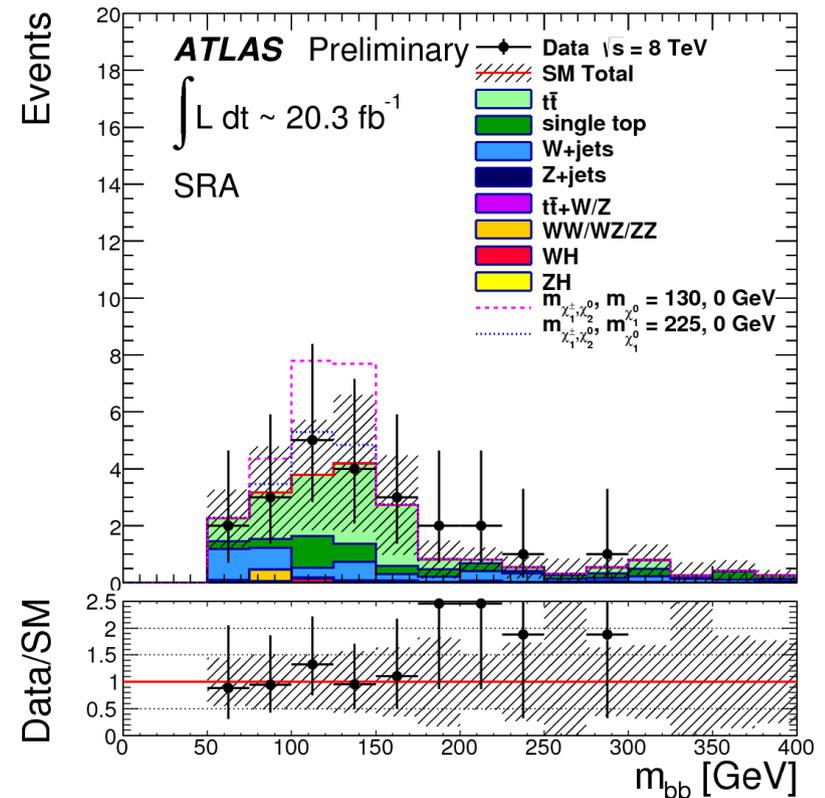
Large E_T^{miss} , m_T and m_{CT}

$$m_{CT}^2 = (E_T^{b1} + E_T^{b2})^2 - |\mathbf{p}_T^{b1} - \mathbf{p}_T^{b2}|^2$$

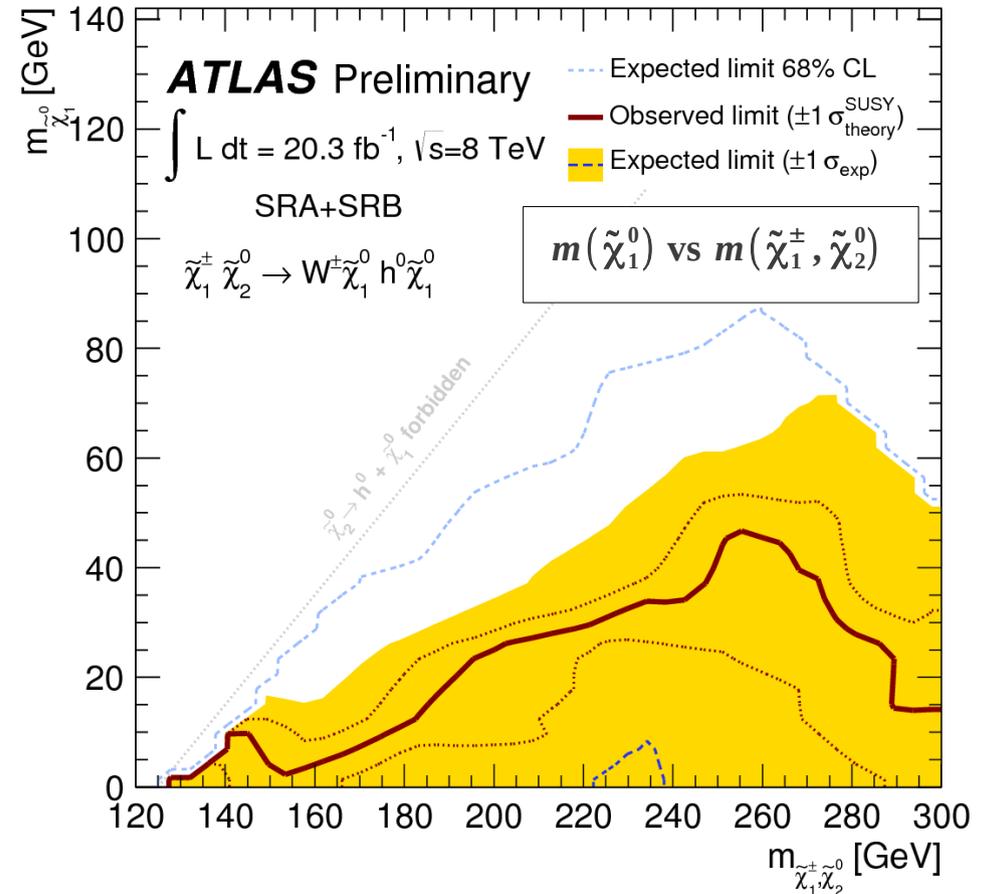
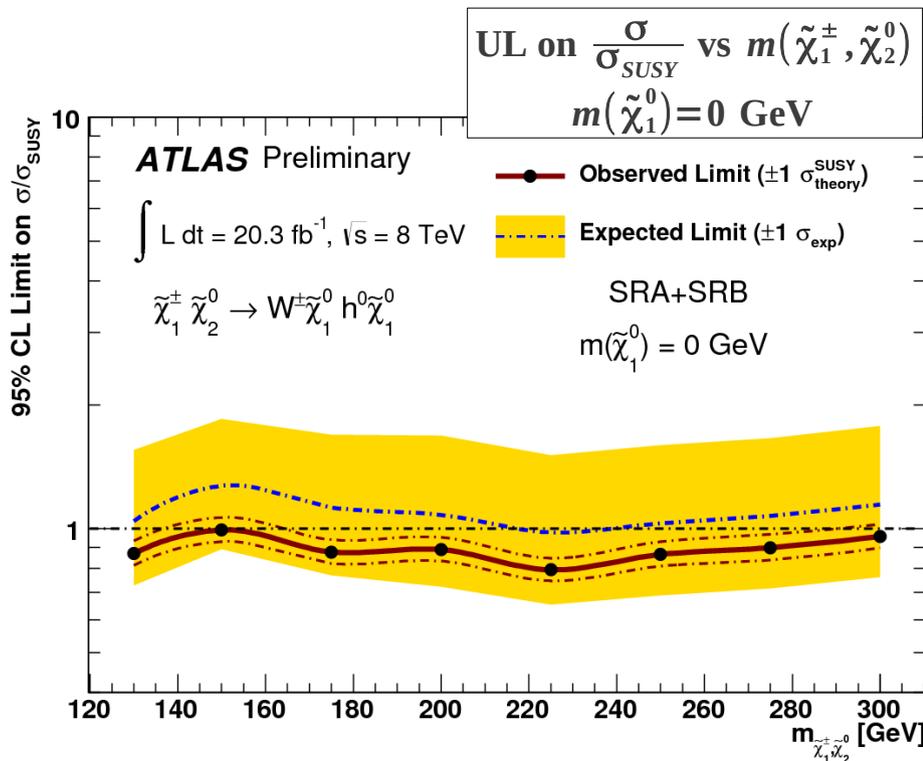
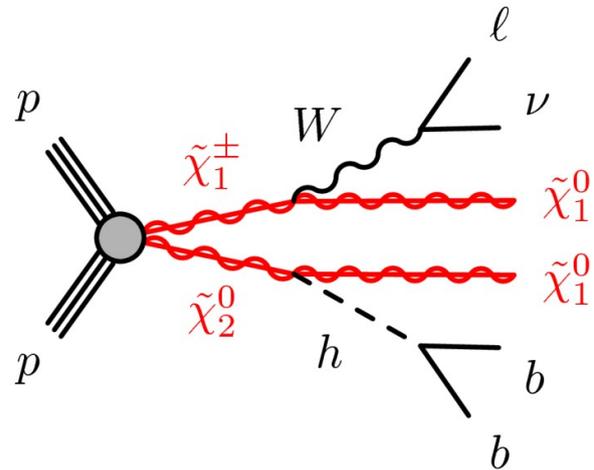
Ttbar and W+jets dominated background

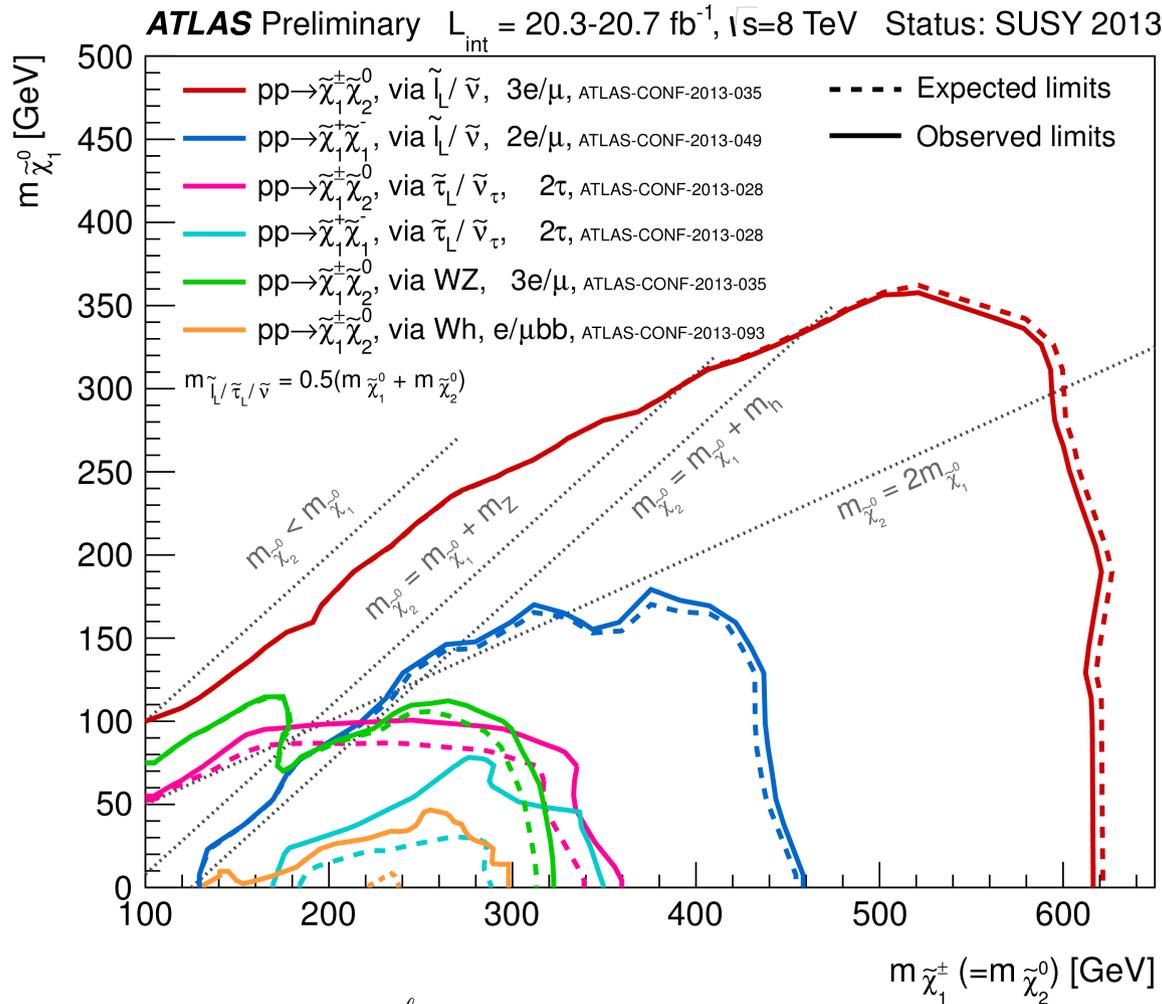
Simultaneous background fit in control and signal regions

Fit using 8 bins in $m_{bb} \rightarrow$ exclude m_{bb} 105–135 GeV to avoid signal contamination (for background only fit, not for discovery or exclusion fit)

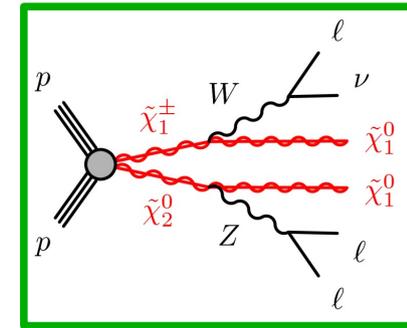
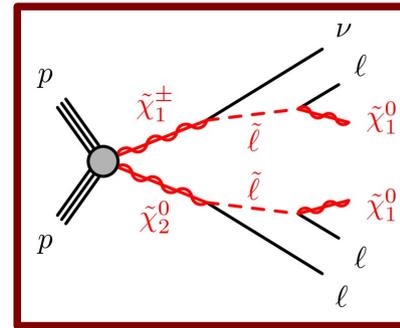


No significant excess seen

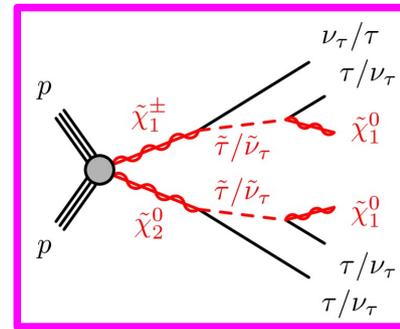




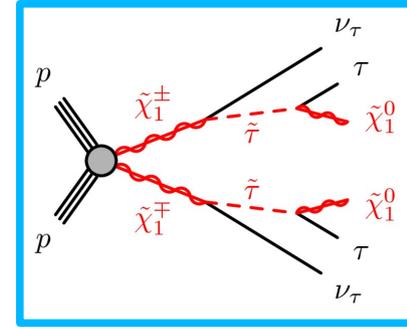
3-leptons $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm$



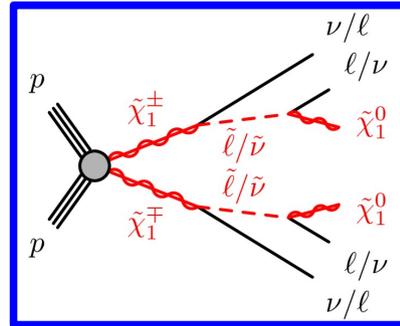
2-taus $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm$



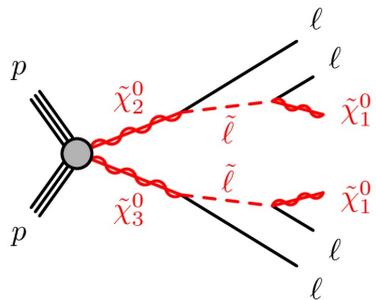
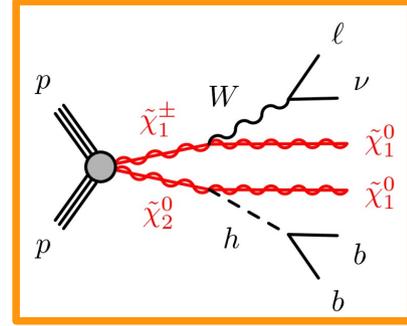
$\tilde{\chi}_1^+ \tilde{\chi}_1^-$



2-leptons $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm$



1-lepton + bb $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm$



+ 4-lepton interpretation
for $\tilde{\chi}_2^0 \tilde{\chi}_3^0$

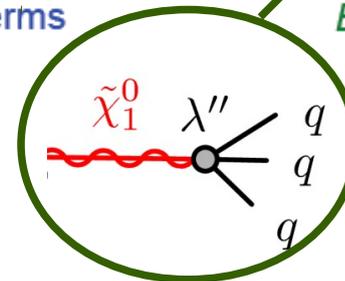
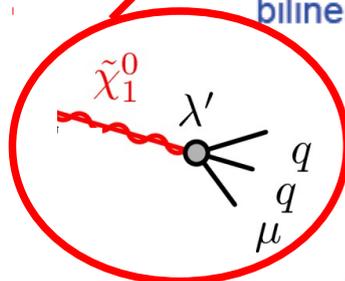
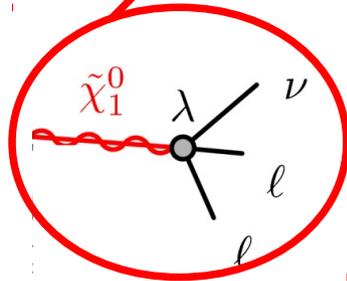
Recall

The MSSM potential

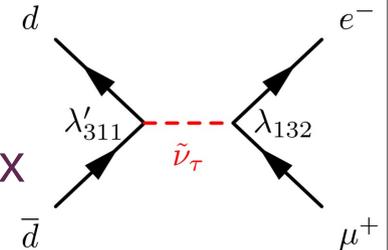
$$W_{RPV} = \lambda_{ijk} \hat{L}_i \hat{L}_j \hat{E}_k^C + \lambda'_{ijk} \hat{L}_i \hat{Q}_j \hat{D}_k^C + \epsilon_i \hat{L}_i \hat{H}_u + \lambda''_{ijk} \hat{U}_i^C \hat{D}_j^C \hat{D}_k^C$$

L-number violating terms
bilinear terms
B-number violating terms

RPV couplings $\lambda, \lambda', \lambda''$



RPV can be at the production vertex and/or at decay vertices

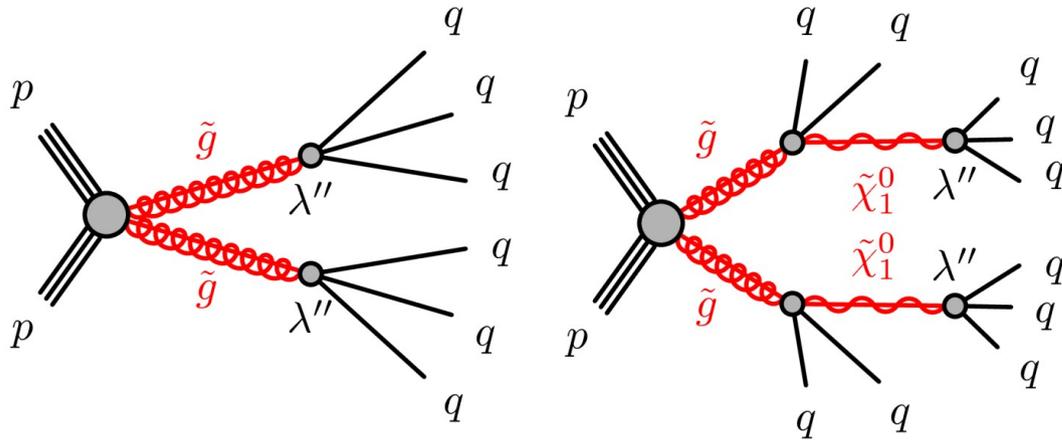


RPV search channels

- » Multijets (2x3 jets)
- Heavy resonance to $e\mu$, $e\tau$, $\mu\tau$ PLB 723 (2013) 15
- 4-leptons ATLAS-CONF-2013-036

>> new result to be discussed

(since last ATLAS SUSY seminar on 26th March 2013)



Consider gluino or neutralino LSP decaying to jets

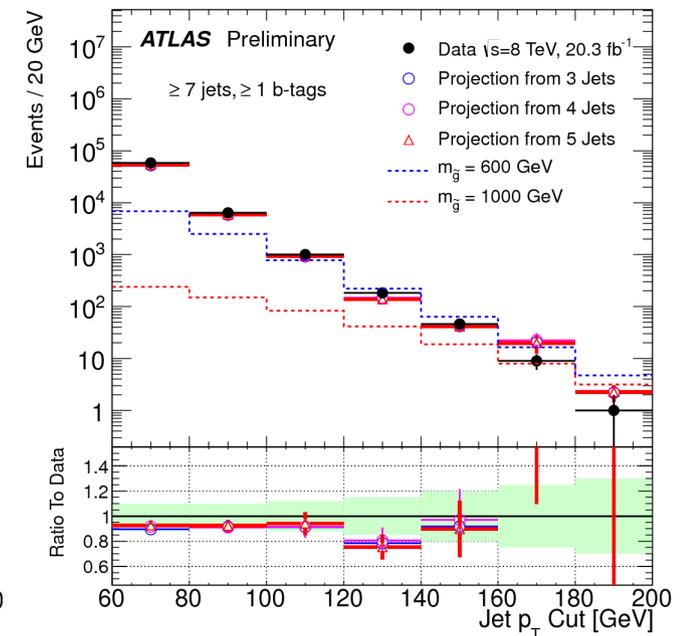
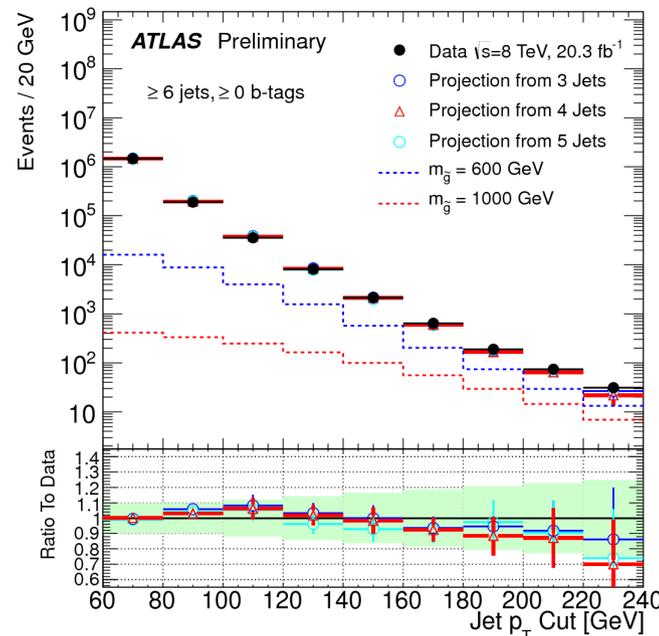
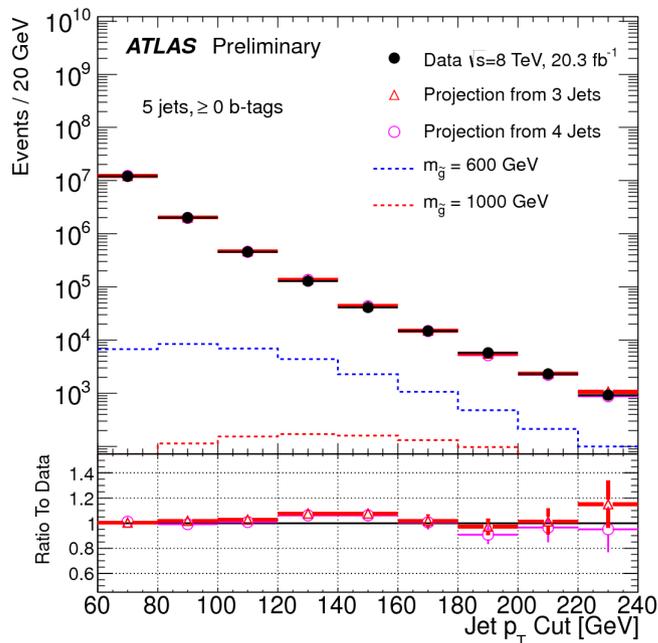
Searching for resonances is difficult due to combinatorics

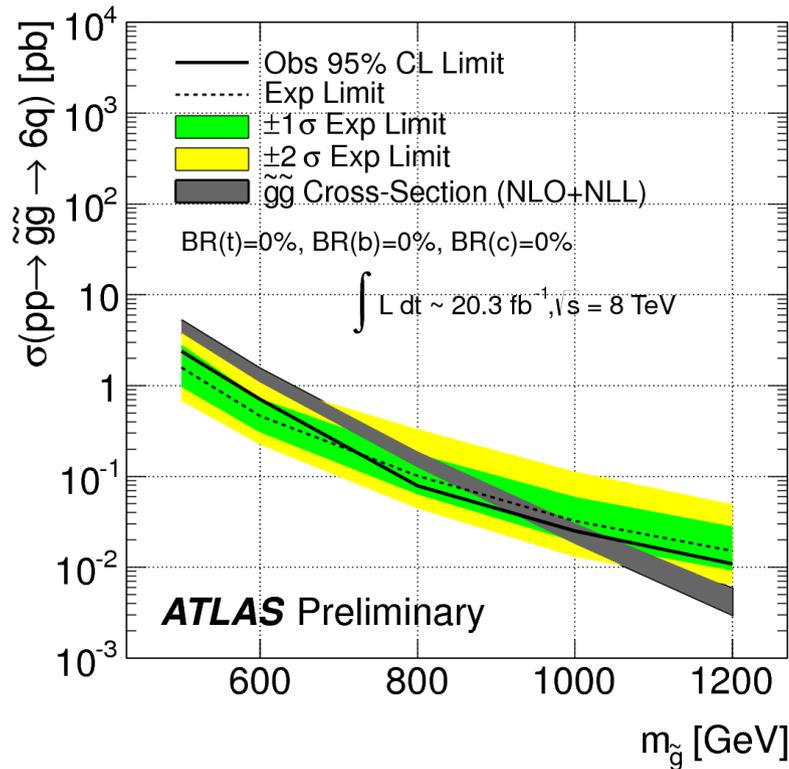
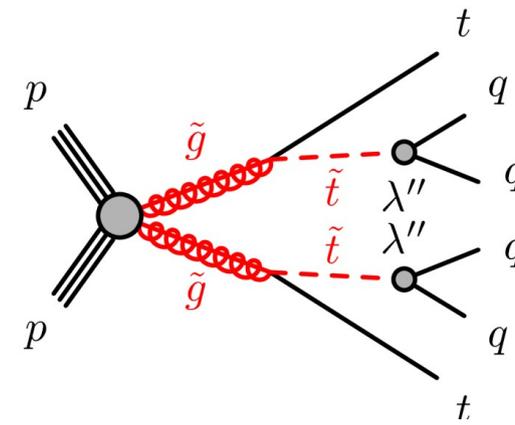
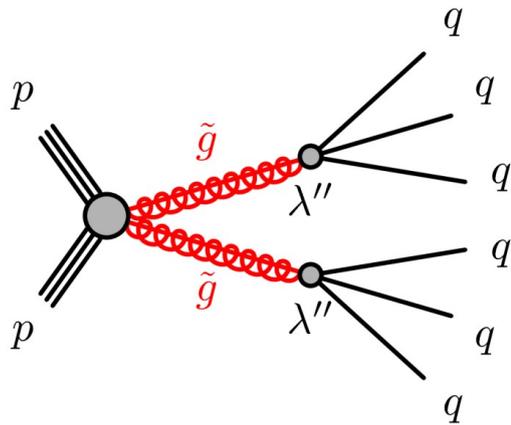
Search for events with ≥ 6 or 7 high p_T jets
 0–2 b-tagged jets to estimate BR to heavy flavour quarks

SM multijet background normalised to data in lower jet multiplicity regions

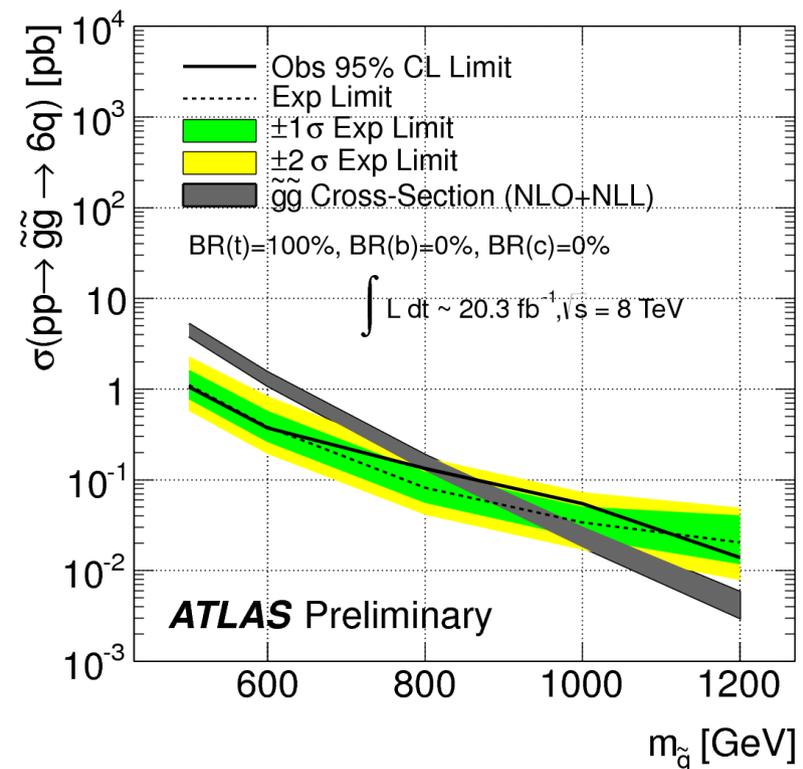
Signal regions optimised for many different models
 Jet $p_T \geq 80$ –220 GeV, $N_{\text{jet}} \geq 6$ –7, $N_{\text{bjet}} 0$ –2

No significant excess seen





Glueinos excluded up to ~ 900 GeV
for $BR(\text{heavy quarks}) = 0\%$



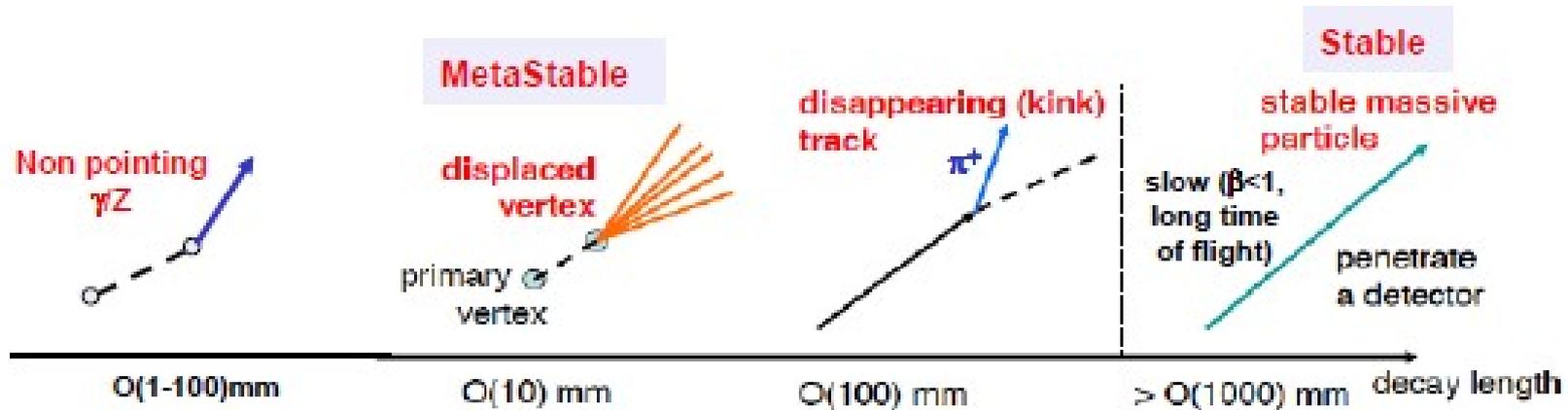
Glueinos excluded up to ~ 800 GeV
for $BR(\text{top}) = 100\%$

Recall

RPV also lead to **non-prompt decays** if λ couplings are small

Long-lived SUSY particles can also arise from

- Heavy mediator sparticles e.g. Split SUSY
- Mass degeneracy
- Weak couplings



RPV & LL search channels

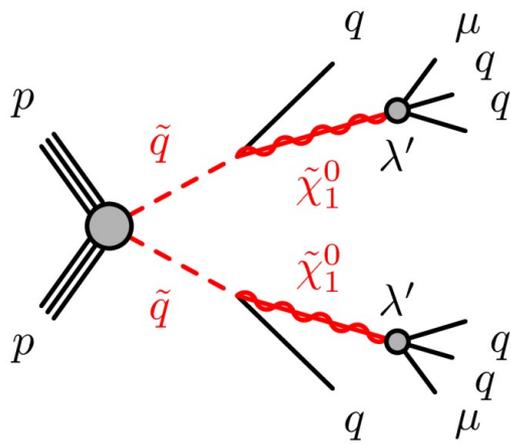
Look for RPV signatures and long-lived signatures

Cover wide coverage of lifetimes

- » Disappearing track
- » Stopped gluino
- » Long lived slepton
- » Displaced vertex
- Non-pointing photon PRD 88, 012001 (2013)

>> new result to be discussed

(since last ATLAS SUSY seminar on 26th March 2013)

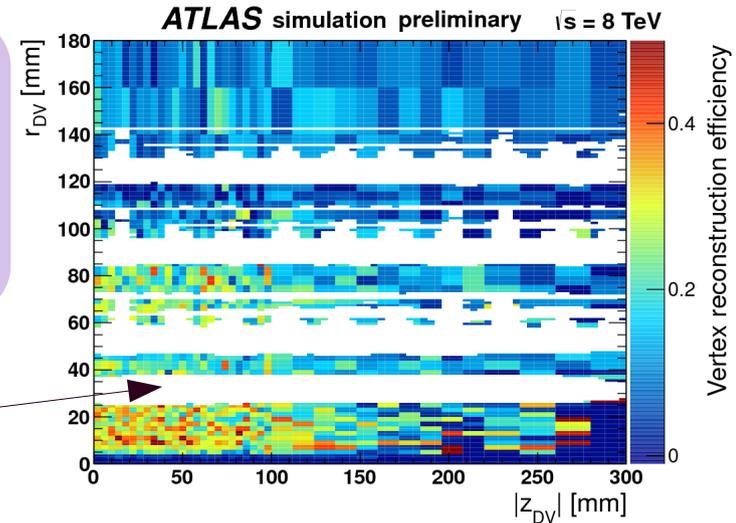


Consider the LSP to be long-lived, decaying to a muon and jets
Dedicated reconstruction of tracks and vertices

Trigger with one high- p_T muon

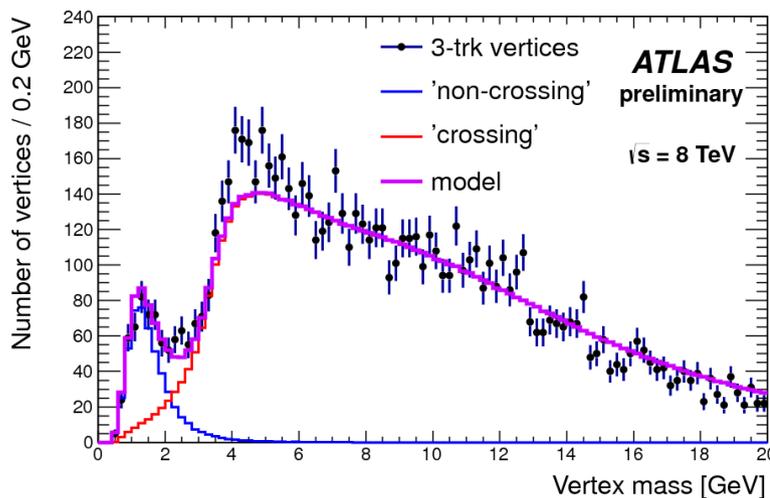
Search for a displaced vertex (DV)
within $r < 180$ mm and $|z| < 300$ mm
 $m_{DV} > 10$ GeV and > 4 tracks

To suppress hadronic interactions, veto
vertices from regions of high density



Dominating background from hadronic interactions with
gas molecules (outside beampipe)

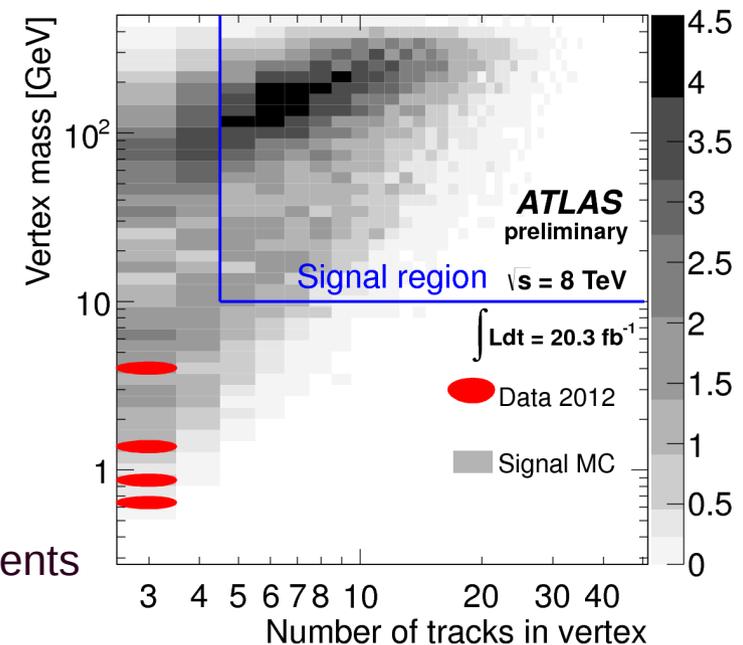
Usually low mass, but random track crossing can give high mass



Model m_{DV} with jet-
triggered events

Random track combination
background negligible

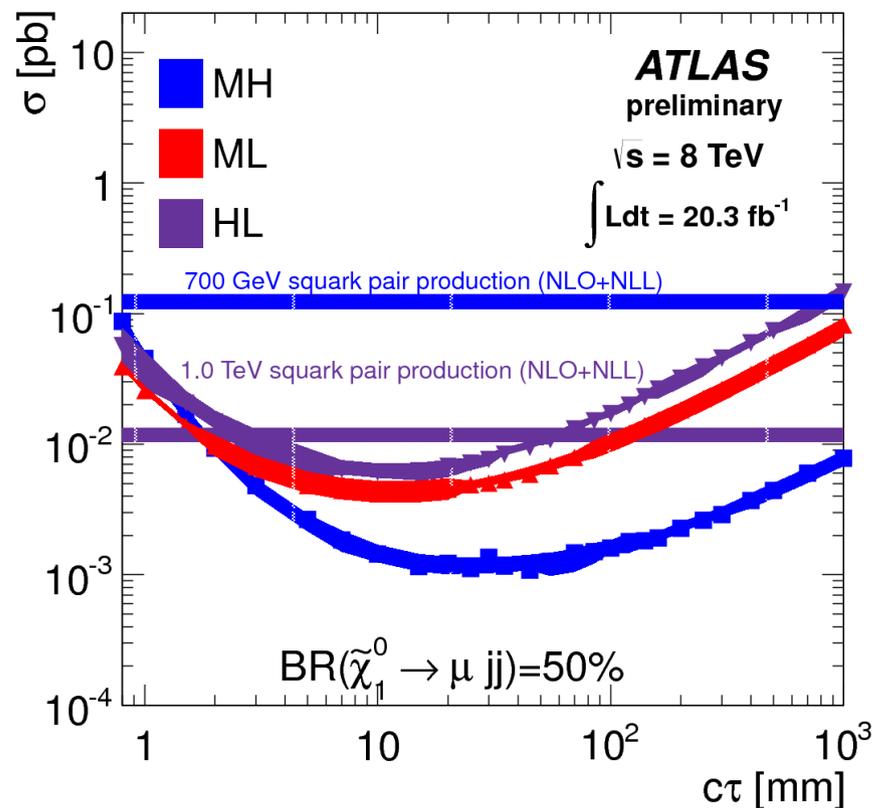
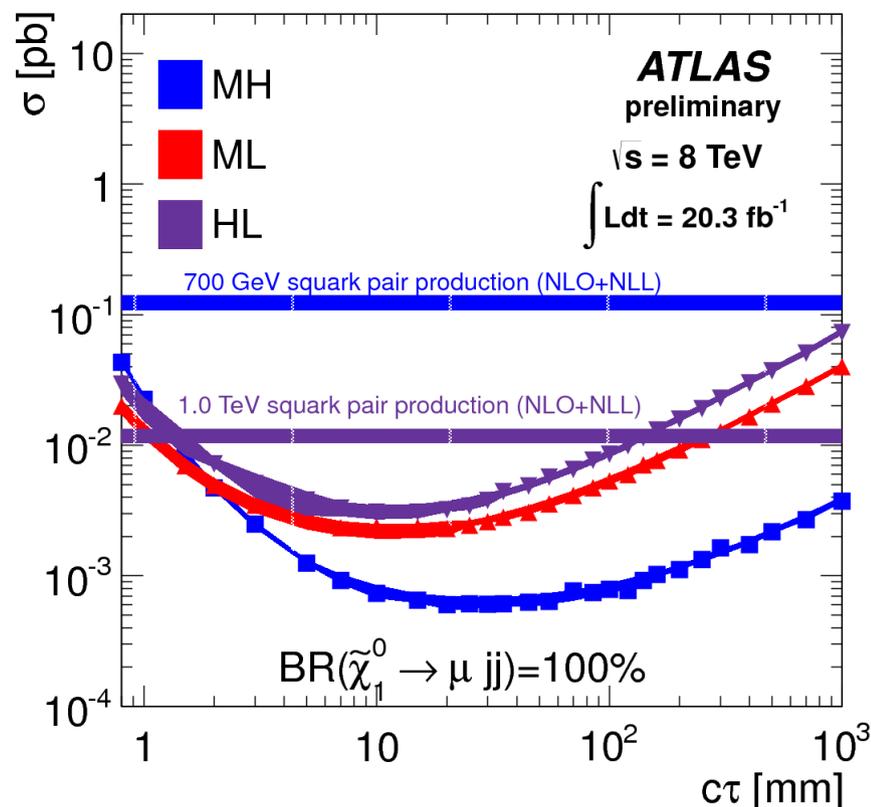
Expected 0.02 ± 0.02 events
Observed 0



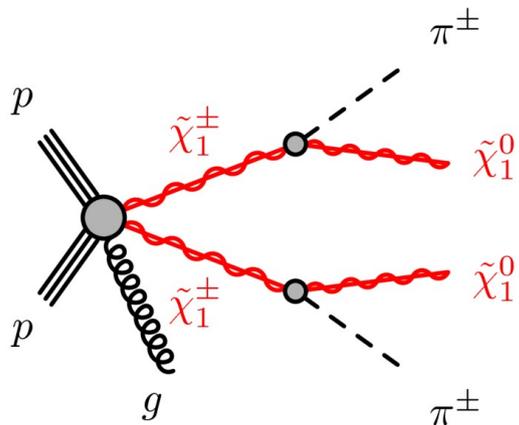
Use efficiency maps to set limits on range of $c\tau$

Medium-Heavy
Medium-Light
Heavy-Light

Sample	$m_{\tilde{q}}$ [GeV]	$m_{\tilde{\chi}_1^0}$ [GeV]	σ [fb]	$\langle\gamma\beta\rangle_{\tilde{\chi}_1^0}$	$c\tau_{MC}$ [mm]	λ'_{211}
MH	700	494	124.3	1.0	175	0.2×10^{-5}
ML	700	108	124.3	3.1	101	1.5×10^{-5}
HL	1000	108	11.9	5.5	220	20.0×10^{-5}



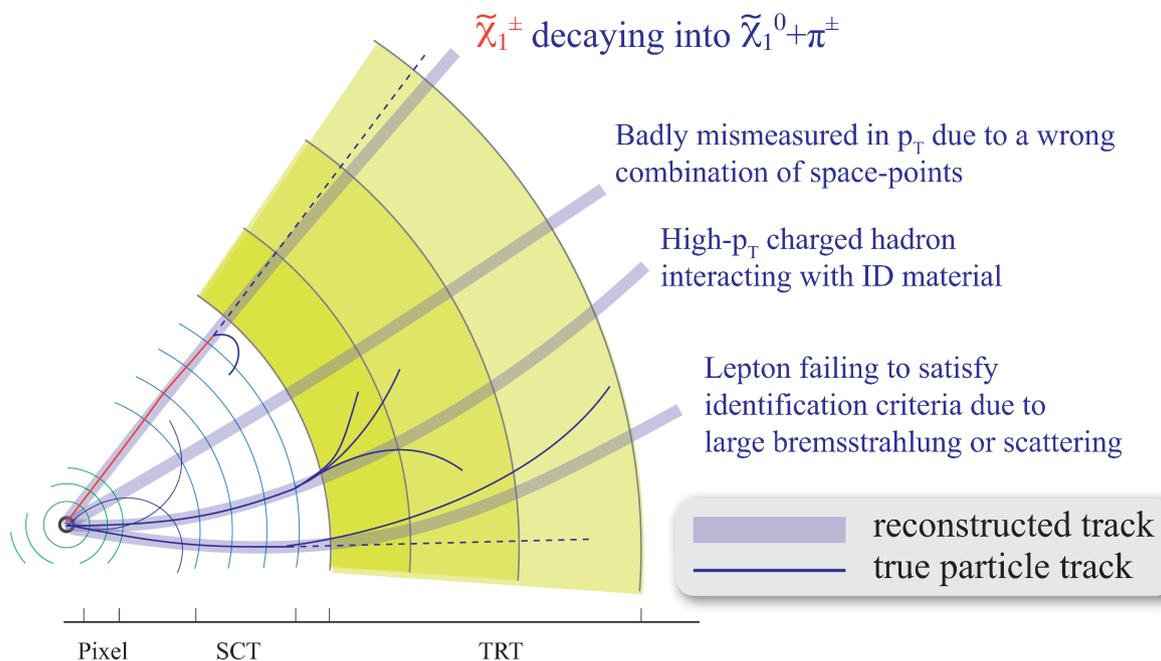
1 TeV squarks excluded for $1.5 < c\tau < 156 \text{ mm}$
100% BR to 108 GeV LSP



Many SUSY model e.g. AMSB have almost mass degenerate chargino and LSP → long-lived chargino

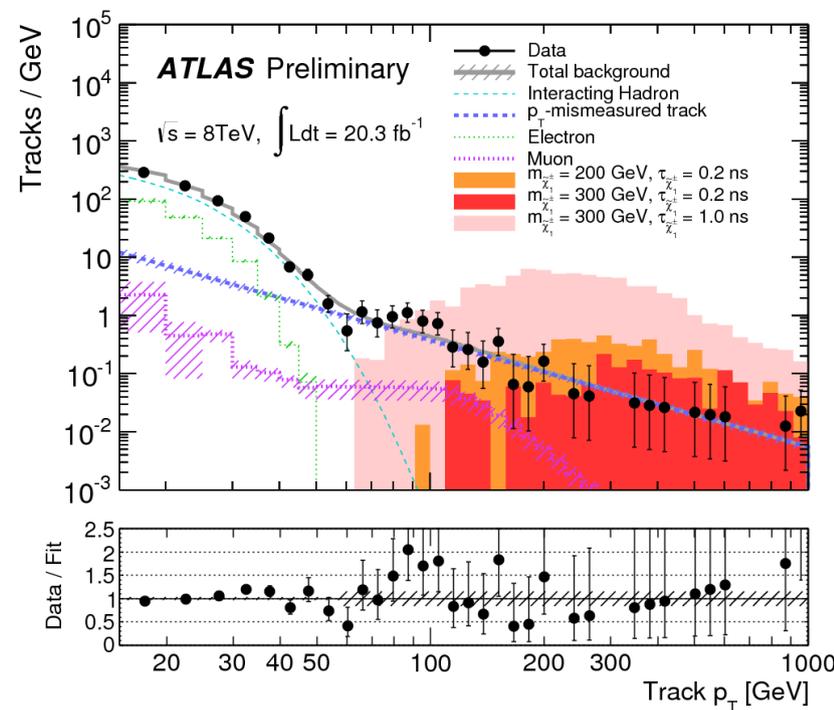
Chargino travels into detector before decaying to soft pion + LSP → disappearing track

Trigger on ISR jet
Look for isolated, high p_T tracks with < 5 TRT hits



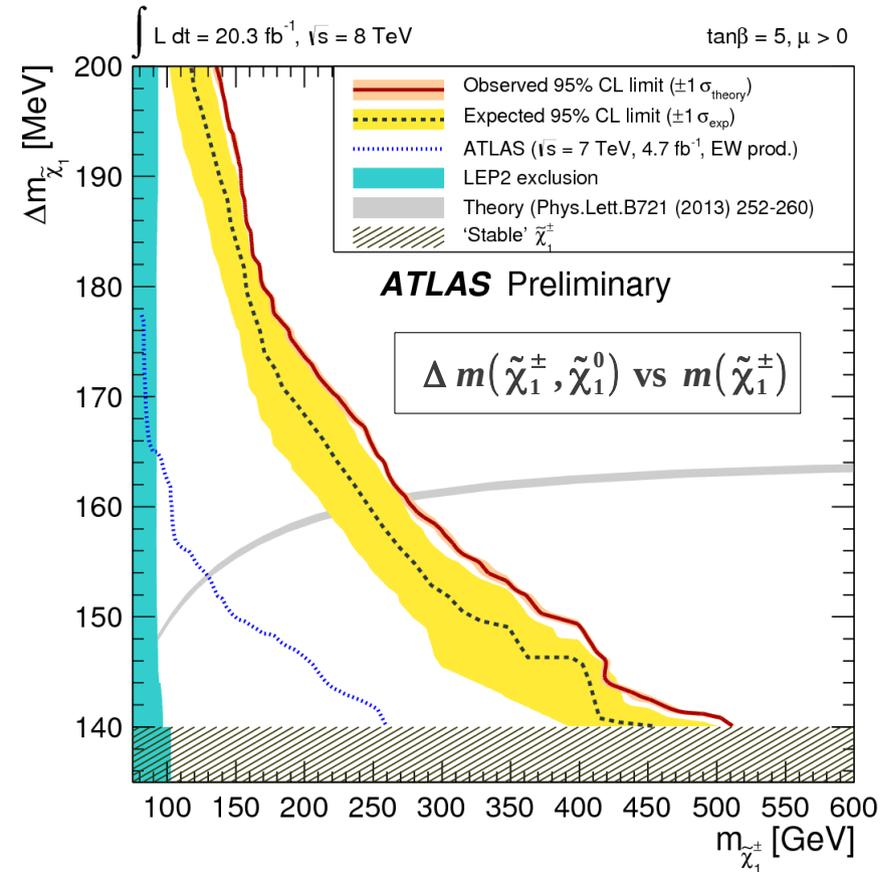
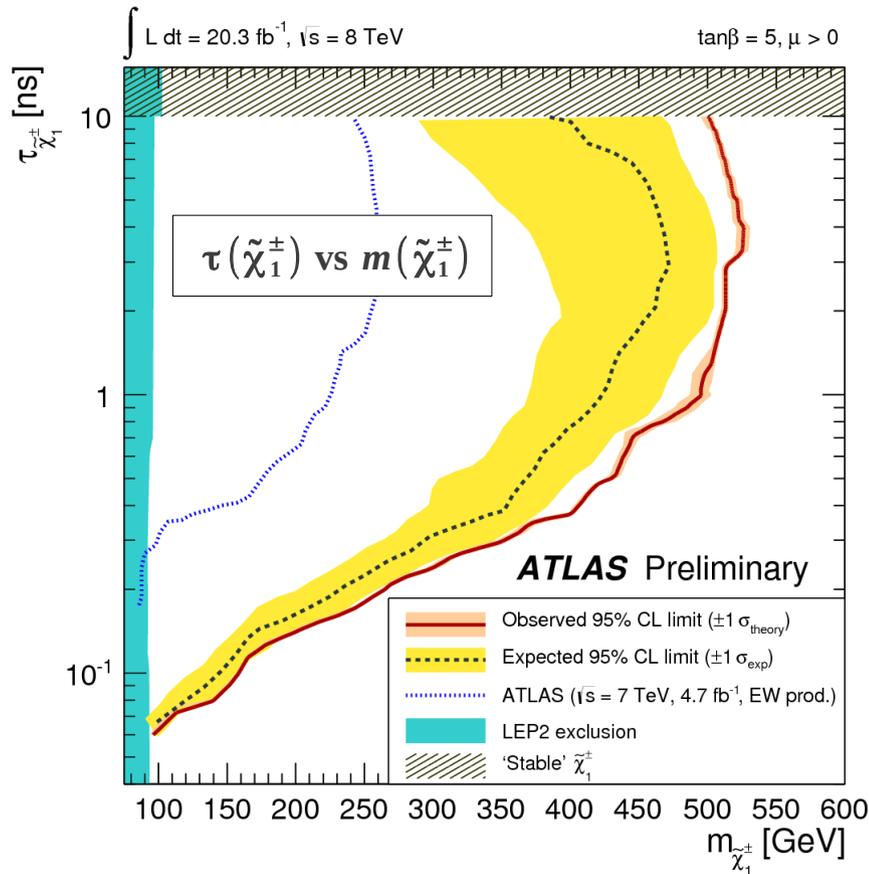
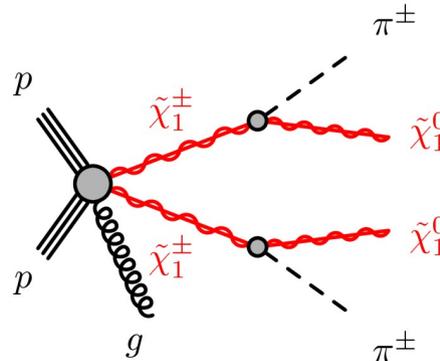
Background track p_T shape taken from data

Signal + background template fit for candidate tracks

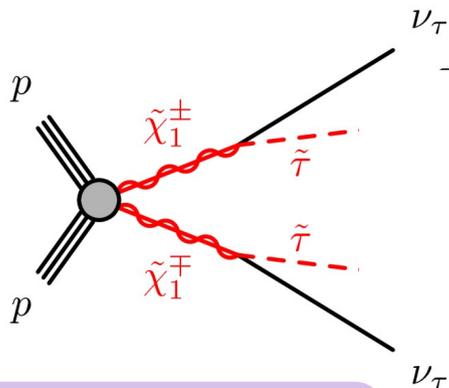


No significant excess observed

AMSB model with $\Delta m \sim 160$ MeV.



For $\tau \sim 0.2$ ns, charginos excluded up to 270 GeV
 For $\tau \sim 1\text{--}10$ ns, chargino excluded up to 520 GeV



Many SUSY model e.g. GMSB have long-lived stau NLSP

If $c\tau >$ few metres, LLP looks like a heavy muon

Low β

Large dE/dx

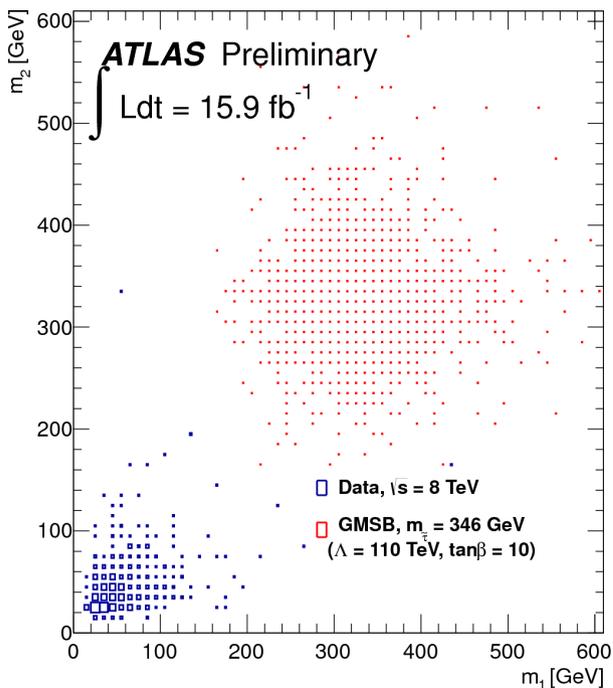
2 LLP per event

$$m = \frac{p}{\beta \gamma}$$

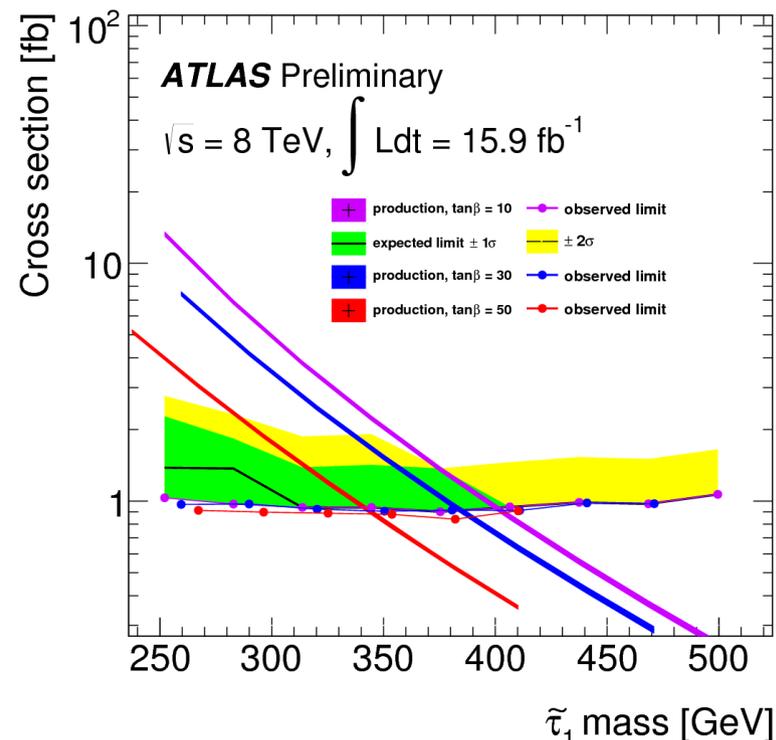
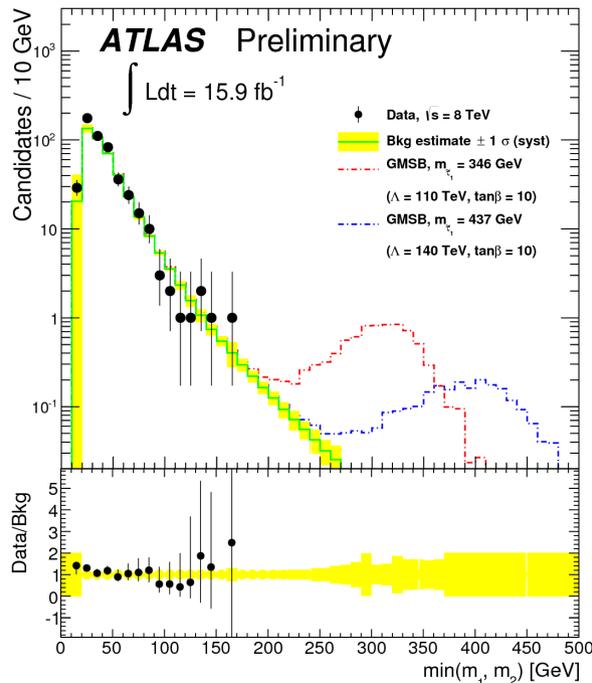
from track
from pixel (dE/dx) + calorimeter + muon detector (Time of Flight)

Signal region
2 high p_T muons
 $0.2 < \beta < 0.95$

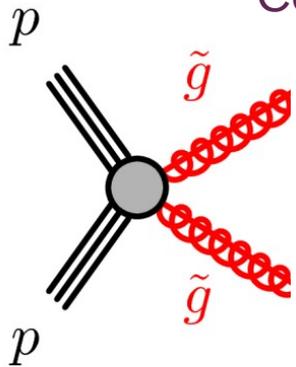
Background dominated by high p_T muons with mis-measured β
Taken from data



No significant excess observed



Sensitivity to stau mass depends on GMSB model parameters

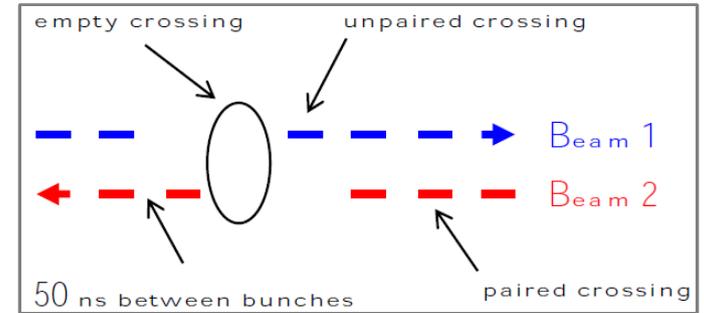


Consider gluino with long lifetime that forms R-hadrons with SM vacuum quarks

R-hadron travels through detector
 Gets stuck in detector via dE/dx energy loss and nuclear scattering
 Can decay at a much later time!

Also relevant for Split SUSY with a high mass intermediate squark

Use empty bunches in LHC beams to look for hadronic activity



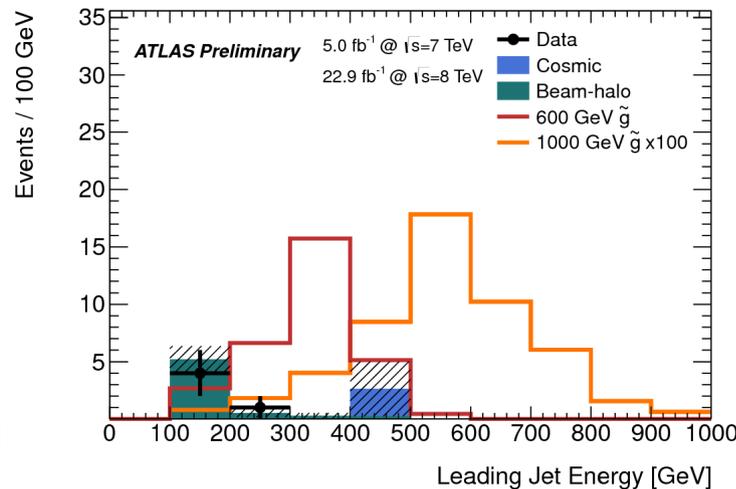
- Detection efficiency depends on
- stopping fraction
 - probability to decay in empty bunch
 - reconstruction efficiency

Signal region

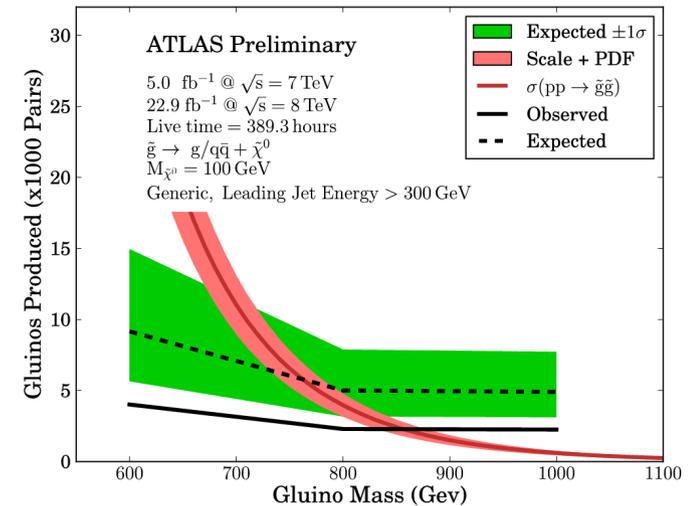
Jet + $E_{\text{T}}^{\text{miss}}$ empty bunch trigger
 Jet $E > 100$ or 300 GeV
 Veto muon activity

Cosmic muon background from low-luminosity run period

Beam halo background from unpaired crossings

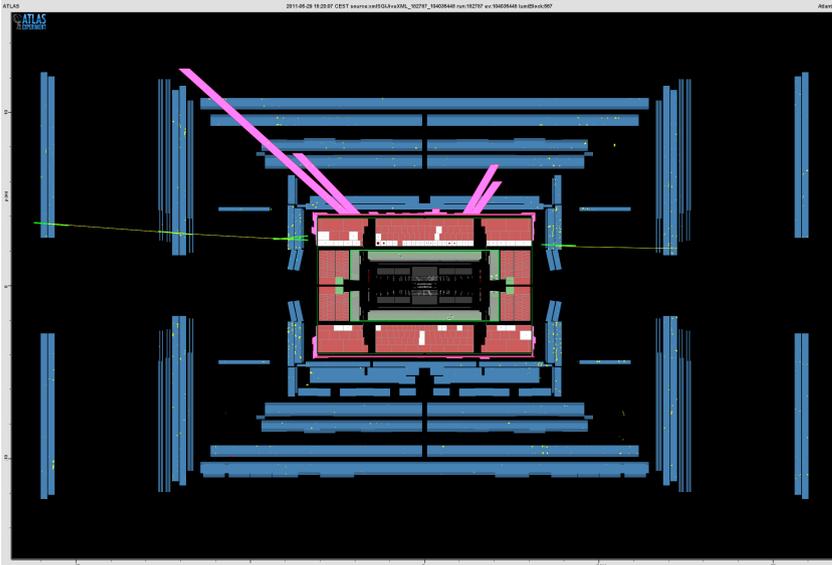


No significant excess seen

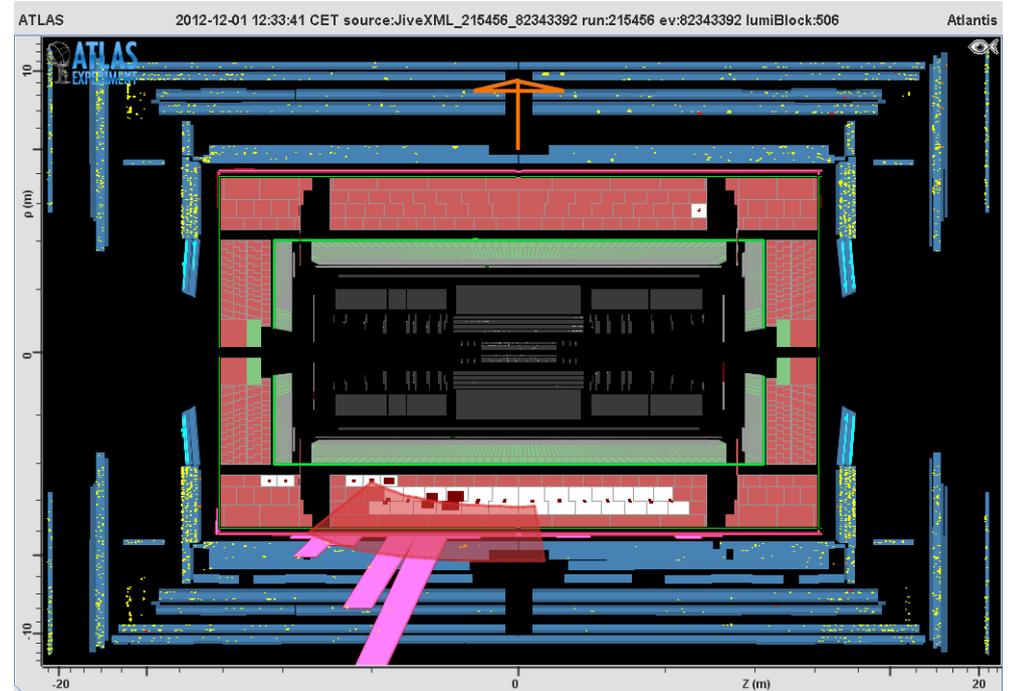


Stopped gluino R-hadrons excluded up to 832 GeV for $10 \mu\text{s} < \tau(\text{gl}\tilde{g}) < 1000 \text{ s}$ and 100 GeV LSP
 Longest lifetime excluded: 2 years!

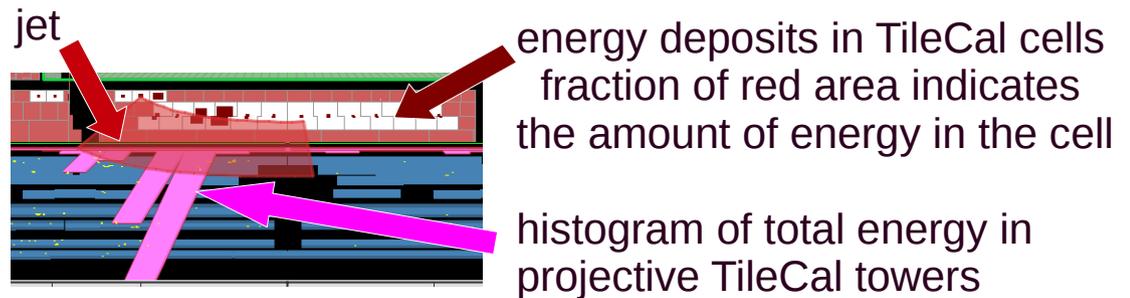
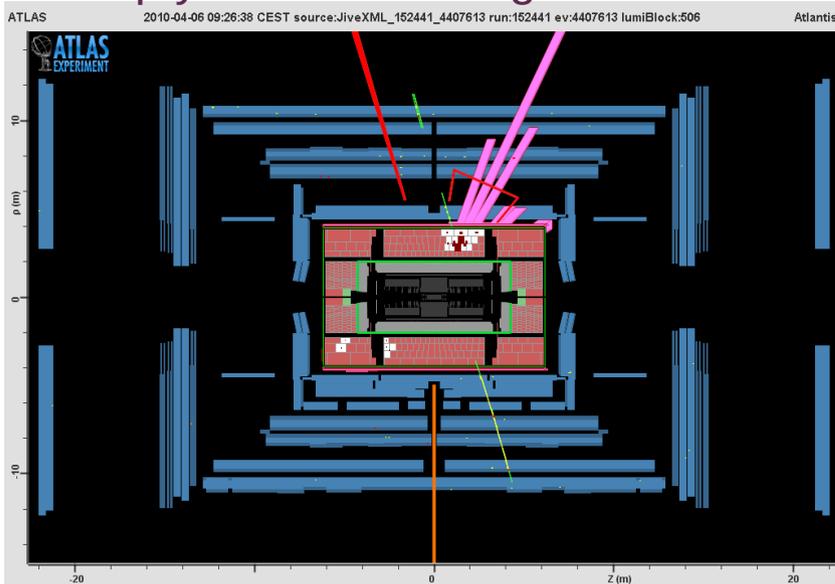
A beam-halo candidate event during an unpaired bunch crossing in data.



A candidate event display from 2011 data passing all selections



A cosmic ray muon candidate event during an empty bunch crossing in data



Muon segments are drawn but not reconstructed

Grand Summary of ATLAS SUSY Search Results

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: SUSY 2013

ATLAS Preliminary

$$\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1} \quad \sqrt{s} = 7, 8 \text{ TeV}$$

Increasing σ

Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference	
Inclusive Searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	\tilde{q}, \tilde{g} 1.7 TeV	$m(\tilde{q})=m(\tilde{g})$ ATLAS-CONF-2013-047
	MSUGRA/CMSSM	1 e, μ	3-6 jets	Yes	20.3	\tilde{g} 1.2 TeV	any $m(\tilde{q})$ ATLAS-CONF-2013-062
	MSUGRA/CMSSM	0	7-10 jets	Yes	20.3	\tilde{g} 1.1 TeV	any $m(\tilde{q})$ 1308.1841
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow \tilde{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{q} 740 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-047
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{q}\tilde{q}\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{g} 1.3 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-047
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{q}\tilde{q}\tilde{\chi}_1^0 \rightarrow \tilde{q}\tilde{q}W^\pm\tilde{\chi}_1^0$	1 e, μ	3-6 jets	Yes	20.3	\tilde{g} 1.18 TeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}, m(\tilde{\chi}^\pm)=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$ ATLAS-CONF-2013-062
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{q}\tilde{q}(\ell\ell/\ell\nu/\nu\nu)\tilde{\chi}_1^0$	2 e, μ	0-3 jets	-	20.3	\tilde{g} 1.12 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-089
	GMSB ($\tilde{\ell}$ NLSP)	2 e, μ	2-4 jets	Yes	4.7	\tilde{g} 1.24 TeV	$\tan\beta < 15$ 1208.4688
	GMSB ($\tilde{\ell}$ NLSP)	1-2 τ	0-2 jets	Yes	20.7	\tilde{g} 1.4 TeV	$\tan\beta > 18$ ATLAS-CONF-2013-026
	GGM (bino NLSP)	2 γ	-	Yes	4.8	\tilde{g} 1.07 TeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$ 1209.0753
GGM (wino NLSP)	1 e, $\mu + \gamma$	-	Yes	4.8	\tilde{g} 619 GeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$ ATLAS-CONF-2012-144	
GGM (higgsino-bino NLSP)	γ	1 b	Yes	4.8	\tilde{g} 900 GeV	$m(\tilde{\chi}_1^0) > 220 \text{ GeV}$ 1211.1167	
GGM (higgsino NLSP)	2 e, μ (Z)	0-3 jets	Yes	5.8	\tilde{g} 690 GeV	$m(\tilde{H}) > 200 \text{ GeV}$ ATLAS-CONF-2012-152	
Gravitino LSP	0	mono-jet	Yes	10.5	\tilde{g} 645 GeV	$m(\tilde{g}) > 10^{-4} \text{ eV}$ ATLAS-CONF-2012-147	
3 rd gen. \tilde{g} med.	$\tilde{g} \rightarrow b\tilde{\chi}_1^0$	0	3 b	Yes	20.1	\tilde{g} 1.2 TeV	$m(\tilde{\chi}_1^0) < 600 \text{ GeV}$ ATLAS-CONF-2013-061
	$\tilde{g} \rightarrow t\tilde{\chi}_1^0$	0	7-10 jets	Yes	20.3	\tilde{g} 1.1 TeV	$m(\tilde{\chi}_1^0) < 350 \text{ GeV}$ 1308.1841
	$\tilde{g} \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ	3 b	Yes	20.1	\tilde{g} 1.34 TeV	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$ ATLAS-CONF-2013-061
	$\tilde{g} \rightarrow b\tilde{\chi}_1^0$	0-1 e, μ	3 b	Yes	20.1	\tilde{g} 1.3 TeV	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}$ ATLAS-CONF-2013-061
3 rd gen. squarks direct production	$\tilde{b}_1, \tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$	0	2 b	Yes	20.1	\tilde{b}_1 100-620 GeV	$m(\tilde{\chi}_1^0) < 90 \text{ GeV}$ 1308.2631
	$\tilde{b}_1, \tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^\pm$	2 e, μ (SS)	0-3 b	Yes	20.7	\tilde{b}_1 275-430 GeV	$m(\tilde{\chi}_1^\pm)=2m(\tilde{\chi}_1^0)$ ATLAS-CONF-2013-007
	\tilde{t}_1, \tilde{t}_1 (light), $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$	1-2 e, μ	1-2 b	Yes	4.7	\tilde{t}_1 110-167 GeV	$m(\tilde{\chi}_1^0)=55 \text{ GeV}$ 1208.4305, 1209.2102
	\tilde{t}_1, \tilde{t}_1 (light), $\tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$	2 e, μ	0-2 jets	Yes	20.3	\tilde{t}_1 130-220 GeV	$m(\tilde{\chi}_1^0)=m(\tilde{t}_1)-m(W)-50 \text{ GeV}, m(\tilde{t}_1) < m(\tilde{\chi}_1^\pm)$ ATLAS-CONF-2013-048
	\tilde{t}_1, \tilde{t}_1 (medium), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	2 e, μ	2 jets	Yes	20.3	\tilde{t}_1 225-525 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-065
	\tilde{t}_1, \tilde{t}_1 (medium), $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$	0	2 b	Yes	20.1	\tilde{t}_1 150-580 GeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}, m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)=5 \text{ GeV}$ 1308.2631
	\tilde{t}_1, \tilde{t}_1 (heavy), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^\pm$	1 e, μ	1 b	Yes	20.7	\tilde{t}_1 200-610 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-037
	\tilde{t}_1, \tilde{t}_1 (heavy), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	0	2 b	Yes	20.5	\tilde{t}_1 320-660 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-024
	$\tilde{t}_1, \tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$	0	mono-jet/c-tag	Yes	20.3	\tilde{t}_1 90-200 GeV	$m(\tilde{t}_1)-m(\tilde{\chi}_1^0) < 85 \text{ GeV}$ ATLAS-CONF-2013-068
	\tilde{t}_1, \tilde{t}_1 (natural GMSB)	2 e, μ (Z)	1 b	Yes	20.7	\tilde{t}_1 500 GeV	$m(\tilde{\chi}_1^0) > 150 \text{ GeV}$ ATLAS-CONF-2013-025
$\tilde{t}_2, \tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 e, μ (Z)	1 b	Yes	20.7	\tilde{t}_2 271-520 GeV	$m(\tilde{t}_1)=m(\tilde{\chi}_1^0)+180 \text{ GeV}$ ATLAS-CONF-2013-025	
EW direct	$\tilde{\ell}_L, \tilde{\ell}_L, \tilde{\ell}_L, \tilde{\ell}_L \rightarrow \tilde{\ell}\tilde{\chi}_1^0$	2 e, μ	0	Yes	20.3	$\tilde{\ell}$ 85-315 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$ ATLAS-CONF-2013-049
	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\ell}\nu(\tilde{\ell}\bar{\nu})$	2 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm$ 125-450 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$ ATLAS-CONF-2013-049
	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}\nu(\tilde{\tau}\bar{\nu})$	2 τ	-	Yes	20.7	$\tilde{\chi}_1^\pm$ 180-330 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\tau}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$ ATLAS-CONF-2013-028
	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\ell}\nu\tilde{\ell}(\tilde{\nu}\bar{\nu}), \tilde{\ell}\tilde{\nu}\tilde{\ell}(\tilde{\nu}\bar{\nu})$	3 e, μ	0	Yes	20.7	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^0$ 600 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$ ATLAS-CONF-2013-035
	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0 Z\tilde{\chi}_1^0$	3 e, μ	0	Yes	20.7	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^0$ 315 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0)=0, \text{ sleptons decoupled}$ ATLAS-CONF-2013-035
	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0 h\tilde{\chi}_1^0$	1 e, μ	2 b	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^0$ 285 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_1^0), m(\tilde{\chi}_1^0)=0, \text{ sleptons decoupled}$ ATLAS-CONF-2013-093
Long-lived particles	Direct $\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	Yes	20.3	$\tilde{\chi}_1^\pm$ 270 GeV	$m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)=160 \text{ MeV}, \tau(\tilde{\chi}_1^\pm)=0.2 \text{ ns}$ ATLAS-CONF-2013-069
	Stable, stopped \tilde{g} R-hadron	0	1-5 jets	Yes	22.9	\tilde{g} 832 GeV	$m(\tilde{\chi}_1^0)=100 \text{ GeV}, 10 \mu\text{s} < \tau(\tilde{g}) < 1000 \text{ s}$ ATLAS-CONF-2013-057
	GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$	1-2 μ	-	-	15.9	$\tilde{\chi}_1^0$ 475 GeV	$10 < \tan\beta < 50$ ATLAS-CONF-2013-058
	GMSB, $\tilde{\chi}_1^0 \rightarrow \tilde{\gamma}G$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	4.7	$\tilde{\chi}_1^0$ 230 GeV	$0.4 < \tau(\tilde{\chi}_1^0) < 2 \text{ ns}$ 1304.6310
	$\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow \tilde{q}\tilde{q}\mu$ (RPV)	1 μ , displ. vtx	-	-	20.3	\tilde{q} 1.0 TeV	$1.5 < c\tau < 156 \text{ mm}, \text{BR}(\mu)=1, m(\tilde{\chi}_1^0)=108 \text{ GeV}$ ATLAS-CONF-2013-092
RPV	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e + \mu$	2 e, μ	-	-	4.6	$\tilde{\nu}_\tau$ 1.61 TeV	$\lambda_{311}^e=0.10, \lambda_{132}=0.05$ 1212.1272
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e(\mu) + \tau$	1 e, $\mu + \tau$	-	-	4.6	$\tilde{\nu}_\tau$ 1.1 TeV	$\lambda_{311}^e=0.10, \lambda_{1(2)33}=0.05$ 1212.1272
	Bilinear RPV CMSSM	1 e, μ	7 jets	Yes	4.7	\tilde{q}, \tilde{g} 1.2 TeV	$m(\tilde{q})=m(\tilde{g}), c\tau_{LSP} < 1 \text{ mm}$ ATLAS-CONF-2012-140
	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^\pm \rightarrow ee\tilde{\nu}_\mu, e\mu\tilde{\nu}_e$	4 e, μ	-	Yes	20.7	$\tilde{\chi}_1^\pm$ 760 GeV	$m(\tilde{\chi}_1^0) > 300 \text{ GeV}, \lambda_{121} > 0$ ATLAS-CONF-2013-036
	$\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^\pm \rightarrow \tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$	3 e, $\mu + \tau$	-	Yes	20.7	$\tilde{\chi}_1^\pm$ 350 GeV	$m(\tilde{\chi}_1^0) > 80 \text{ GeV}, \lambda_{133} > 0$ ATLAS-CONF-2013-036
	$\tilde{g} \rightarrow \tilde{q}\tilde{q}\tilde{q}$	0	6-7 jets	-	20.3	\tilde{g} 916 GeV	$\text{BR}(t)=\text{BR}(b)=\text{BR}(c)=0\%$ ATLAS-CONF-2013-091
$\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$	2 e, μ (SS)	0-3 b	Yes	20.7	\tilde{g} 880 GeV	ATLAS-CONF-2013-007	
Other	Scalar gluon pair, $sgluon \rightarrow q\bar{q}$	0	4 jets	-	4.6	$sgluon$ 100-287 GeV	incl. limit from 1110.2693 1210.4826
	Scalar gluon pair, $sgluon \rightarrow t\bar{t}$	2 e, μ (SS)	1 b	Yes	14.3	$sgluon$ 800 GeV	ATLAS-CONF-2013-051
	WIMP interaction (D5, Dirac χ)	0	mono-jet	Yes	10.5	M^* scale 704 GeV	$m(\chi) < 80 \text{ GeV}, \text{limit of } < 687 \text{ GeV for D8}$ ATLAS-CONF-2012-147

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1σ theoretical signal cross section uncertainty.

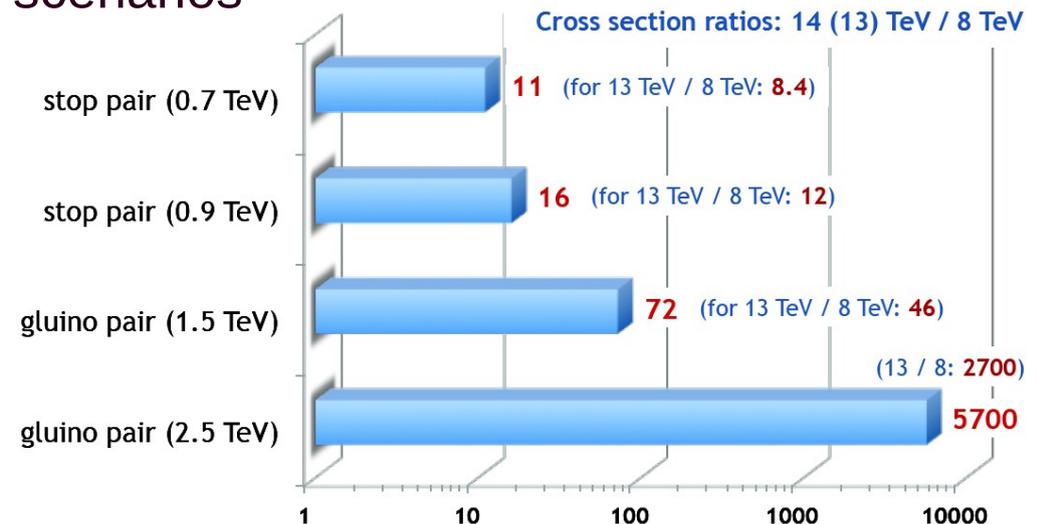
Summary and Outlook

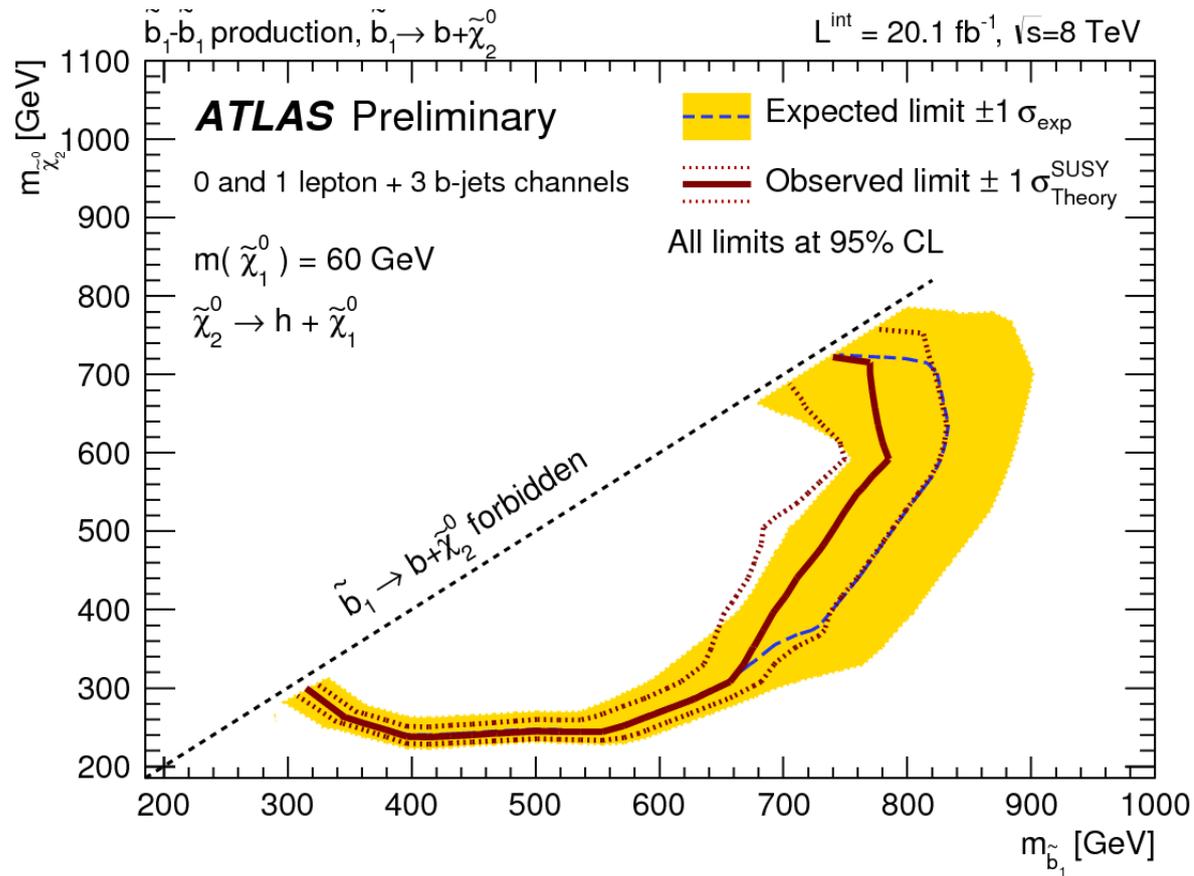
LHC run-1 dataset

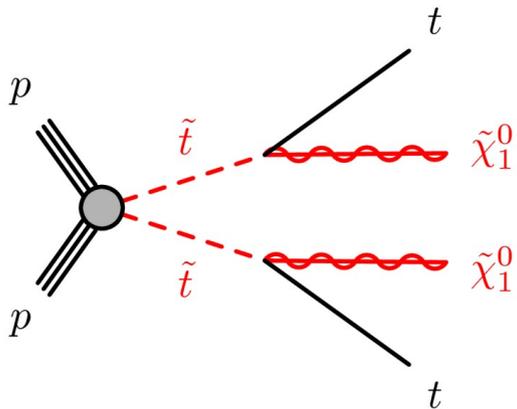
- ★ Broad SUSY programme developed
- ★ Effort to probe maximum area of SUSY parameter space possible
Using simplified models, pheno models and full models
- ★ Detailed and thorough searches, wide range of signatures covered
Focus on natural SUSY, strong production, RPV, long-lived SUSY searches
- ★ No sign of SUSY yet

What is next?

- ★ Increase sensitivity to difficult SUSY scenarios
- ★ Explore new channels,
probe more parameter space
- ★ **Prepare for $\sqrt{s} = 13$ TeV**
LHC run-2 in 2015
Increased sensitivity
to many SUSY scenarios







Consider SUSY processes with decays to two leptons + jets + LSP

Signal regions

- 2 leptons
- ≥ 2 jets
- High E_T^{miss}, m_{eff}

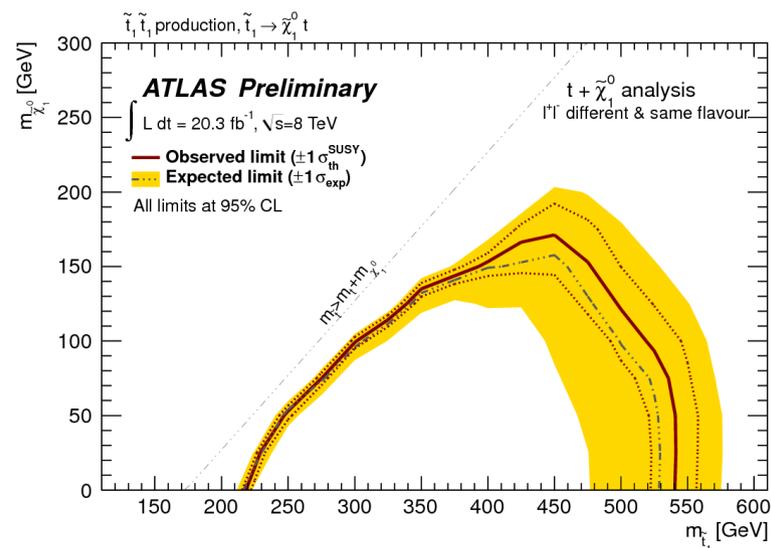
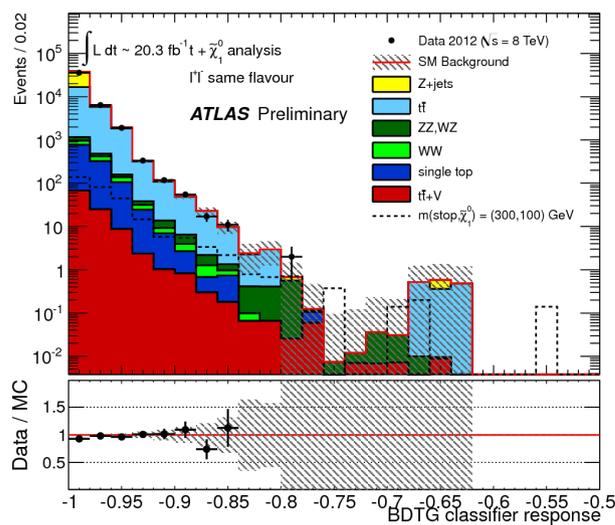
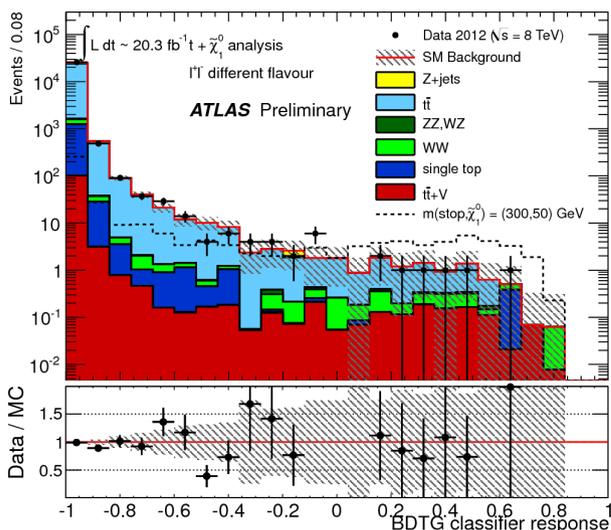
Train a BDT using

- E_T^{miss}, m_{ll}, m_{T2}(l₁, l₂, E_T^{miss})
- Δθ_{ll}, Δφ_{ll}
- Δφ(E_T^{miss}, j₁), Δφ(l₁, j₁)

Train Same Flavour (SF) and Different Flavour (DF) separately

Ttbar normalised to data in low-BDT control regions
Fake lepton background taken from data

4 SF and 7 DF signal regions

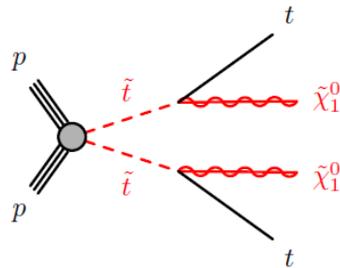


channel	SR ₁ ^{DF}	SR ₂ ^{DF}	SR ₃ ^{DF}	SR ₄ ^{DF}	SR ₅ ^{DF}	SR ₆ ^{DF}	SR ₇ ^{DF}
Observed events	9	3	12	5	3	2	1
Total (constrained) bkg events	4.7 ± 2.0	2.5 ± 1.9	11 ± 5	6.3 ± 2.5	1.0 ± 0.8	0.33 ^{+1.1} _{-0.33}	1.6 ± 1.4

High-energy LHC running in 2015 will significantly increase our sensitivity to many SUSY scenarios

- expect $\sim x10$ for 600 GeV stops, $\sim x200$ for 2 TeV gluinos

With $\sqrt{s} = 14$ TeV and 300 fb^{-1} , we expect to significantly improve our reach!



ATL-PHYS-PUB-2013-011

